

Contrasting interannual and multidecadal NAO variability

Article

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1 **Contrasting interannual and multidecadal NAO**
2 **variability**

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7 **Abstract** Decadal and longer timescale variability in the winter North At-
8 lantic Oscillation (NAO) has considerable impact on regional climate, yet it
9 remains unclear what fraction of this variability is potentially predictable. This
10 study takes a new approach to this question by demonstrating clear physical
11 differences between NAO variability on interannual-decadal (<30 year) and
12 multidecadal (>30 year) timescales. It is shown that on the shorter timescale
13 the NAO is dominated by variations in the latitude of the North Atlantic
14 jet and storm track, whereas on the longer timescale it represents changes in
15 their strength instead. NAO variability on the two timescales is associated
16 with different dynamical behaviour in terms of eddy-mean flow interaction,
17 Rossby wave breaking and blocking. The two timescales also exhibit different
18 regional impacts on temperature and precipitation and different relationships
19 to sea surface temperatures. These results are derived from linear regression
20 analysis of the Twentieth Century and NCEP-NCAR reanalyses and of a high-
21 resolution HiGEM General Circulation Model control simulation, with addi-
22 tional analysis of a long sea level pressure reconstruction. Evidence is presented
23 for an influence of the ocean circulation on the longer timescale variability of
24 the NAO, which is particularly clear in the model data. As well as provid-
25 ing new evidence of potential predictability, these findings are shown to have
26 implications for the reconstruction and interpretation of long climate records.

27 **Keywords** North Atlantic Oscillation · Jet variability · Atmosphere-ocean
28 interaction · Climate reconstructions

1 Introduction

As the leading pattern of atmospheric circulation variability over the North Atlantic, the North Atlantic Oscillation¹ (NAO) has a strong influence on surface climate across the Atlantic basin and beyond (Thompson and Wallace 2001). Interest in the NAO has been partly motivated by the prominence of its decadal-scale variability in winter (Stephenson *et al.* 2000). The increase of the winter NAO index from the 1960s to the 1990s gained particular attention (Hurrell 1995) but decadal variability is also evident in longer records of the NAO (Pinto and Raible 2012). In the likely absence of atmospheric memory from one winter season to the next, influences from other components of the climate system may have played a role. Evidence has been provided of possible influences such as the extratropical (Rodwell *et al.* 1999; Czaja and Frankignoul 1999; Mosedale *et al.* 2006; Gastineau and Frankignoul 2012) or tropical oceans (Hoerling *et al.* 2001; Selten *et al.* 2004; Greatbatch *et al.* 2012), the sea-ice (Deser *et al.* 2004; Bader *et al.* 2011) and also forcings acting via the stratosphere, such as changes in stratospheric water vapour (Joshi *et al.* 2006) or solar variability (Ineson *et al.* 2011). The associated potential for predictability of NAO variability continues to drive research in this area (Folland *et al.* 2012).

Much recent work has focused on shorter, intraseasonal timescales in attempts to understand the atmospheric dynamics underlying NAO variability,

¹ Or equivalently, the Arctic Oscillation or Northern Annular Mode (Feldstein and Franzke 2006).

50 following Feldstein (2003) and Benedict *et al.* (2004). It is clear that some of
51 the NAO variability on decadal timescales could arise from so-called climate
52 noise, in which seasonal sampling of the strong intraseasonal variability can
53 lead to apparent power on interannual and longer timescales (Wunsch 1999;
54 Feldstein 2000b; Schneider *et al.* 2003; Raible *et al.* 2005). Various statistical
55 methods have been applied to estimate the fraction of variance on interannual
56 and longer timescales which could be explained simply as climate noise. How-
57 ever, these methods differ widely in their findings (Feldstein 2000a; Keeley
58 *et al.* 2009; Franzke and Woollings 2011), so that the statistical significance of
59 low-frequency NAO variability, and hence the potential for seasonal-decadal
60 predictability is still unclear. In this paper we take a different and complemen-
61 tary approach, by searching for physical differences between NAO variability
62 on short and long timescales.

63 The NAO is essentially a description of the preferred structure of variabil-
64 ity in the North Atlantic eddy-driven jet stream (Thompson *et al.* 2002). This
65 deep tropospheric jet represents the net effect of westerly wind forcing by the
66 transient atmospheric eddies (Li and Wettstein 2012), and variations in its
67 strength and position affect regional temperatures and precipitation via vari-
68 ations in the prevailing westerly winds and associated storm tracks. The NAO
69 is usually defined via patterns in surface pressure or geopotential height, using
70 methods such as principal component analysis. Physical quantities such as the
71 latitude and speed of the jet are generally not separable by these methods
72 (Monahan *et al.* 2009), and the NAO reflects variations in both of these quan-

73 tities (Woollings *et al.* 2010). When height fields are linearly regressed onto
74 time series of the jet latitude and speed, both of the resulting spatial pat-
75 terns project onto the NAO pattern (Woollings and Blackburn 2012). Despite
76 this, the jet latitude and speed are clearly distinct, having different annual cy-
77 cles, power spectra and interannual variability (Woollings and Blackburn 2012;
78 Woollings *et al.* 2014). This suggests that variations in jet latitude and speed
79 have different physical mechanisms and drivers, and yet they are combined in
80 standard NAO analyses.

81 Here we highlight the contrasting nature of the jet variability associated
82 with the NAO on two different timescales, namely multidecadal and interannual-
83 decadal, with periods greater or less than 30 years respectively. This study
84 is related to other approaches which focus on the non-stationarity of the
85 NAO pattern over time (Jung *et al.* 2003; Lu and Greatbatch 2002; Raible
86 *et al.* 2006; Wang *et al.* 2012; Moore *et al.* 2013), or multi-decadal changes
87 in regime activity (Casty *et al.* 2005; Franzke *et al.* 2011). Other studies have
88 highlighted non-stationary relationships between the NAO and regional im-
89 pacts on temperature (Pozo-Vázquez *et al.* 2001; Haylock *et al.* 2007; Comas-
90 Bru and McDermott 2013), precipitation (Vicente-Serrano and López-Moreno
91 2008; Raible *et al.* 2014) and storm activity (Luo *et al.* 2011; Lee *et al.*
92 2012), and the timescale dependence shown here may help in interpreting
93 this non-stationarity. Finally, there is evidence for distinct patterns of ocean-
94 atmosphere variability on decadal/multi-decadal timescales in observations
95 and models (Deser and Blackmon 1993; Delworth and Mann 2000; Sutton and

96 Hodson 2003; Shaffrey and Sutton 2006), with non-stationarity or frequency-
97 dependence in the relationship between the NAO and sea surface temperatures
98 (Raible *et al.* 2001; Walter and Graf 2002; Raible *et al.* 2005; Alvarez-Garcia
99 *et al.* 2008). Hence we also examine the NAO-SST relationship on the two
100 timescales as a preliminary study of the associated ocean-atmosphere interac-
101 tion.

102 **2 Methods**

103 In this study we focus on the variability in wintertime (DJF) mean data from
104 atmospheric reanalyses. We use both the NCEP-NCAR reanalysis (Kalnay
105 and coauthors 1996) and the Twentieth Century reanalysis (20CR) (Compo
106 *et al.* 2011). The latter uses mean sea level pressure (MSLP) observations
107 only and takes an ensemble approach to quantify uncertainty, providing 56
108 estimates of atmospheric flow from 1871 to 2012. Unless otherwise stated, the
109 analyses presented here were performed individually for each of these ensemble
110 members and only averaged over the ensemble at the end of the analysis.

111 The NCEP-NCAR reanalysis is used over the period 1950-2012. Since this
112 only provides 62 winters of data this record is short to examine the multi-
113 decadal behaviour, but as will be shown results are qualitatively similar to
114 those in 20CR. This reanalysis is used particularly to investigate transient
115 features such as the storm track and the transient eddy fluxes. These quanti-
116 ties have been examined in 20CR but found to give unphysical results in the
117 low-frequency regressions, in particular at high latitudes. This is likely due

118 to issues in the data-sparse period before 1920 (see e.g. Krueger *et al.* (2013)
119 for a general discussion). The storm track is characterised using data filtered
120 with a 2-6 day Chebyshev recursive filter to select only the synoptic timescales
121 (Cappellini 1978).

122 We also make use of a 100 year present-day control simulation of the high-
123 resolution coupled General Circulation Model HiGEM (Shaffrey *et al.* 2009).
124 This has an atmospheric resolution of $0.833^\circ \times 1.25^\circ$ in longitude-latitude
125 with 38 levels, and an ocean resolution of $\frac{1}{3}^\circ$ with 40 levels. This model shows
126 improved simulation of the climatology and variability of North Atlantic cli-
127 mate compared to the standard resolution HadGEM1.2 (Shaffrey *et al.* 2009;
128 Keeley *et al.* 2012; Hodson and Sutton 2012). Some limited transient diagnos-
129 tics have been derived from the HiGEM simulation and these agree well with
130 the results from the NCEP-NCAR reanalysis.

131 The NAO was defined in all datasets as the leading Empirical Orthogonal
132 Function (EOF) and associated principal component time series of monthly
133 mean wintertime (DJF) mean sea level pressure over the Atlantic sector (90°W -
134 30°E , 30 - 90°N). In the 20CR data the NAO was calculated separately in each
135 ensemble member, and the resulting average spatial pattern is shown in Fig-
136 ure 5a of Woollings *et al.* (2014, W14 hereafter). The monthly NAO index was
137 averaged up to seasonal mean values for analysis. As described in W14, indices
138 of jet latitude and speed were derived using the zonal wind at 850 hPa. The
139 method essentially averages the daily zonal wind over 0 - 60°W and smoothes
140 it using a 10-day low pass filter before locating the maximum wind speed

141 (Woollings *et al.* 2010). The resulting daily values of jet latitude and speed
142 were averaged over each winter season to derive seasonal mean values.

143 To separate the different timescales we apply Empirical Mode Decomposi-
144 tion (EMD), as in Franzke and Woollings (2011), to the seasonal mean time
145 series of the NAO and jet indices. This approach empirically decomposes a
146 time series into Intrinsic Mode Functions (IMFs) of different average periods.
147 See Franzke and Woollings (2011) for more description and an example of the
148 method. Here we focus on two timescales: the interannual-decadal, formed by
149 isolating the IMFs with average periods less than 30 years, and the multi-
150 decadal, with IMF periods greater than 30 years. The sum of the two filtered
151 time series is exactly equal to the full unfiltered series. These two timescales
152 were chosen after experimentation to best represent the contrasting NAO be-
153 haviour (for example the 10-30 year band of timescales behave similarly to the
154 1-10 year band). Note that the general results presented here are reproduced
155 using other filtering methods such as running means, but the EMD results are
156 presented due to their smoothness and objectivity.

157 The general approach taken here is to linearly regress various fields onto the
158 NAO time series at the two different timescales. After averaging the monthly
159 data up to seasonal means, the NAO series is re-normalised so that the series of
160 winter mean values has a mean of zero and a standard deviation of one. Maps
161 therefore show the anomalies associated with one standard deviation of the
162 full unfiltered winter NAO. As described below, this makes the magnitudes
163 of the patterns on the two timescales comparable. However, it is important

164 to note that the long timescale anomaly patterns then have larger amplitude
165 than is experienced in practise.

166 **3 Jet Characteristics**

167 We begin by comparing the NAO and jet indices from 20CR in Figure 1. The
168 decadal variability of the NAO is clear, with high NAO values dominating
169 in the early and late twentieth century, and low NAO values dominating in
170 the middle of the century. In contrast, the jet latitude shows mostly interan-
171 nual variability, and as shown by W14 it is the jet speed which exhibits the
172 strongest decadal variability. W14 used a Monte Carlo statistical test to assess
173 the probability that the observed variability in the decadal means of these jet
174 series could arise from a white noise process. The results showed that this was
175 quite plausible for the jet latitude ($p=0.19$) but very unlikely for the jet speed
176 ($p=0.01$).

177 W14 also found that the jet latitude and speed series are uncorrelated ($r=-$
178 0.07), yet both are related to the NAO. This is shown in Figure 2 which cor-
179 relates these series with the NAO series on the different timescales obtained
180 using the EMD filtering. Both 20CR and NCEP-NCAR results are plotted,
181 with errorbars reflecting the uncertainty across the ensemble in 20CR. On
182 timescales shorter than 30 years the NAO is dominated by variations in jet
183 latitude. However, on the multidecadal timescale the reverse is true for 20CR
184 at least; the jet speed is more highly correlated with the NAO. NCEP-NCAR
185 shows high correlations for both jet speed and jet latitude on this timescale. If

186 the 20CR analysis is restricted to the time period of the NCEP-NCAR reanal-
187 ysis, this gives very similar correlations to the NCEP-NCAR data (asterisks
188 in Figure 2), suggesting that this difference is largely due to the short time
189 period.

190 These correlations suggest a change in the nature of NAO variability on
191 long timescales, with variations in jet speed becoming more important. This
192 impression is confirmed in Figure 3 which shows the 850 hPa zonal wind
193 anomalies associated with NAO variability on the two timescales. On the
194 shorter timescale the wind anomalies generally straddle the mean jet, indi-
195 cating a meridional shift, although the anomalies exhibit weaker meridional
196 tilt and are focused downstream of the mean wind maximum. On the mul-
197 tidecadal timescale, however, the anomalies overlie the mean jet, indicating
198 a clear increase in jet speed during positive NAO variations. The increase in
199 speed is also shifted towards the eastern end of the jet, highlighting an ex-
200 tension of the jet towards central Europe. Similar patterns are seen in the
201 NCEP-NCAR data, despite the difference in correlation on the long timescale
202 in Figure 2. Although the multidecadal anomalies are weaker in NCEP-NCAR,
203 the zonal wind is strengthened along the jet core as seen in 20CR. The same
204 behaviour is also seen in the HiGEM model control simulation, which suggests
205 the result is not a coincidence of the recent observed period. The similarity
206 of the patterns across the three datasets adds considerable confidence to the
207 result. These other datasets will be used in particular to analyse the storm
208 activity and ocean-atmosphere interaction on the two timescales, since these

209 are two aspects which have proved problematic in the analysis of the 20CR
210 data.

211 Following the jet analyses of W14, we performed a simple statistical test
212 of the decadal NAO variability against a white noise hypothesis. For each of
213 the 56 ensemble members, 1000 surrogate white noise NAO time series were
214 generated with the same standard deviation, and then decadal means were
215 calculated from these. The surrogate series were then used to determine the
216 likelihood of the observed level of variability in decadal means occurring from
217 the noise. This analysis showed that the decadal NAO variability in 20CR
218 is very unlikely to occur in a white noise model ($p=0.01$). We then applied
219 multiple linear regression to express the NAO as a linear combination of the jet
220 indices (which explained 71% of the NAO variance). This enabled us to remove
221 the influence of jet latitude and speed in turn and recalculate the likelihood
222 of the resulting decadal NAO variability. Removing the contribution of jet
223 latitude variations resulted in an NAO series which was still very unlikely in
224 the noise model ($p=0.01$), but removing the contribution of jet speed gave a
225 value of $p=0.19$, so that the resulting decadal NAO variability was no longer
226 significantly different from that expected from white noise. This exercise shows
227 that it is the variations in jet speed which are responsible for the elevated power
228 of the NAO on decadal timescales.

229 It would clearly be beneficial to verify the contrasting NAO behaviour
230 on the two timescales in a longer observational dataset. Several attempts have
231 been made to reconstruct atmospheric flow fields beyond the last century using

instrumental records, with considerable success. We have analysed the Küttel
et al. (2010) reconstruction of MSLP over Europe and the eastern North At-
lantic back to 1750, which uses both terrestrial pressure and marine wind data.
The result is that this dataset does not exhibit the distinct nature of multi-
decadal variability shown in the other datasets. Our analysis (described in
section 8) suggests that this may be at least partly an artefact of an assump-
tion of stationarity in the method used to derive the reconstruction. Given
the strong agreement between the other datasets, we conclude that the recon-
struction likely underestimates the timescale dependence of NAO variability.

4 Regional Impacts

NAO variations are of particular interest because of their strong influence on
regional surface climate. These connections have obvious societal impact and
are also often used to reconstruct indices of the NAO back in time. Figure 4
shows the patterns of near surface air temperature and precipitation associated
with the NAO. These impacts are notably different on the two timescales, espe-
cially over Europe. On interannual timescales this analysis gives the canonical
patterns of a quadrupole in temperature anomalies and a north-south dipole
in precipitation. On the decadal timescale, however, these patterns are shifted
south, so that both temperature and precipitation anomalies are focused on
western-central Europe. This southward shift is consistent with the role of the
East Atlantic pattern (the second EOF) which also describes changes in jet
speed and can be interpreted as acting to shift the NAO circulation pattern

254 north and south (Woollings *et al.* 2010). Figure 4 also shows strong differences
255 in Arctic climate on the two timescales, though this should be treated with
256 caution. If the analysis is restricted to the first 100 years of the period (1871-
257 1970) then the Arctic signal is greatly reduced. This suggests that this signal
258 arises from a correlation with the recent Arctic warming trend, which may be
259 unrelated to the NAO.

260 Some consideration of the variance of the NAO on the different timescales
261 should be taken in interpreting these impacts. This applies to all of the re-
262 gression maps shown in this paper. As described above, the NAO time series
263 was normalised and then split into the two components. The variance of the
264 full time series is 1.0, but the two components have very different variances:
265 0.90 for the interannual-decadal series and only 0.08 for the multidecadal se-
266 ries. This means that anomalies of the magnitude of those in the multidecadal
267 regressions are never realised in practise; the units of these regressions are
268 changes per standard deviation of the full NAO series, while the multidecadal
269 changes are much smaller than this. Despite this, the anomalies are of consid-
270 erable importance when compared to the level of variability on this timescale.
271 Over western Europe the multidecadal NAO regression accounts for up to
272 50% of the variance in decadal mean zonal wind, and similarly 30% of the
273 temperature and 60% of the precipitation variance.

274 5 Hemispheric connections

275 In this section we investigate whether the NAO on the two timescales has
276 different links to remote regions, in particular the Pacific. This is motivated
277 by discussion over the hemispheric or regional nature of the NAO (Wallace
278 2000) and also by evidence that interaction between the Atlantic and Pacific
279 sectors might be non-stationary or timescale dependent (Raible *et al.* 2001;
280 Castanheira and Graf 2003; Pinto *et al.* 2011; Lee *et al.* 2012).

281 Figure 5 shows the MSLP associated with NAO variations on the two
282 timescales. Despite the clear differences in jet behaviour, the MSLP patterns
283 in the Atlantic are only subtly different, indicating that MSLP anomalies are
284 hard to interpret in terms of jet characteristics. The most robust difference
285 is an eastward shift of the equatorward Atlantic centre of action on the long
286 timescale. The two reanalyses and the model are in good agreement over the
287 structure of variability on the shorter timescale, including a weak centre of
288 action in the eastern North Pacific. Climate models have historically overes-
289 timated the NAO teleconnection to the North Pacific (McHugh and Rogers
290 2005), but HiGEM appears to perform well in this regard (at least on the short
291 timescale). On the longer timescale there is little agreement between the three
292 datasets in the Pacific sector (though again the two reanalyses are similar if
293 only the NCEP-NCAR period is used; not shown). This lack of agreement be-
294 tween datasets limits the confidence we can have in hemispheric connections
295 on the long timescale.

296 To further investigate the Atlantic-Pacific links in the observations, we
297 analyse the storm track variability in the NCEP-NCAR reanalysis. The upper
298 level storm track is summarised by the mean of v'^2 at 200 hPa, which is re-
299 gressed onto the NAO in Figure 6. In the Atlantic sector the contrast in storm
300 track behaviour is very clear. The positive phase of the NAO is associated with
301 a northward shift and extension of the storm track on interannual-decadal
302 timescales, whereas it is associated with a strengthening of the storm track on
303 multidecadal timescales. These features do extend upstream into the Pacific
304 on both timescales. This is particularly clear on the multidecadal timescale,
305 where a strong increase in storm activity is seen over the eastern North Pacific.
306 This is consistent with the results of Lee *et al.* (2012) who found similar long
307 term changes in the Atlantic and Pacific storm tracks over recent decades.
308 Although weaker, the Pacific storm track signal on the shorter timescale is
309 again consistent with the Atlantic flow, since it indicates a weakening on the
310 southern side of the storm track. It is also consistent with studies which have
311 noted a latitudinal shift of the Pacific storm track accompanying an Atlantic
312 shift (Franzke *et al.* 2004; Strong and Magnusdottir 2008). Figure 6 also shows
313 corresponding results from HiGEM. As in the MSLP analysis, there seems to
314 be a Pacific-Atlantic storm track link on the shorter timescale which agrees
315 well with that in the reanalysis. On the long timescale there is good agreement
316 between the model and the reanalysis over the Atlantic sector, but not over
317 the Pacific.

318 To summarise, there is evidence that NAO variability on both timescales
319 has links to the Pacific sector. This is particularly clear on the shorter timescale,
320 where both MSLP and storm tracks show good agreement between the dif-
321 ferent datasets. In contrast, confidence in Atlantic-Pacific links on the longer
322 timescale is limited by the large differences between the three datasets.

323 **6 Eddy-mean flow interaction**

324 In this section we present further dynamical diagnostics of the atmospheric cir-
325 culation differences on the two timescales. We use the NCEP-NCAR reanalysis
326 for this analysis because of higher confidence in its transient fields. Figure 7
327 shows the vertical structure of the zonal wind anomalies along a section at
328 30°W . On both timescales the wind anomalies are equivalent barotropic with
329 maxima in the upper troposphere. The differences between the two timescales
330 seen at 850 hPa are clearly evident in the eddy-driven jet through the depth
331 of the troposphere, and are not just surface features. This suggests that the
332 wind anomalies are accompanied by changes in transient eddy driving of the
333 zonal flow, as expected from the storm track changes shown in Figure 6. Fig-
334 ure 7 also shows an interesting contrast in subtropical jet variability. This is
335 opposite to the eddy-driven jet, in that the subtropical jet strengthens and
336 weakens on the short timescale but shifts meridionally on the long timescale.

337 Transient baroclinic eddies influence the large-scale flow via both heat
338 and momentum fluxes. The top panels of Figure 8 show the lower tropo-
339 spheric transient eddy heat fluxes ($v'T'$). As in the other fields there is a clear

340 change between the timescales from a largely shifting pattern of variability to
341 a strengthening one. The transient eddy heat fluxes drive a residual overturn-
342 ing circulation with Coriolis torque acting to accelerate the westerly flow at
343 the latitude of the maximum in $v'T'$. Figure 8 then shows that the changes in
344 transient eddy heat fluxes act to support the zonal flow variations in each case,
345 helping to shift the surface westerlies in the interannual-decadal variability and
346 strengthen them in the multidecadal variability.

347 To summarise the effects of the transient eddy momentum fluxes, we follow
348 Raible *et al.* (2010) in calculating $\mathbf{E} \cdot \mathbf{D}$ where $\mathbf{E} = ((v'^2 - u'^2)/2, -u'v')$ is
349 similar to the E-vector of Hoskins *et al.* (1983) and $\mathbf{D} = (U_x - V_y, V_x + U_y)$
350 is the deformation vector of the time mean flow. Here u' and v' are the 2-6
351 day band-pass filtered wind components and U and V are the wind compo-
352 nents averaged over the relevant winter season. This diagnostic describes the
353 exchange of kinetic energy between the eddies and the background flow (Mak
354 and Cai 1989). Regressions of $\mathbf{E} \cdot \mathbf{D}$ on the NAO are shown in the lower panels
355 of Figure 8. The climatology of $\mathbf{E} \cdot \mathbf{D}$ features positive values over North Amer-
356 ica, implying that eddies grow there at the expense of the background state.
357 Over the Atlantic Ocean the climatology is negative, showing that the eddies
358 lose kinetic energy to the background state there. The regression patterns are
359 again very different on the two timescales. The multidecadal regression shows
360 a strengthening of the conversion from eddy to background state kinetic en-
361 ergy, consistent with increased eddy driving of the stronger jet stream. On
362 the interannual-decadal timescale the pattern is more complicated. While the

363 region of maximum eddy forcing is shifted northward by the anomalies, the
364 pattern also shows a meridional tightening of the eddy forcing over the ocean
365 and a general strengthening downstream. The upstream part of this pattern
366 may be related to the strengthening of the subtropical jet seen in Figure 7.

367 The effect of eddy forcing on the mean flow of the NAO is increasingly de-
368 scribed with regard to the breaking of transient Rossby waves (Benedict *et al.*
369 2004; Franzke *et al.* 2004; Rivière and Orlanski 2007; Martius *et al.* 2007; Kunz
370 *et al.* 2009; Archambault *et al.* 2010). Here we use the index of Barnes and
371 Hartmann (2012) which identifies wave breaking via the latitudinal overturn-
372 ing of vorticity contours. The index outputs the centroid of the wave breaking
373 event, counting each event once only (with a median lifetime of events of two
374 days) and discriminates between cyclonic and anticyclonic breaking based on
375 the morphology of the overturning region. Regressions of the occurrence of
376 wave breaking on the NAO are shown in Figure 9. On the interannual-decadal
377 timescale, the positive NAO is associated with a reduction in cyclonic wave
378 breaking on the poleward flank of the jet (to the south of Greenland) and an
379 increase in wave breaking of both types on the equatorward flank of the jet.
380 There is also a decrease in anticyclonic breaking in the subtropics, suggest-
381 ing that this region of wave breaking shifts north along with the jet. These
382 patterns are consistent with the picture that Rossby wave breaking acts to de-
383 accelerate the westerly winds locally, so that breaking on the equatorward side
384 pushes the jet polewards and vice versa (Gabriel and Peters 2008).

385 On the multidecadal timescale the positive NAO is instead associated with
386 increased wave breaking on both sides of the jet, which is consistent with
387 the strengthening and extension of the jet. Such large-scale conditions are
388 known to foster the occurrence of extreme windstorms over Western Europe
389 (Hanley and Caballero 2012; Gómara *et al.* 2014). The strongest signals are
390 increased cyclonic breaking to the north and increased anticyclonic breaking
391 to the south, though the two types of breaking also show weaker increases
392 on the opposite side of the jet. On both timescales the behaviour is therefore
393 consistent with Strong and Magnusdottir (2008), in that the latitude of the
394 breaking seems more important than its direction (e.g. the breaking on the
395 equatorward side of the jet may be cyclonic as well as anticyclonic).

396 Finally in this section, we analyse the relationship between the NAO and
397 blocking on both timescales. Blocking is a synoptic situation in which the west-
398 erly winds and storm tracks are blocked by a persistent, usually anticyclonic,
399 flow anomaly. Blocking is itself related to wave-breaking (Pelly and Hoskins
400 2003; Altenhoff *et al.* 2008), though the requirements of spatial scale and per-
401 sistence separate it from more transient wave breaking (Masato *et al.* 2009).
402 Here we define blocking using the index of Scherrer *et al.* (2006) which is a
403 two-dimensional extension of the classical Tibaldi and Molteni (1990) index.
404 A blocking pattern is identified at a point if 1) the meridional 500 hPa geopo-
405 tential height gradient is reversed and 2) the flow is westerly to the north of
406 the point, with a height gradient stronger than 10 m per degree of latitude. A
407 5-day persistence criterion is then applied at each gridpoint.

408 Figure 10 shows the regressions of blocking activity on the NAO. On the
409 interannual-decadal timescale, a positive NAO is associated with strongly re-
410 duced blocking over Greenland and the northern North Atlantic, as in Shabbar
411 *et al.* (2001); Croci-Maspoli *et al.* (2007) and Woollings *et al.* (2008). This is
412 also consistent with the reduction in cyclonic wave-breaking seen in Figure 9.
413 The increase in blocking to the south of the jet and over western Europe is
414 also consistent with previous studies (Davini *et al.* 2013). Essentially the jet
415 shifts southward due to blocking on its northern flank and northward due to
416 blocking on its southern flank.

417 On the multidecadal timescale, blocking anomalies are less strongly re-
418 lated to the NAO, with only very weak anomalies at high and low latitudes.
419 The implication is that the effect of blocking is largely to shift the jet stream
420 whereas transient wave breaking can act both to shift or strengthen the jet
421 depending on its position. Interestingly, there is an increase in blocking at the
422 jet exit over the British Isles, despite the strengthening of the westerly winds
423 there under the positive NAO. This may be a consequence of the storm track
424 changes, since a strong storm track upstream is favourable for block main-
425 tenance (Shutts 1983; Nakamura and Wallace 1993). Häkkinen *et al.* (2011)
426 demonstrated multidecadal variability of Atlantic-European blocking associ-
427 ated with Atlantic Ocean variability. Such basin-wide variations in blocking
428 do not appear in the analysis presented here, suggesting that the NAO is not
429 a good description of that variability.

430 **7 Ocean-atmosphere interaction**

431 The distinct physical characteristics of decadal NAO variability suggest an
432 influence external to the atmosphere. The slowly varying ocean circulation is
433 one potential forcing and a natural candidate is the Atlantic variability de-
434 scribed by the Atlantic Multidecadal Oscillation (Knight *et al.* 2005; Sutton
435 and Dong 2012). Correlations of the low-frequency NAO and jet indices with
436 a smoothed AMO index are given in Figure 1. The AMO index was obtained
437 from the NOAA ESRL website and was derived as in Enfield *et al.* (2001), in-
438 cluding detrending and smoothing with a 121 month smoother. Annual means
439 are plotted in Figure 1. The correlations show that the NAO is weakly anti-
440 correlated with the AMO and that most of this correlation likely comes from
441 the decadal variability in jet speed, which gives a slightly higher correlation of
442 -0.48. Another potential candidate for ocean forcing of decadal NAO variabil-
443 ity is the slow evolution of ocean temperatures in the tropical western Pacific
444 (Kucharski *et al.* 2006; Manganello 2008).

445 These potential links are investigated in Figure 11 by correlating winter
446 mean sea surface temperatures (SSTs) with the NAO at the two timescales.
447 The SST data comes from the HadISST dataset (Rayner *et al.* 2003) and the
448 correlation uses the ensemble mean NAO from the complete period of the
449 20CR data. Only gridpoints where the correlation is significant at the 95%
450 level are shown. On the short timescale the SSTs show the familiar tripole
451 pattern of anomalies which is largely a response to NAO variability. On the
452 longer timescale the SSTs show a more global pattern, with significant values

453 outside of the Atlantic basin. The North Atlantic is generally cool, as expected
454 from the negative correlation with the AMO, but the pattern is noisy and the
455 large values elsewhere are hard to interpret and may not be physically related.
456 The tropical western Pacific does, however, show a perturbed meridional SST
457 gradient, as found by Kucharski *et al.* (2006).

458 Taking a similar approach in correlating the SSTs with the time series of
459 jet speed from 20CR gives a clearer pattern on the multidecadal timescale,
460 comprising a cold subpolar gyre in the North Atlantic and warm anomalies
461 elsewhere, largely confined to the southern hemisphere (lower panels of Fig-
462 ure 11). Both of these features are reminiscent of AMO behaviour, and lends
463 support to the potential role of Atlantic Ocean circulation in influencing NAO
464 changes on the multidecadal timescale (e.g. Omrani *et al.* 2014; Peings and
465 Magnusdottir 2014). There is also indication of a potential influence from the
466 tropical Indian Ocean as suggested by Bader and Latif (2003).

467 In order to provide further evidence of an ocean influence, we examine the
468 ocean-atmosphere coupling in more detail in the HiGEM simulation, where
469 data availability and quality is not a limiting factor. Hodson and Sutton (2012)
470 previously investigated the North Atlantic ocean-atmosphere coupling in this
471 model with a focus on the shorter timescale. Figure 12 shows the correlation
472 of winter (DJF) mean SST with the long timescale NAO for the model. This
473 shows a distinct cold North Atlantic subpolar gyre; a pattern which is at least
474 qualitatively similar to that related to jet speed in the observations (Figure 11).
475 Figure 13a shows time series of the subpolar gyre temperature and the long

476 timescale NAO variability, which show reasonable covariability on the long
477 timescale.

478 In order to determine an influence of the ocean on the atmosphere, Fig-
479 ure 13b presents the heat content budget for the subpolar gyre region. The
480 sum of the heat contributions due to individual fluxes is shown by the dotted
481 black line, and this agrees well with the total heat content in the thick black
482 line, showing that all terms have been accounted for. The budget shows that
483 it is the ocean heat flux convergence (in red) which is driving the heat content
484 changes of the subpolar gyre, with the atmospheric fluxes (latent, sensible
485 heating) acting to damp the changes in heat content. This agrees with the
486 observational study of Gulev *et al.* (2013) which shows that ocean-atmosphere
487 surface heat fluxes are driven by the ocean on multidecadal timescales.

488 The variations in ocean heat flux convergence into the subpolar gyre may
489 be driven by a number of factors. Figure 13c demonstrates that these variations
490 are closely related to variations in the Meridional Overturning Circulation at
491 45°N, suggesting that variations in meridional ocean transport are responsi-
492 ble. These variations in turn arise in response to west to east ocean pressure
493 gradient across the Atlantic basin, which is dominated by ocean density vari-
494 ations on the deep western Atlantic boundary (Figure 13c: green line). The
495 resulting picture is that the decadal variations in the SST of the subpolar gyre
496 (see Figure 13a) are driven by changes in the MOC in the model, which are
497 in turn driven by variations in the density within the deep western boundary
498 current. Such density variations are ultimately generated at the ocean surface

499 in small regions of intense ocean cooling, such as the Labrador Sea (Marshall
500 and Schott 1999). NAO variability is thought to be a significant factor in driv-
501 ing variations in ocean cooling in these regions (Eden and Willebrand 2001).
502 Density anomalies generated by this process then slowly propagate southwards
503 along the western Atlantic boundary.

504 In summary, subpolar SST anomalies in HiGEM arise due to changes in
505 ocean heat convergence. These subpolar SST anomalies then in turn influence
506 the atmosphere, likely by changing the meridional temperature gradient and
507 hence the baroclinicity across the storm track. With a cold subpolar gyre the
508 meridional gradient is strengthened which is expected to lead to stronger storm
509 activity as seen in Figures 6 and 8. This in turn leads to increased acceleration
510 of the westerly flow (Figure 8) and a stronger jet. Evidence of this mechanism
511 of ocean influence on the atmosphere has been found in the natural variability
512 of other models (Gastineau and Frankignoul 2012), in the context of model
513 biases (Keeley *et al.* 2012) and in the response of models to climate change
514 (Woollings *et al.* 2012).

515 **8 Implications for long climate reconstructions**

516 In this section we analyse the Küttel *et al.* (2010) reconstruction of Atlantic/European
517 MSLP back to 1750 for evidence of contrasting NAO behaviour on the short
518 and long timescales. The NAO in the reconstructed data is defined as the
519 first EOF of winter mean MSLP over the region (0-40 °W, 20-70 °N), which
520 roughly comprises the North Atlantic portion of the data domain. The surface

521 geostrophic zonal wind u_g is then derived from the MSLP and this is regressed
522 onto the NAO at the two different timescales. The results are shown in the
523 top panels of Figure 14. In contrast to the other datasets, the differences be-
524 tween the two regression patterns are small, with the NAO largely describing
525 a jet shift on both timescales. To test whether this is due to the use of u_g
526 rather than u_{850} , we apply the same procedure to derive u_g from 20CR and
527 the results are shown in the middle panels of Figure 14. The results for 20CR
528 resemble the difference in u_{850} found between a jet shift on short timescales
529 and an increase in speed on longer timescales (Figure 3), suggesting that u_g
530 from the reconstruction should be capable of capturing this behaviour.

531 The reconstruction method used in Küttel *et al.* (2010) is the multivariate
532 principal component regression technique which relies on the assumption of
533 stationarity. EOFs of both observed MSLP fields and pressure-sensitive proxy
534 data (e.g., early measurements and documentary data such as ship log books)
535 are combined with a multiple linear regression technique for the observational
536 period to project local proxy information onto regional patterns. The linear
537 relation is then assumed to be stationary over time and used to reconstruct
538 MSLP fields further back in time. It is possible that the similarity of the
539 NAO regressions on the two timescales is a consequence of this assumption of
540 stationarity. To investigate this possibility we have performed a simple test on
541 the 20CR data, treating it in an analogous way to the reconstruction method.

542 Firstly an EOF analysis of the MSLP is performed on the last 30 years of
543 the 20CR data. Only the first four EOFs are retained, which explain 91% of

544 the variance of this sample. These EOFs are then used as a basis to truncate
545 the full 136 year dataset: a multiple linear regression technique considering the
546 anomaly maps for each year, the four leading EOFs and the associated prin-
547 cipal component time series are used to derive MSLP fields by only including
548 the projection on these four EOFs. The resulting pseudo-reconstructed MSLP
549 fields are then regressed on the NAO at the two timescales and finally u_g is
550 calculated from these. The results, shown in the lower panels of Figure 14, in-
551 dicate some differences between the two timescales but these are substantially
552 weaker than those in the original 20CR analysis (middle panels of Figure 14).
553 This is particularly clear over Europe where the latitude of the wind anomalies
554 is quite different in the middle panels but not in the lower panels. Retaining
555 more than four EOFs (e.g. 10 EOFs which comprises 99% of the variance) does
556 not significantly alter these findings (not shown). While this procedure is anal-
557 ogous but not identical to the technique of Küttel *et al.* (2010), it does suggest
558 that the lack of a distinct multidecadal NAO signature in the reconstruction
559 could be at least partly due to the assumption of stationarity in the method.
560 This effect could be compounded by the relatively low density of proxy data in
561 the jet stream region over the ocean considered in the reconstruction of Küttel
562 *et al.* (2010), and non-climatic noise intrinsic to proxy data.

563 9 Conclusions

564 This study shows that the multidecadal variability of the NAO represents very
565 different variations in atmospheric circulation from the interannual-decadal

566 variability. The faster variability is dominated by meridional shifts of the jet
567 stream and associated storm track, while the slower variability is dominated
568 by changes in the speed of the jet and the strength of the storm track.

569 Variations on both timescales are supported by forcing from the transient
570 eddies, but the nature of this forcing is different. The interannual-decadal
571 variations are associated with shifts of the transient eddy forcing and with the
572 occurrence of blocking weather patterns. Other work suggests this variability
573 represents variations in the occurrence of different synoptic circulation regimes
574 such as preferred jet positions (W14 and references therein). In contrast, the
575 multidecadal variability is associated with changes in strength of the eddy
576 forcing and with in-phase changes in the occurrence of transient Rossby wave
577 breaking on both sides of the jet.

578 The patterns of influence of the NAO on regional temperatures, wind
579 speeds and precipitation are different on the two timescales, and this has clear
580 implications for the interpretation of proxy or reconstructed records of past at-
581 mospheric variability in this region. The variations on multidecadal timescales
582 may not be well represented by the canonical NAO pattern, especially since
583 the shorter timescale variability dominates the variance of the NAO index. A
584 potential example of this has been given, by analysing a long MSLP recon-
585 struction. In contrast to the other datasets, this does not exhibit a difference
586 in NAO character on short and long timescales, and it is suggested that the
587 stationarity assumption commonly used in reconstruction methods is at least

588 partly responsible for this. These findings have implications for the interpre-
589 tation of climate reconstructions and long climate records.

590 These results also provide strong evidence for the presence of some forc-
591 ing on the decadal NAO from more slowly varying components of the climate
592 system than the atmosphere. Some evidence of links to Atlantic Ocean variabil-
593 ity were revealed, although other factors may also contribute. This evidence
594 is particularly clear in the HiGEM GCM, where variations in the Atlantic
595 Meridional Overturning Circulation lead to significant SST anomalies in the
596 subpolar gyre region which are then damped by the heat fluxes to the atmo-
597 sphere.

598 For the emerging discipline of decadal prediction these results are an en-
599 couraging sign of potential predictability of the winter NAO on multidecadal
600 timescales. Furthermore, the multidecadal component of NAO variability has
601 a clear and distinct influence on surface temperatures and precipitation, es-
602 pecially in Europe, so that decadal forecasts of this variability could be of
603 practical use. However, the contrasting behaviour on interannual and decadal
604 timescales suggests that the potential sources of skill may be different for
605 decadal forecasts than for seasonal forecasts.

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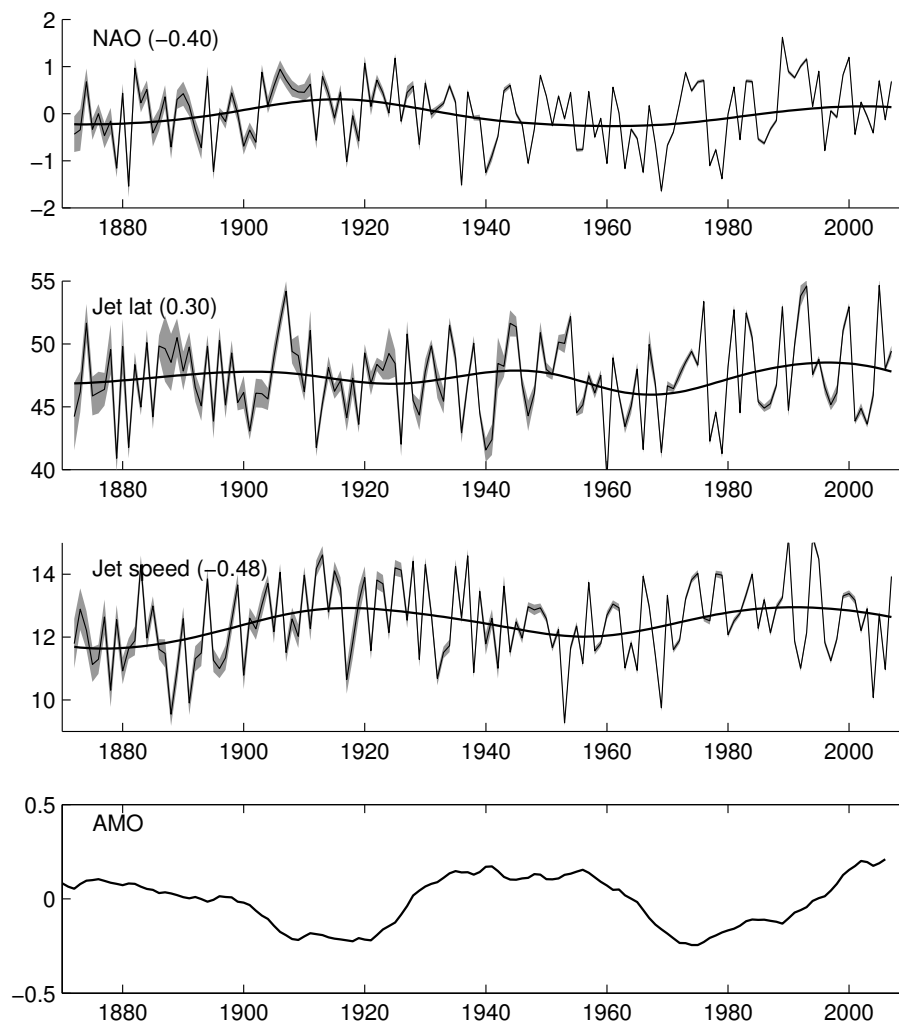


Fig. 1 Ensemble mean indices of the winter mean NAO, jet latitude and jet speed from 20CR, with the multidecadal (>30 year) component also shown. The shading indicates the ± 2 standard deviation range across the ensemble. The AMO is shown in the bottom panel, and in each other panel the correlation of the respective low-frequency timescale with this is given.

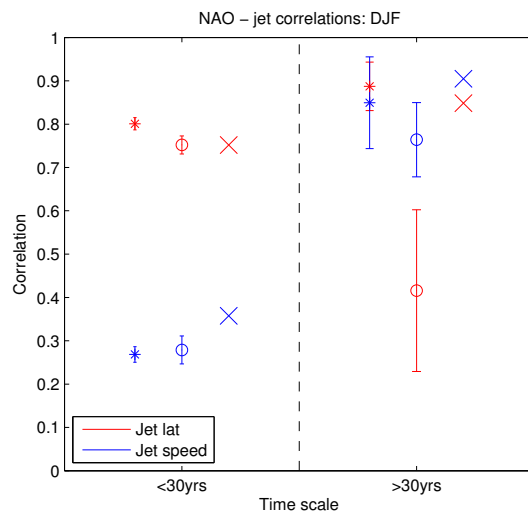


Fig. 2 Correlations of jet indices with the NAO on short and long timescales, showing NCEP-NCAR results as crosses and 20CR as circles, with errorbars giving the ± 2 standard deviation range across the ensemble. Asterisks give 20CR results for the NCEP-NCAR period.

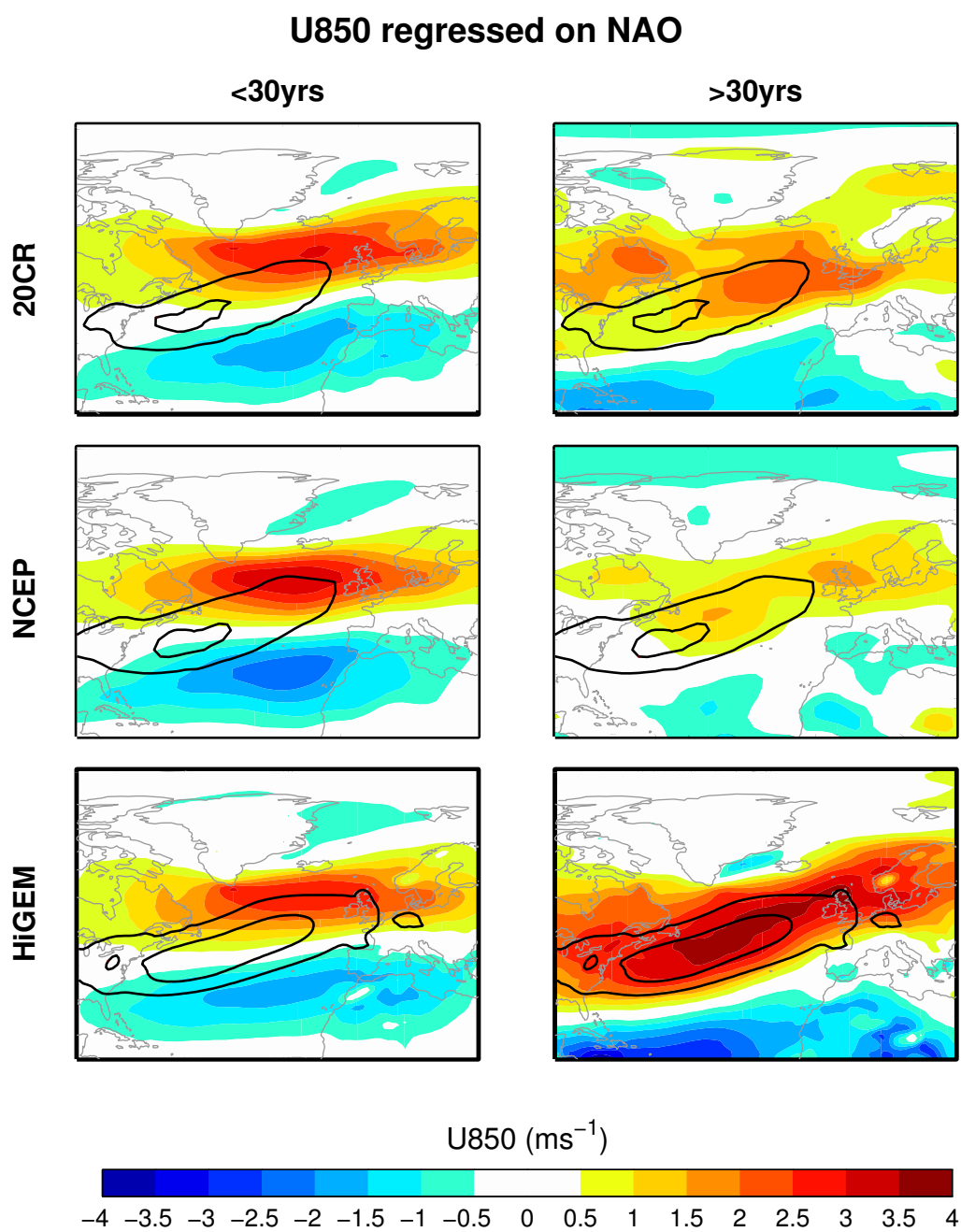


Fig. 3 Regression patterns of anomalies in 850hPa zonal wind on the NAO at the two timescales, using 20CR, NCEP-NCAR and HiGEM. The wind climatology is shown in black contours at 7.5 and 10.5 ms^{-1} .

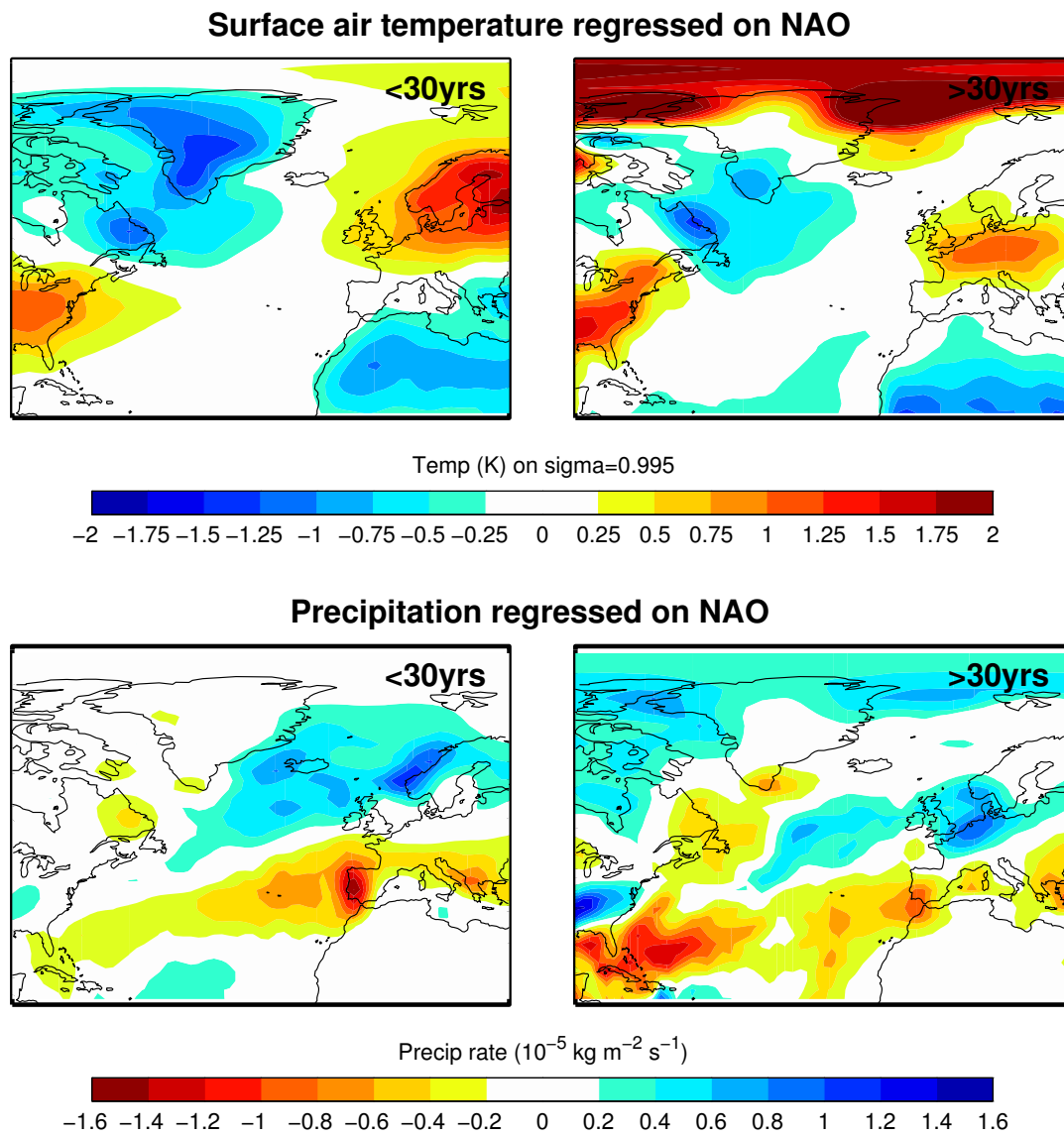


Fig. 4 Regression patterns of near surface air temperature (on the $\sigma = 0.995$ level) and precipitation on the NAO at the two timescales. This analysis was performed on the ensemble mean fields from 20CR.

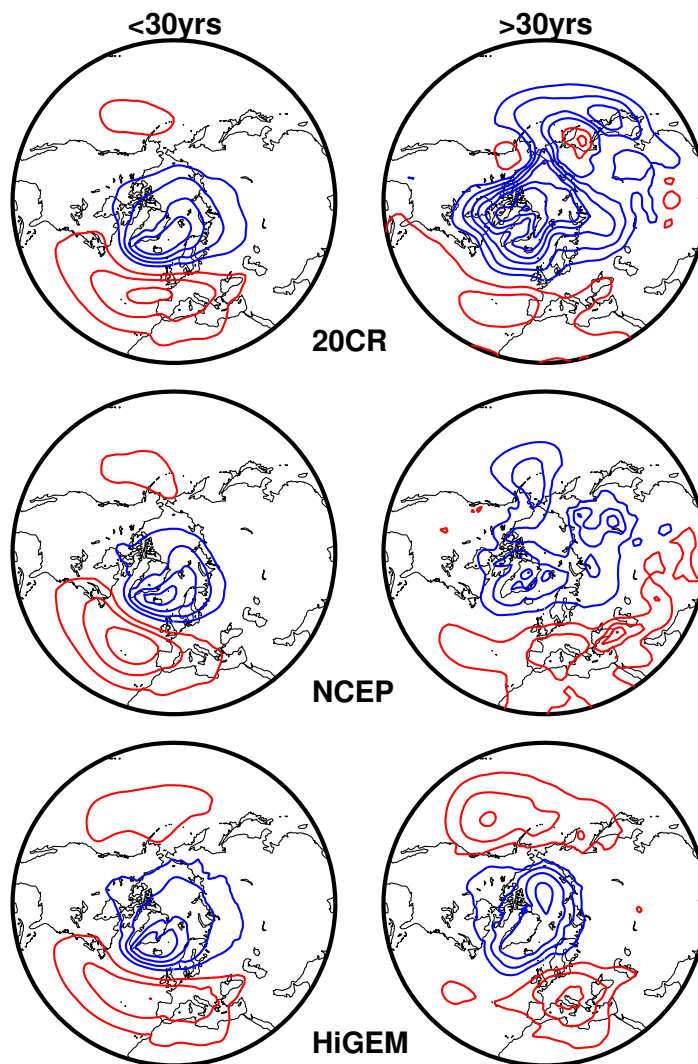


Fig. 5 MSLP from both reanalyses and from the HiGEM model regressed onto the NAO on both timescales. Contours are drawn every 1 hPa with negative contours in blue. All ensemble members are used for 20CR and the results averaged.

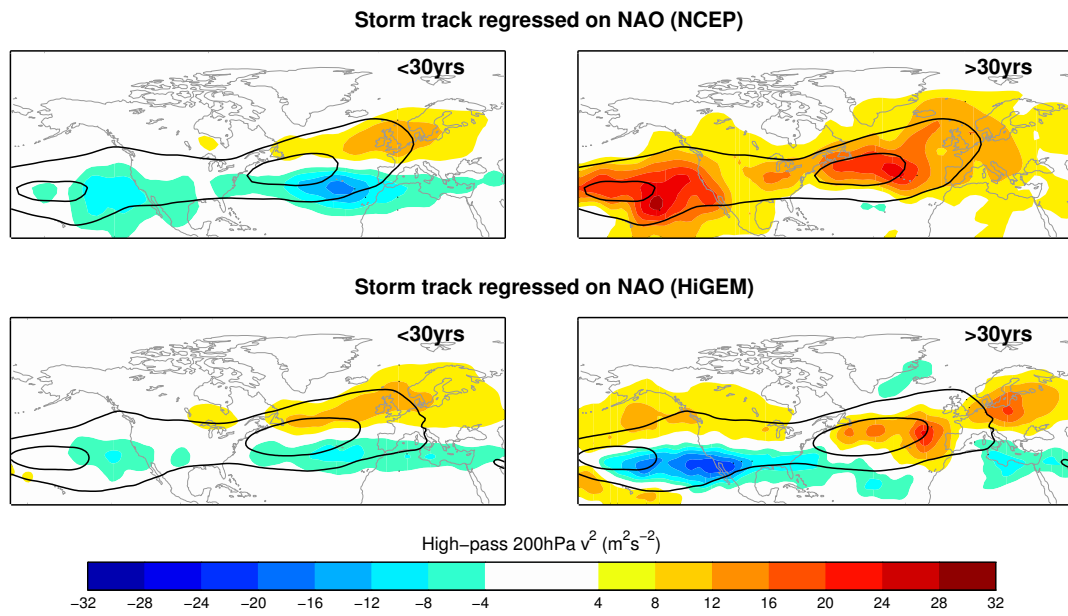


Fig. 6 Regressions of the storm track activity on the NAO in the NCEP-NCAR reanalysis and HiGEM, using the square of the meridional wind anomalies after applying a 2-6 day filter. The climatology is contoured at 60 and 100 m^2s^{-2} .

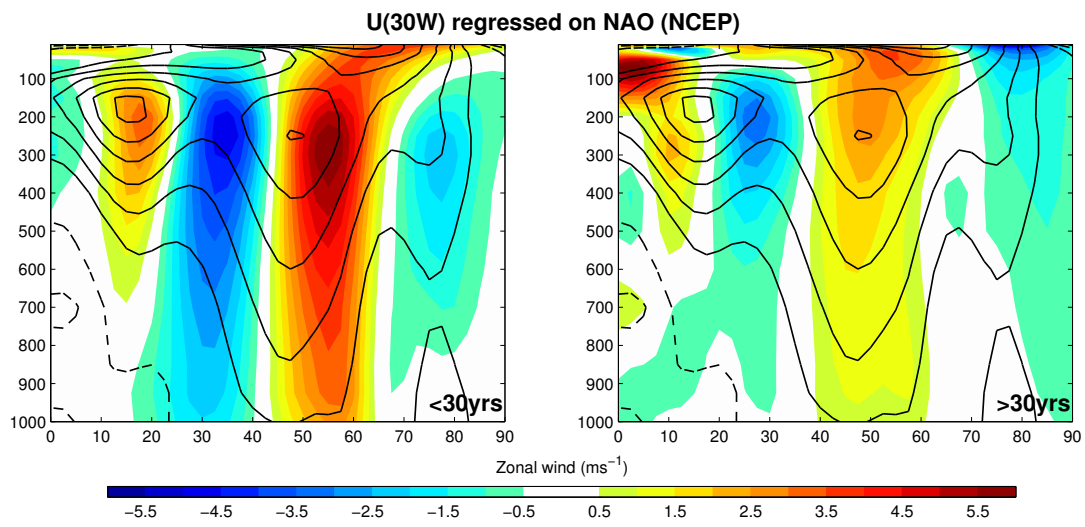


Fig. 7 Regression patterns of anomalies in zonal wind at 30°W on the NAO at the interannual and decadal timescales, using the NCEP reanalysis. The wind climatology is contoured every 5 ms^{-1} .

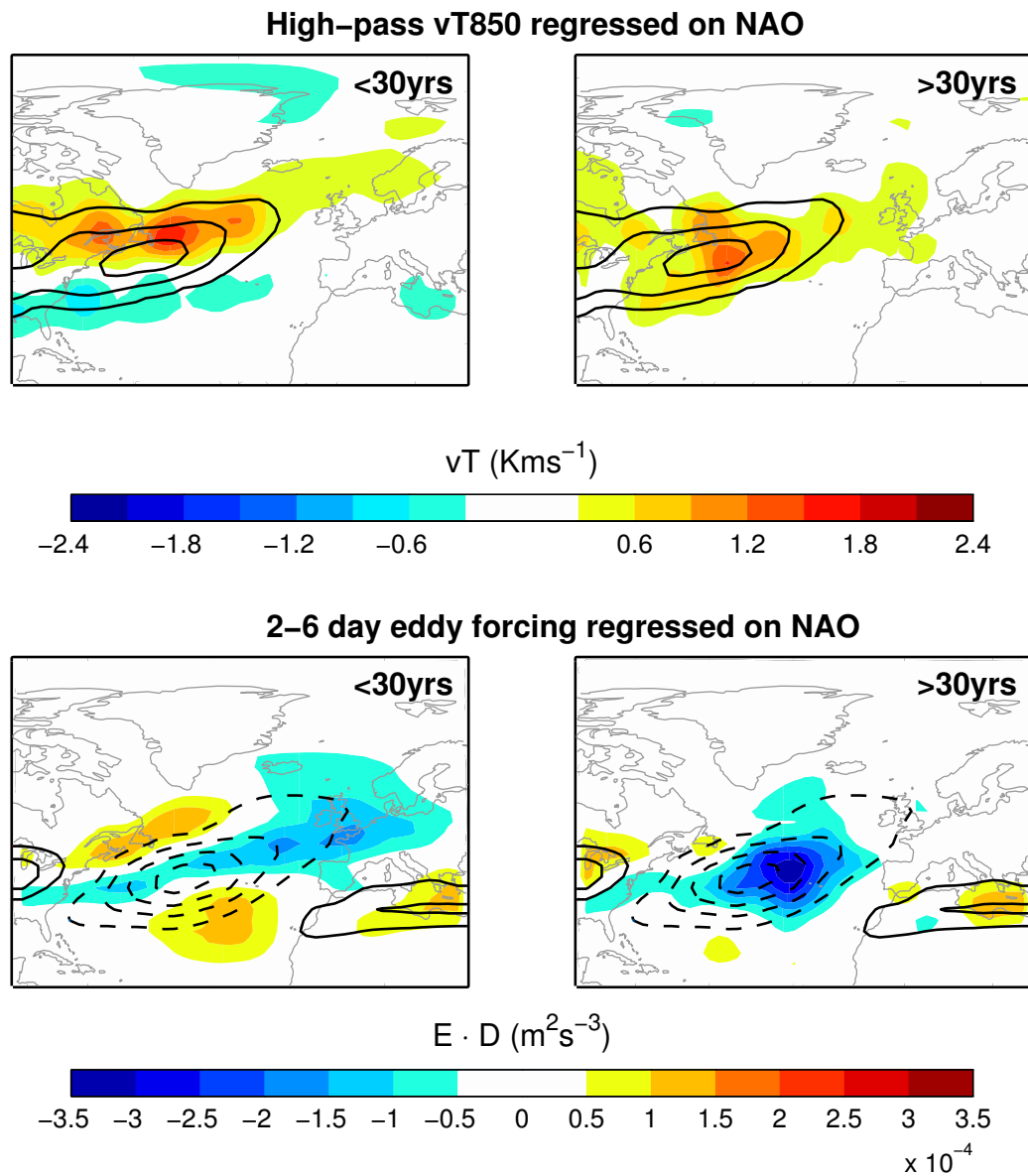


Fig. 8 Top: Regressions on the NAO of the 2–6 day $v'T'$ at 850 hPa. The climatology is contoured at 4, 7 and 10 Kms^{-1} . Bottom: Regressions on the NAO of the eddy forcing diagnostic $E \cdot D$ at 250 hPa. The climatology is contoured every $5 \text{m}^2\text{s}^{-3}$, with negative contours dashed and the zero contour omitted. In all cases the data is from the NCEP–NCAR reanalysis.

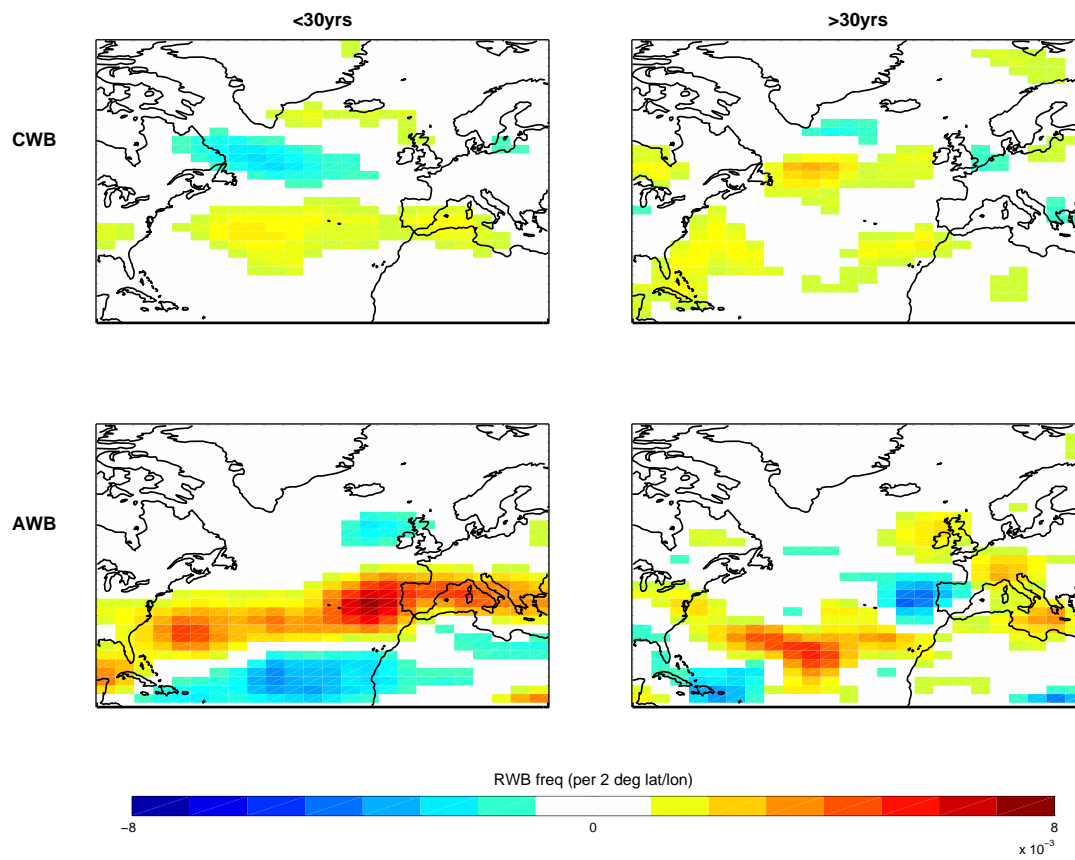


Fig. 9 Regressions on the NAO of the transient Rossby wave breaking occurrence, split into cyclonic (CWB) and anticyclonic (AWB). The NCEP-NCAR reanalysis is used.

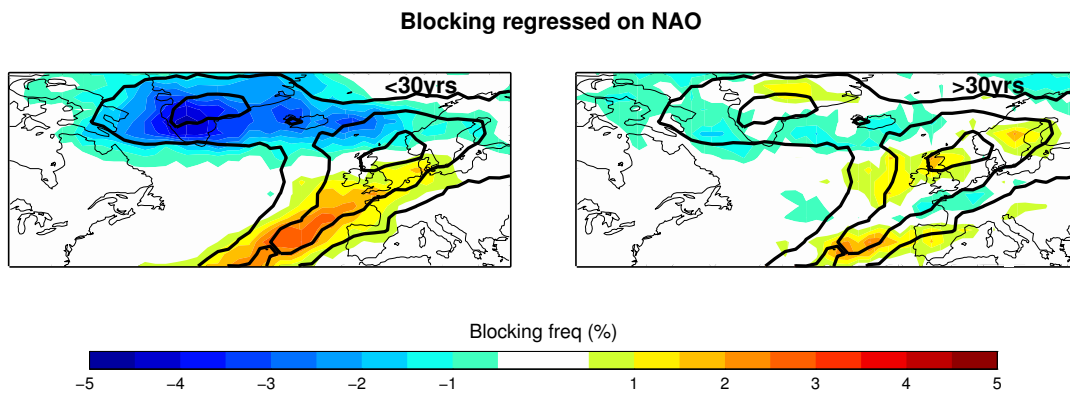


Fig. 10 Regressions of blocking occurrence on the NAO, using the NCEP-NCAR reanalysis.

The climatology is shown in black contours every 2%.

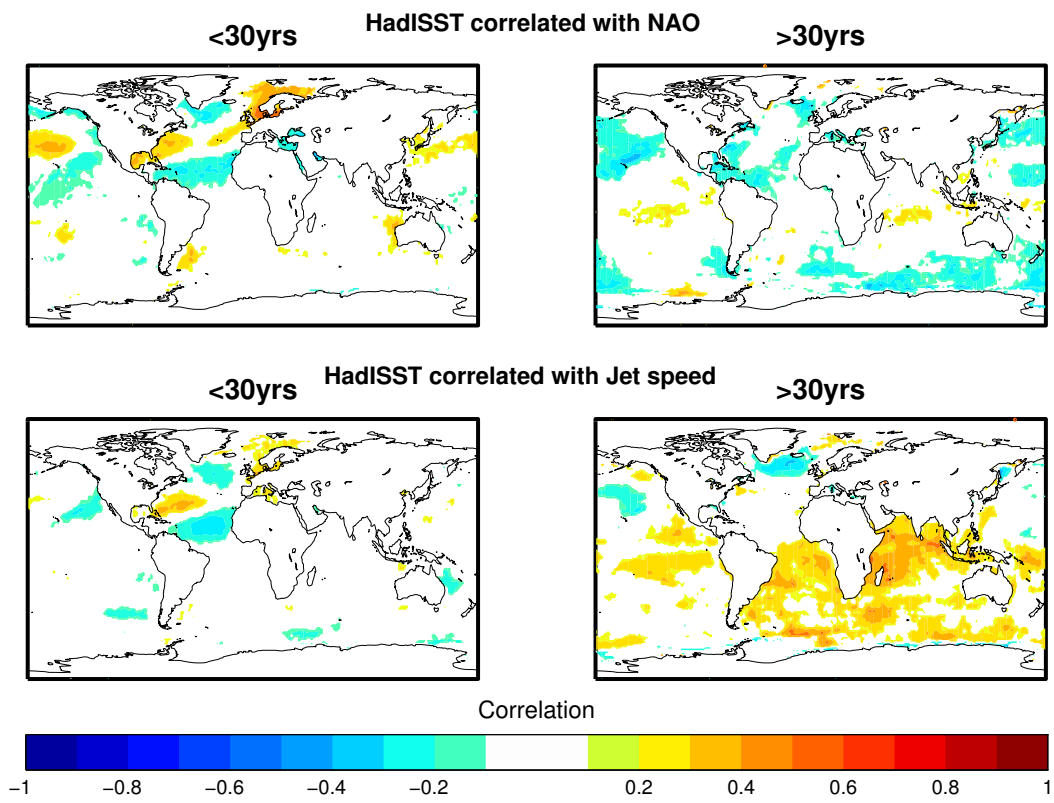


Fig. 11 Correlation value r of winter mean SSTs on the NAO (top) and jet speed (bottom) at the two timescales. Only values which are significant at the 95% level have been plotted.

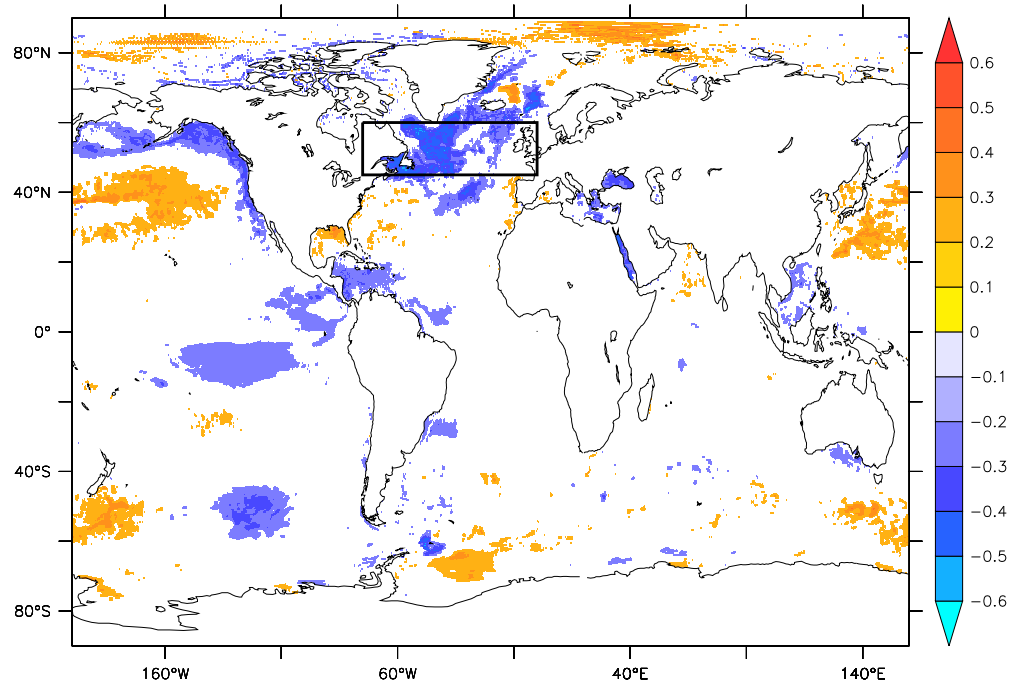


Fig. 12 Correlation value r of Sea Surface Temperature (SST) correlated with multidecadal NAO index in HiGEM. Shaded areas are significant at the 95% level ($p < 0.05$). Both SST and NAO were detrended before correlation.

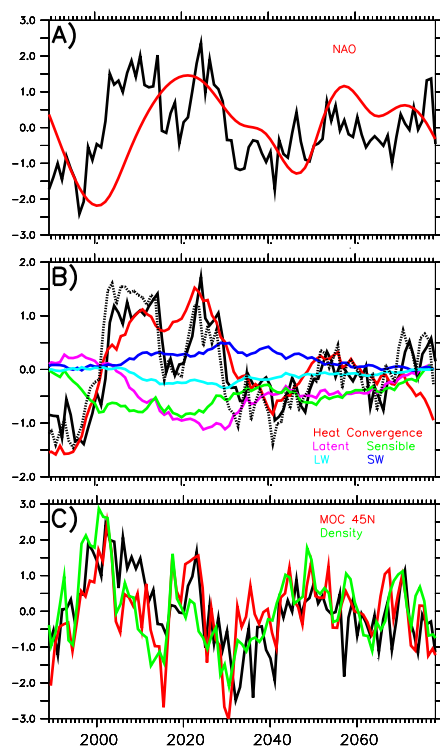


Fig. 13 Decadal NAO variability and Sub Polar Gyre heating in HiGEM Control simulation. A) Black: Mean Atlantic Sub Polar Gyre (SPG) Sea Surface Temperature (SST) ($75:0^{\circ}\text{W}$, $45:60^{\circ}\text{N}$ - box in Figure 12). Red: Detrended decadal component of the NAO in HiGEM multiplied by -1, extracted using EMD as for observations. Both indices have been standardized to have unit variance. B) Heat budget for the SPG region. Black solid: upper ocean heat content within the SPG region (0:500 m depth). Other lines - Heat content in the SPG due to: Ocean Heat convergence (Red), Surface Latent (Purple) and Sensible (Green) Heat fluxes and Longwave (Light Blue) and Shortwave (Dark Blue) surface radiation fluxes. All surface fluxes are defined positive *into* the ocean. Black dotted line: the sum of all contributions to the heat content. All indices have been detrended. Units are 10^7 PJ. Black (solid and dotted) lines have been multiplied by 2 to aid comparison with SPG SST in panel A. C) Black: Ocean Heat Convergence Flux into the SPG region ($45:60^{\circ}\text{N}$). Red: Atlantic Meridional Overturning Circulation (AMOC) at 45°N (AMOC is the integral of southward meridional ocean velocity between 1000:7000 m across the Atlantic Basin). Green: Mean Ocean Density on the Deep western Atlantic Boundary (1500:3000 m $59:58^{\circ}\text{W}$ $44:45^{\circ}\text{N}$). All indices have been detrended and standardized to have unit variance.

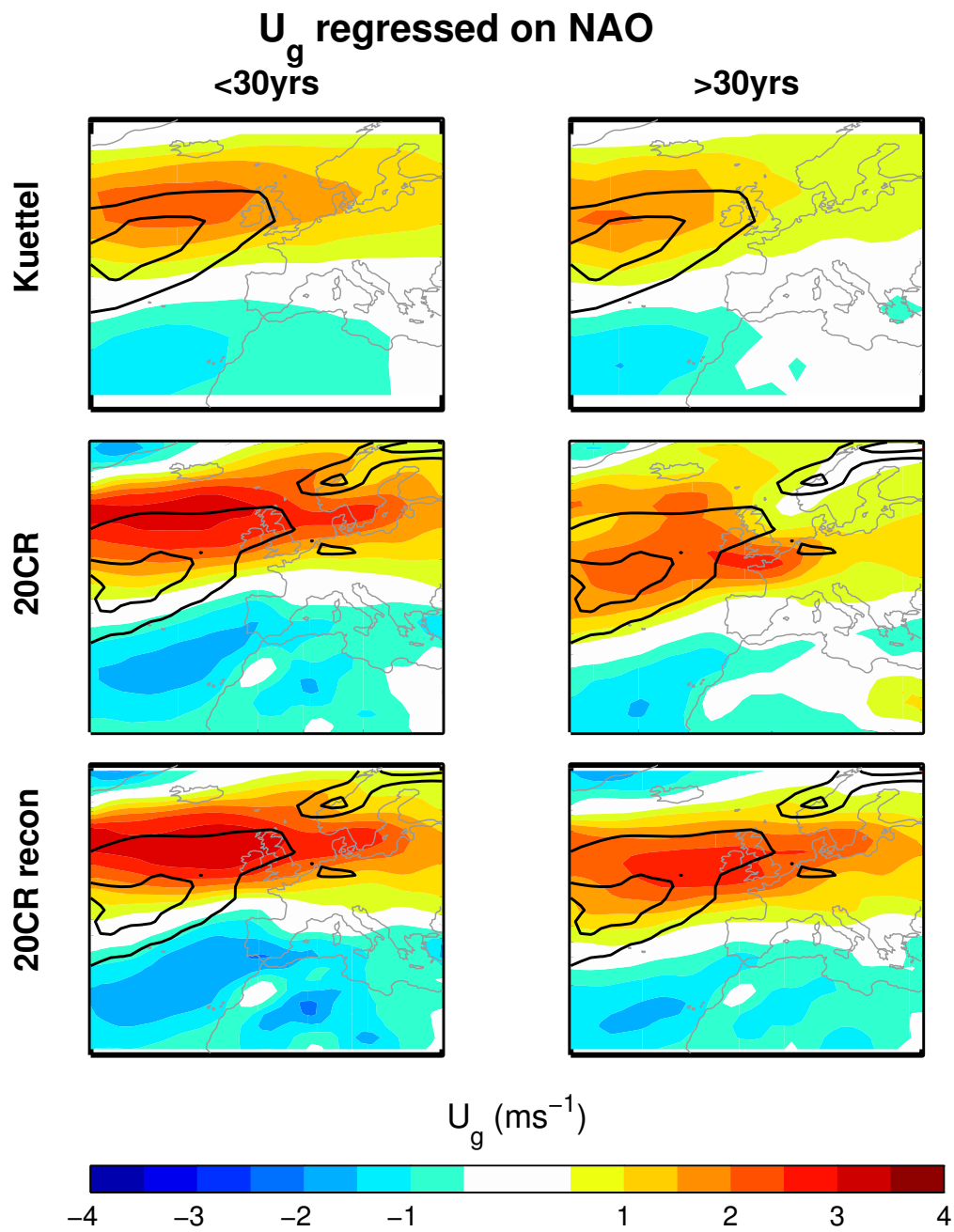


Fig. 14 As Figure 3 but showing the surface geostrophic zonal wind using the Küttel reconstruction and the 20CR data. The wind climatology is contoured in black at ± 5 , 7 ms^{-1} . The reconstruction only covers the region shown. See text for further details.