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Hagger, R. (2015) The eigenvalues of tridiagonal sign matrices are dense in the spectra of periodic tridiagonal sign operators. *Journal of Functional Analysis*, 269 (5). pp. 1563-1570. ISSN 0022-1236 doi: <https://doi.org/10.1016/j.jfa.2015.01.019>
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To link to this article DOI: <http://dx.doi.org/10.1016/j.jfa.2015.01.019>

Publisher: Elsevier

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The Eigenvalues of Tridiagonal Sign Matrices are Dense in the Spectra of Periodic Tridiagonal Sign Operators

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October 17, 2018

Abstract

Chandler-Wilde, Chonchaiya and Lindner conjectured that the set of eigenvalues of finite tridiagonal sign matrices (i.e. plus and minus ones on the first sub- and superdiagonal, zeroes everywhere else) is dense in the set of spectra of periodic tridiagonal sign operators on the usual Hilbert space of square summable bi-infinite sequences. We give a simple proof of this conjecture. As a consequence we get that the set of eigenvalues of tridiagonal sign matrices is dense in the unit disk. In fact, a recent paper further improves this result, showing that this set of eigenvalues is dense in an even larger set.

2010 Mathematics Subject Classification: Primary 15B35, 15A18; Secondary 47A10, 47B36

Keywords: sign matrix, tridiagonal, periodic operator, eigenvalues, spectrum

1 Introduction

Let $n, m \in \mathbb{N}$. For $k, l \in \{\pm 1\}^n$ we define the corresponding (finite) tridiagonal matrix

$$A_{fin}^{k,l} := \begin{pmatrix} 0 & l_1 & & \\ k_1 & \ddots & \ddots & \\ & \ddots & \ddots & l_n \\ & & k_n & 0 \end{pmatrix} \in \mathbb{C}^{(n+1) \times (n+1)}.$$

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Similarly we define the m -periodic tridiagonal operator

$$A_{per}^{k,l} := \left(\begin{array}{ccccccc} \ddots & & \ddots & & & & \\ & \ddots & & & & & \\ & & 0 & l_m & & & \\ & & k_m & 0 & l_1 & & \\ & & & k_1 & \ddots & \ddots & \\ & & & & \ddots & \ddots & l_m \\ & & & & & k_m & 0 & l_1 \\ & & & & & & k_1 & 0 & \ddots \\ & & & & & & & \ddots & \ddots \end{array} \right) \in \mathcal{L}(\ell^2(\mathbb{Z}))$$

for $k, l \in \{\pm 1\}^m$. In the following we will compare the spectra, denoted by $\text{sp}(\cdot)$, of the finite matrices with the spectra of the periodic operators as follows. In accordance with the notation in [3] and [4], we define the sets σ_∞ and π_∞ :

$$\sigma_\infty := \bigcup_{n \in \mathbb{N}} \bigcup_{k, l \in \{\pm 1\}^n} \text{sp}(A_{fin}^{k,l}), \quad \pi_\infty := \bigcup_{m \in \mathbb{N}} \bigcup_{k, l \in \{\pm 1\}^m} \text{sp}(A_{per}^{k,l}).$$

These sets are of interest because they are closely related to the spectrum of the Feinberg-Zee Random Hopping Matrix, denoted by Σ , see e.g. [3]-[5], [7], [9]-[11]. It was shown in [4] that

$$\sigma_\infty \subset \pi_\infty \subset \Sigma$$

holds. Note that these inclusions have to be proper because σ_∞ is a countable set, π_∞ is an uncountable null set (w.r.t. 2-dimensional Lebesgue measure) and Σ contains the closed unit disk (see below). The authors of [4] conjectured that σ_∞ is actually dense in Σ . We make a step in this direction by proving that σ_∞ is at least dense in π_∞ .

Moreover, it was shown in [5] that π_∞ is dense in the unit disk $\mathbb{D} := \{z \in \mathbb{C} : |z| \leq 1\}$, i.e. $\text{clos}(\pi_\infty) \cap \mathbb{D} = \mathbb{D}$. This result was improved to a larger set in [10]. Our new result here implies that σ_∞ is dense in this particular set as well. On the other hand, using the numerical range, one can show that Σ is contained in the square with vertices $\{2, 2i, -2, -2i\}$ ([4], [9]). In [9] this upper bound was further improved using second order numerical ranges. Furthermore, π_∞ has an infinite number of symmetries that also carry over to σ_∞ (see [4] for rotations and reflections, [5] for square roots and [10] for higher roots). Combining all these results mentioned above, we have a pretty accurate description of the set σ_∞ . However, many questions still remain open. For example the seemingly fractal boundary is yet unexplained (cf. Figure 1, borrowed from [4]). Some comments on this particular issue can be found in [10].

2 Three Lemmas and a Theorem

As mentioned in the introduction it was shown in [4] that σ_∞ is a subset of π_∞ . The following theorem confirms the conjecture by Chandler-Wilde, Chonchaiya and Lindner ([3],[4]) that σ_∞ is even dense in π_∞ .

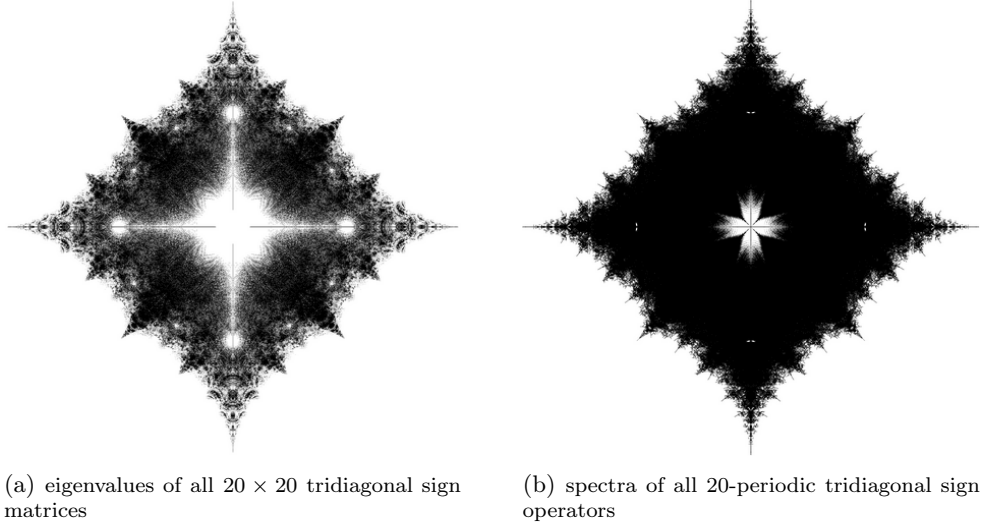


Figure 1

Theorem 1. σ_∞ is a dense subset of π_∞ .

First observe that it is enough to consider pairs (k, l) with $l = (1, \dots, 1)$ since every $A_{fin}^{k, l}$ is unitarily equivalent, via a diagonal so-called gauge transform ([12]), to some $A_{fin}^{\tilde{k}, \tilde{l}}$ with $\tilde{k} \in \{\pm 1\}^n$ and $\tilde{l} = (1, \dots, 1)$. Similarly, every $A_{per}^{k, l}$ is unitarily equivalent to some $A_{per}^{\tilde{k}, \tilde{l}}$ with $\tilde{k} \in \{\pm 1\}^m$ and $\tilde{l} = (1, \dots, 1)$. Thus we omit the index l from now on.

In the proof of Theorem 1 we will use three auxiliary lemmas. The first one is well-known in the theory of block Laurent operators. It generalizes the symbol calculus for Laurent operators to periodic operators. Just observe that an m -periodic operator can be interpreted as a block Laurent operator with block size $m \times m$. Using a block Fourier transform it is then not hard to see that the operator A_{per}^k is unitarily equivalent to the (generalized) multiplication operator $M_{a^k} : L^2([0, 2\pi), \mathbb{C}^m) \rightarrow L^2([0, 2\pi), \mathbb{C}^m)$ defined by the matrix valued function

$$a^k : [0, 2\pi) \rightarrow \mathbb{C}^{m \times m}, \quad \varphi \mapsto a^k(\varphi) := \begin{pmatrix} 0 & 1 & & k_m e^{i\varphi} \\ k_1 & \ddots & \ddots & \\ & \ddots & \ddots & 1 \\ e^{-i\varphi} & & k_{m-1} & 0 \end{pmatrix}.$$

The first lemma is then an easy consequence. For a proof and more material on this subject see e.g. [2, Theorem 2.93], [6, Theorem 4.4.9] or [8, Chap. VIII, Theorem 5.1].

Lemma 1. Let $m \in \mathbb{N}$ and $k \in \{\pm 1\}^m$. Then we have

$$\text{sp}(A_{per}^k) = \{\lambda \in \mathbb{C} : \det(a^k(\varphi) - \lambda I_m) = 0 \text{ for some } \varphi \in [0, 2\pi)\}.$$

In the case of tridiagonal periodic operators, only the constant term of $\det(a^k(\varphi) - \lambda I_m)$ (as a polynomial of λ) depends on φ as the next lemma shows. This leads to the fact that the spectrum of every periodic operator A_{per}^k can be written as $p^{-1}([-2, 2])$ for some polynomial p (cf. Figure 2). For an explicit formula of p we refer to [10].

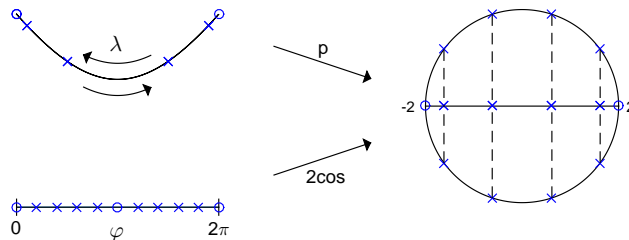


Figure 2: a sketch of the maps involved in the computation of the spectrum of a periodic operator

Lemma 2. *Let $m \in \mathbb{N}$ and $k \in \{\pm 1\}^m$. Then there exists a polynomial $p: \mathbb{C} \rightarrow \mathbb{C}$ of degree m (depending on k) such that*

$$\det(a^k(\varphi) - \lambda I_m) = (-1)^m \left(p(\lambda) - e^{i\varphi} \prod_{j=1}^m k_j - e^{-i\varphi} \right)$$

for all $\varphi \in [0, 2\pi)$.

Proof. Using the Leibniz formula

$$\det(a^k(\varphi) - \lambda I_m) = \sum_{\tau \in S_m} \text{sign}(\tau) \prod_{i=1}^m a_{i, \tau_i}^k(\varphi),$$

it is easily seen that the only surviving term containing $e^{i\varphi}$ but not $e^{-i\varphi}$ is $(-1)^{m+1} e^{i\varphi} \prod_{j=1}^m k_j$ and similar for $e^{-i\varphi}$. Thus the assertion follows. Alternatively one could also apply Laplace's formula twice to get the same result (see [10] for details). \square

The third lemma is a discrete version of Lemma 1. Although the result is well-known (see e.g. [1, Section 2.1]), we include a short proof for the reader's convenience.

Lemma 3. *Let $n, m \in \mathbb{N}$ and $A, B, C \in \mathbb{C}^{m \times m}$. Furthermore, denote by $e^{i\xi_1}, \dots, e^{i\xi_n}$ the n -th roots of unity, i.e. $\xi_j := \frac{2j}{n}\pi$ for $j \in \{1, \dots, n\}$. Then the following block matrices are unitarily*

equivalent:

$$T_1 := \begin{pmatrix} B & C & & A \\ A & \ddots & \ddots & \\ & \ddots & \ddots & C \\ C & & A & B \end{pmatrix} \in \mathbb{C}^{nm \times nm},$$

$$T_2 := \begin{pmatrix} Ae^{i\xi_1} + B + Ce^{-i\xi_1} & & & \\ & \ddots & & \\ & & \ddots & \\ & & & Ae^{i\xi_n} + B + Ce^{-i\xi_n} \end{pmatrix} \in \mathbb{C}^{nm \times nm}$$

$$= \text{diag}(Ae^{i\xi_1} + B + Ce^{-i\xi_1}, \dots, Ae^{i\xi_n} + B + Ce^{-i\xi_n}).$$

Proof. With the help of the Kronecker product \otimes , we can write $T_1 = P \otimes A + I_n \otimes B + P^* \otimes C$, where

$$P := \begin{pmatrix} 0 & & & 1 \\ 1 & \ddots & & \\ & \ddots & \ddots & \\ & & 1 & 0 \end{pmatrix} \in \mathbb{C}^{n \times n}.$$

P is unitarily equivalent to the diagonal matrix

$$D := \begin{pmatrix} e^{i\xi_1} & & \\ & \ddots & \\ & & e^{i\xi_n} \end{pmatrix} \in \mathbb{C}^{n \times n}.$$

Thus T_1 is unitarily equivalent to $D \otimes A + I_n \otimes B + D^* \otimes C$, which is exactly T_2 . \square

Proof. (Theorem)

That $\sigma_\infty \subset \pi_\infty$ holds was proven in [4, Theorem 4.1]. So let $m \in \mathbb{N}$, $k \in \{\pm 1\}^m$ and consider the operator A_{per}^k . W.l.o.g. we can assume that the number of -1 's in k is even because we can always double the period without changing the operator. This implies that we have

$$\det(a^k(\varphi) - \lambda I_m) = (-1)^m (p(\lambda) - 2 \cos(\varphi)) \quad (1)$$

in Lemma 2. Using Lemma 1, we get $\text{sp}(A_{per}^k) = p^{-1}([-2, 2])$. This implies, in particular, that the set

$$S_\infty^k := \{\lambda \in \mathbb{C} : \det(a^k(\varphi) - \lambda I_m) = 0 \text{ for some } \varphi \in \pi(\mathbb{Q} \setminus \mathbb{Z})\} = p^{-1}(2 \cos(\pi(\mathbb{Q} \setminus \mathbb{Z})))$$

is dense in $\text{sp}(A_{per}^k)$, cf. again Figure 2. Let $n \in \mathbb{N}$ and $\xi_j := \frac{2j}{n}\pi$ for $j \in \{1, \dots, n\}$ as above. We will show that for every $n \in \mathbb{N}$ there exists a finite matrix $A_{fin}^l \in \mathbb{C}^{(nm-1) \times (nm-1)}$ such that

$$S_n^k := \bigcup_{j \in \{1, \dots, n-1\} \setminus \{\frac{n}{2}\}} \text{sp}(a^k(\xi_j)) \subset \text{sp}(A_{fin}^l) \subset \sigma_\infty. \quad (2)$$

where $\sigma \in (0, 1)$, as considered in [5] or [9] for example, the conclusion fails. This was of course to be expected because a finite matrix

$$\begin{pmatrix} 0 & l_1 & & & \\ k_1 & \ddots & \ddots & & \\ & \ddots & \ddots & l_n & \\ & & k_n & & 0 \end{pmatrix} \in \mathbb{C}^{(n+1) \times (n+1)}, k \in \{\pm\sigma\}^n, l \in \{\pm 1\}^n$$

is always similar (diagonal gauge transform, see [12]) to a matrix

$$\begin{pmatrix} 0 & \tilde{l}_1 & & & \\ \tilde{k}_1 & \ddots & \ddots & & \\ & \ddots & \ddots & \tilde{l}_n & \\ & & \tilde{k}_n & & 0 \end{pmatrix} \in \mathbb{C}^{(n+1) \times (n+1)}, \tilde{k} \in \{\pm\sqrt{\sigma}\}^n, \tilde{l} \in \{\pm\sqrt{\sigma}\}^n.$$

This is no longer true for periodic operators. For tridiagonal operators on $\ell^2(\mathbb{Z})$ we can only shift phases to the other side. This remaining freedom, however, can be used to prove Theorem 1 for arbitrary alphabets as long as all elements share the same absolute value. This only needs a small modification of Lemma 3 and a refinement of [4, Theorem 4.1]. The details are left to the reader.

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