

Near-Earth solar wind forecasting using corotation from L5: the error introduced by heliographic latitude offset

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1 **Near-Earth solar wind forecasting using corotation**
2 **from L5: The error introduced by heliographic latitude**
3 **offset**

4 **M.J. Owens^{1*}, P. Riley², M. Lang³, M. Lockwood¹**

5 ¹Department of Meteorology, University of Reading, Berkshire, UK, RG6 6BB

6 ²Predictive Science Inc., 9990 Mesa Rim Rd, Suite 170, San Diego, CA 92121, USA

7 ³Le Laboratoire des Sciences du Climat et de l'Environnement, CEA-CNRS-UVSQ, 91191 Gif Sur Yvette,
8 France

9 **Key Points:**

- 10 • Heliographic latitudinal offset can have a significant effect on corotation solar wind
11 forecasting from L5
- 12 • The error introduced by latitudinal offset is maximised at solar minimum and dur-
13 ing the solstices
- 14 • A combination of L5 measurements and global MHD modelling may yield the most
15 accurate solar wind forecasts

*<https://orcid.org/0000-0003-2061-2453>

Corresponding author: M.J. Owens, m.j.owenst@reading.ac.uk

16 **Abstract**

17 Routine in-situ solar wind observations from L5, located 60° behind Earth in its orbit,
 18 would provide a valuable input to space-weather forecasting. One way to utilise such ob-
 19 servations is to assume that the solar wind is in perfect steady state over the 4.5 days
 20 it takes the Sun to rotate 60° and thus near-Earth solar wind in 4.5-days time would be
 21 identical to that at L5 today. This corotation approximation is most valid at solar min-
 22 imum when the solar wind is slowly evolving. Using STEREO data, it has been possi-
 23 ble to test L5-corotation forecasting for a few months at solar minimum, but the var-
 24 ious contributions to forecast error cannot be disentangled. This study uses 40+ years
 25 of magnetogram-constrained solar wind simulations to isolate the effect of latitudinal off-
 26 set between L5 and Earth due to the inclination of the ecliptic plane to the solar rota-
 27 tional equator. Latitudinal offset error is found to be largest at solar minimum, due to
 28 the latitudinal ordering of solar wind structure. It is also a strong function of time of
 29 year; maximum at the solstices and very low at equinoxes. At solstice, the latitudinal
 30 offset alone means L5-corotation forecasting is expected to be less accurate than numer-
 31 ical solar wind models, even before accounting for time-dependent solar wind structures.
 32 Thus, a combination of L5-corotation and numerical solar wind modelling may provide
 33 the best forecast. These results also highlight that three-dimensional solar wind struc-
 34 ture must be accounted for when performing solar wind data assimilation.

35 **1 Introduction**

36 Space weather can disrupt power grids, communications and satellite operations,
 37 and poses a threat to health of humans in space and on high altitude aircraft (Cannon
 38 et al., 2013). Long lead-time (> 1 day) space-weather forecasting requires accurate pre-
 39 diction of near-Earth solar wind conditions. For this purpose, UK and US forecast cen-
 40 tres primarily use numerical magnetohydrodynamic (MHD) coronal and solar wind mod-
 41 els constrained by observations of the photospheric magnetic field (Riley et al., 2001; Odstr-
 42 cil, 2003; Tóth et al., 2005). Transient structures resulting from coronal mass ejections
 43 (CMEs), can be inserted into numerical solar wind models (Odstrcil et al., 2004). Nev-
 44 ertheless, simple empirical solar wind forecasts can serve as useful independent forecasts,
 45 as well as providing a contingency if, e.g., magnetogram observations, are not available
 46 (Owens, Riley, & Horbury, 2017).

47 One such empirical forecast method is solar wind recurrence (sometimes referred
 48 to as 27-day persistence), which assumes corotation of steady-state or quasi-steady-state
 49 solar wind structures. It predicts that the near-Earth solar wind in one solar rotation's
 50 time (approximately 27.27 days relative to Earth's motion) will be the same as that to-
 51 day. This works well for heliospheric magnetic field (HMF) polarity and solar wind speed
 52 (V_R) at solar minimum, often outperforming the MHD models (Owens et al., 2013) for
 53 the ambient solar wind. Except for rare cases when long-lived active regions introduce
 54 27-day periodicity by producing multiple CMEs over multiple rotations, recurrence is
 55 not capable of forecasting transient CMEs. Towards solar maximum, however, the corona
 56 becomes increasingly dynamic and transient CMEs make up an increasing proportion
 57 of the solar wind (Cane & Richardson, 2003; Riley, 2007). Thus the steady-state assump-
 58 tion over 27 days becomes increasingly invalid and recurrence forecasting performs poorly
 59 (Owens et al., 2013). Information about the time evolution of the corona can be incor-
 60 porated to improve recurrence forecasts (Temmer et al., 2018).

61 A proposed operational space-weather mission (Hapgood, 2017) at the Lagrange
 62 L5 point, 60° behind Earth in its orbit, provides an opportunity to make a corotation
 63 forecast with a much shorter (4.5 days) assumption of steady-state conditions (Miyake
 64 et al., 2005). Such advanced knowledge of steady-state solar wind structures is expected
 65 to have wide-ranging space-weather applications (McGranaghan et al., 2014).

The STEREO mission (Kaiser, 2005), with two spacecraft in Earth-like orbits but drifting ahead and behind Earth at a rate of 22.5° per year, provided a unique opportunity to test a corotation forecast from L5 (Simunac et al., 2009; Turner & Li, 2011); the STEREO spacecraft were separated from each other by 60° longitude in early 2008, and from Earth by 60° near the end of 2009. During these few months of data during a particularly deep solar minimum, L5 corotation was shown to be superior to a 27-day recurrence forecast from near-Earth data (Kohutova et al., 2016). The improvement, however, was fairly modest, with only $\sim 20\%$ skill gain in solar wind speed forecast relative to 27-day recurrence (Kohutova et al., 2016; Thomas et al., 2018). As the STEREO spacecraft approached the far side of the Sun (from Earth's viewpoint), they again achieved a brief 60° separation in December 2013. During this time there was considerable difference in the solar wind structures (Jian et al., 2019).

Limitations on L5-corotation forecasting likely include: (1) transient CMEs which only encounter either L5 or Earth; (2) evolution of the large-scale solar wind structures over the 4.5-day corotation period; (3) small-scale stochastic processes such as solar wind turbulence (Bruno & Carbone, 2005); and (4) the heliographic latitudinal offset of L5 relative to Earth resulting from the inclination of the ecliptic plane to the solar rotation axis. From the few months of L5-like STEREO data it is not possible to disentangle these effects. The latitudinal offset was small amplitude, varying between 0 and 7° over the year. Global solar wind simulations, however, suggest that even small latitudinal variations can have a considerable effect on the solar wind structures encountered (Riley et al., 2010).

In this study, we use coupled coronal and heliospheric simulations, constrained by photospheric magnetic field observations, to produce completely steady-state reconstructions of the solar wind. By sampling L5 and Earth-like trajectories through the 1-AU solar wind structure, we isolate and quantify the effect of the L5-Earth heliolatitude offset for a range of solar activity levels.

2 Data

Solar wind structure is determined using the Magnetohydrodynamics Algorithm outside a Sphere (MAS) global coronal and heliosphere model (Linker et al., 1999; Riley et al., 2012). MAS is constrained by photospheric magnetic field observations, which are computed outward to 30 solar radii, while self-consistently solving the plasma and magnetic field parameters on a non-uniform grid in polar coordinates, using the MHD equations and the vector potential \mathbf{A} (where the magnetic field, \mathbf{B} , is given by $\nabla \times \mathbf{A}$, such that $\nabla \cdot \nabla \times \mathbf{A} = 0$ which ensures current continuity, $\nabla \cdot \mathbf{J} = 0$, is conserved to within the models numerical accuracy). The heliospheric version of MAS then propagates solar wind conditions out to 1 AU. We use MAS solutions based on Carrington maps of the photospheric magnetic field and thus assume that the solar wind is completely steady state over a Carrington rotation (CR). Magnetograms from a range of observatories are used to minimise data gaps and provide the longest possible time sequence. See Owens, Lockwood, and Riley (2017) for details of the dataset. The heliospheric MAS used in this study is based on the polytropic approximation (Linker et al., 1999) and uses a grid resolution of 121, 128 and 140 cells in the meridional, azimuthal and radial directions, respectively. As steady state solar wind is assumed and the time series at 1 AU is produced by rotation, the effective time resolution is set by the azimuthal resolution of MAS. In this case, the effective time resolution is approximately 5 hours.

The use of the MAS solar wind solutions in this study does not require the model to accurately forecast or reconstruct specific solar wind intervals. For the purposes of this study, we are only reliant on MAS reproducing the solar wind speed structures and the position of the heliospheric current sheet (HCS) in a statistical sense. Specifically, we require accurate representation of the latitudinal width of the slow wind band, and

117 the shape and inclination to the rotational axis of the slow wind band and HCS. These
 118 features control the proportions and durations of fast/slow wind and inward/outward
 119 magnetic sectors in the ecliptic plane and thus whether spacecraft at Earth and L5 are
 120 likely to observe the same solar wind features. We note that even a systematic bias in
 121 the reconstructed solar wind speed would not necessarily be an issue for the results pre-
 122 sented here

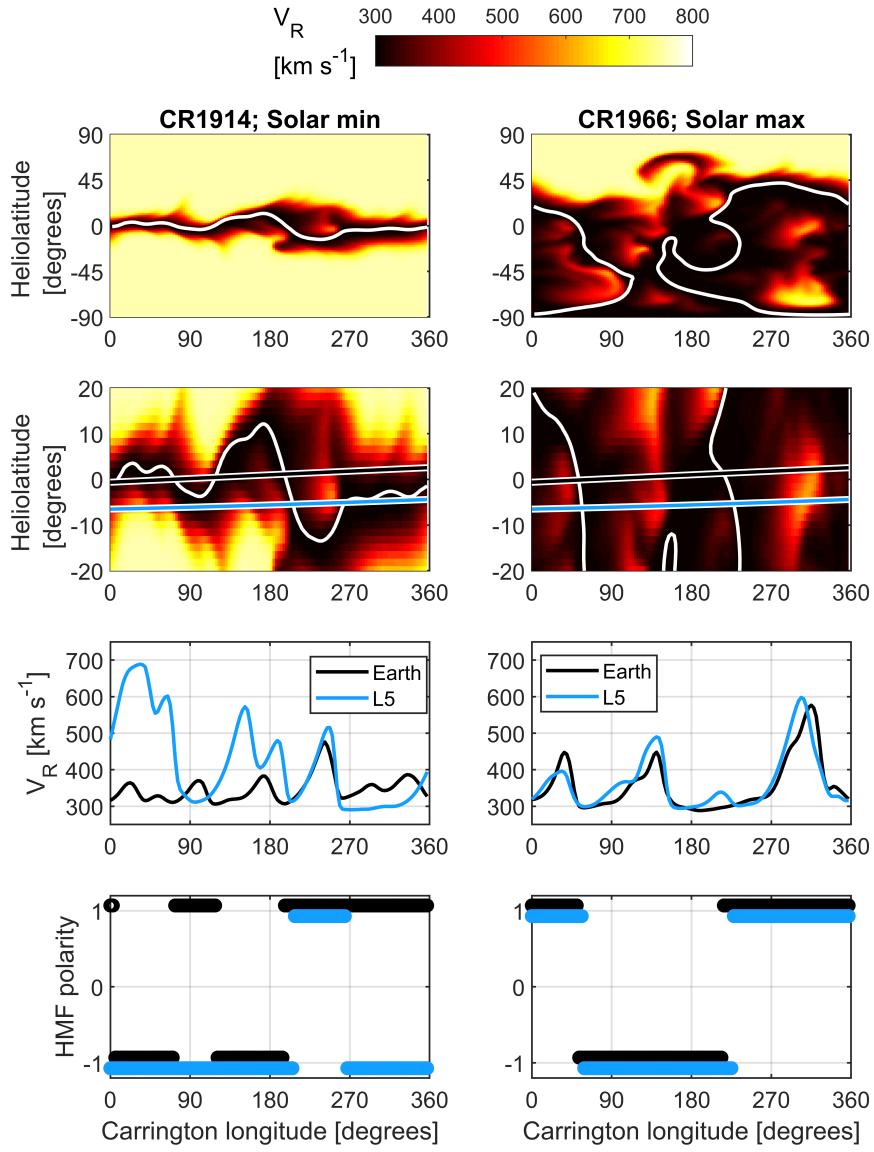
123 The width of the slow solar wind band predicted from the MAS model has been
 124 shown to provide good agreement with the three Ulysses fast latitude scans under dif-
 125 fering solar activity levels (Jian et al., 2011; Owens, Lockwood, & Riley, 2017; Jian et
 126 al., 2016). The latitudinal extent of slow wind predicted by MAS and observed using in-
 127 terplanetary scintillation is also in reasonable agreement (e.g., qualitatively comparing
 128 Figure 6 of Manoharan (2012) and Figure 4a of Owens, Lockwood, and Riley (2017)).
 129 The proportions of fast and slow wind seen in the ecliptic are generally well reproduced
 130 by magnetogram extrapolations when the estimated speed is a function of both the ex-
 131 pansion factor of a magnetic flux tube between the photosphere and the source surface,
 132 and the distance of the flux tube from the coronal hole boundary, as is the case with MAS
 133 (McGregor et al., 2011; Riley et al., 2015). Thus the relative occurrence of fast/slow streams
 134 near the ecliptic plane is reasonably well reproduced in a point-by-point manner (Jian
 135 et al., 2015; Owens et al., 2008; Riley et al., 2010). Taking a feature-based approach by
 136 identifying high-speed enhancements in model and observations and allowing for tim-
 137 ing errors which are not critical to the present study (Owens et al., 2005), MAS produces
 138 good agreement with observations, once transient solar wind structures are discounted
 139 (Owens et al., 2008; Jian et al., 2015, 2016).

140 Ulysses observations of the latitudinal extent of the HCS are well reproduced by
 141 magnetogram-constrained models such as MAS (e.g., see Figure 13 in Owens and Forsyth
 142 (2013)). This is also true in the ecliptic plane, with a strong correspondence between magnetogram-
 143 based models and in-situ magnetic sector structure, even at solar maximum (e.g., see Fig-
 144 ure 5 in Owens and Forsyth (2013)).

145 Given models such as MAS generally produce “smoother” solar wind speed and he-
 146 liospheric magnetic field structures than those observed, the differences between L5 and
 147 Earth are likely to be underestimated. Therefore the values estimated in this study should
 148 be treated as a lower limit.

149 3 Results

150 Examples of 1-AU solar wind structure are shown in Figure 1. For the solar min-
 151 imum example (CR1914), slow wind is confined to the equatorial region, while fast wind
 152 fills the heliosphere for latitudes more than 25° from the equator. The latitudinal gra-
 153 dients in solar wind speed are large close to the equator (and ecliptic plane). Similarly,
 154 the HCS lies very close the equator. While CR 1914 spanned mid-September to mid-October
 155 1996, we consider how the corotation forecast would have fared if such a solar wind con-
 156 figuration was encountered in June. Around the summer solstice, Earth lies close to the
 157 helioequator and consequently remains primarily within the slow solar wind, except for
 158 a moderate increase in solar wind speed around 230° Carrington longitude. Earth crosses
 159 the HCS three and thus sees four different magnetic sectors. That the polarity of the field
 160 at Earth on either side of the HCS reflects that of solar polar field was first noted by Rosenberg
 161 and Coleman (1969) and this “Rosenberg-Coleman effect” means that L5 and Earth will
 162 see opposite polarity HMF when their difference in heliographic latitudes places them
 163 on opposite sides of the HCS. The time Earth spends above/below the HCS in Figure
 164 1 is approximately equal, resulting in roughly equal proportions of inward and outward
 165 polarity HMF. Conversely, L5 is well below the equator during June, approximately -
 166 6.5 to -4.5° heliolatitude. As a result, L5 encounters significantly more fast wind than
 167 Earth, particularly around Carrington longitudes of 0 to 80°, and L5 remains predom-
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150 **Figure 1.** Examples of 1-AU solar wind structure from solar minimum (left)
151 and solar maximum (right). Top: The global solar wind structure at 1 AU, shown as a heliolatitude-Carrington
152 longitude map. The colour map extends from 300 (black) to 800 km s^{-1} (white). The white line
153 shows the heliospheric current sheet. Second row: Same as top, but with latitude restricted to
154 20 degrees about the equator. Black and blue lines show paths of Earth and L5, respectively, if
155 these solar wind structures had been encountered in June. Third row: Solar wind speed at Earth
156 (black) and L5 (blue). Bottom: Heliospheric magnetic field polarity in the same format.

175 inantly below the HCS and thus sees primarily inward polarity HMF. Thus if near-Earth
 176 conditions were predicted using L5-corotation for this solar wind configuration during
 177 June, the Mean Absolute Error (MAE) in V_R would be 106 km s^{-1} and the HMF po-
 178 larity would be incorrect 44% of the time.

179 At solar maximum (CR1966, approximately spanning August 2000), the picture
 180 is very different. Slow solar wind dominates, with fast wind confined to the north pole
 181 (at latitudes above 45°). Within 20° of the equator, the latitudinal speed gradients are
 182 greatly reduced compared to the solar minimum case, and the HCS is essentially verti-
 183 cal. As a result, L5 and Earth see almost identical solar wind conditions, and a corota-
 184 tion forecast would give a V_R MAE of only 26 km s^{-1} and only 5% of the HMF polar-
 185 ities would be incorrect. Of course, these errors are purely for the latitudinal offset be-
 186 tween L5 and Earth and at solar maximum the corona is far more dynamic. Thus the
 187 steady-state assumption would become the primary source of error in a L5-corotation
 188 forecast at this time.

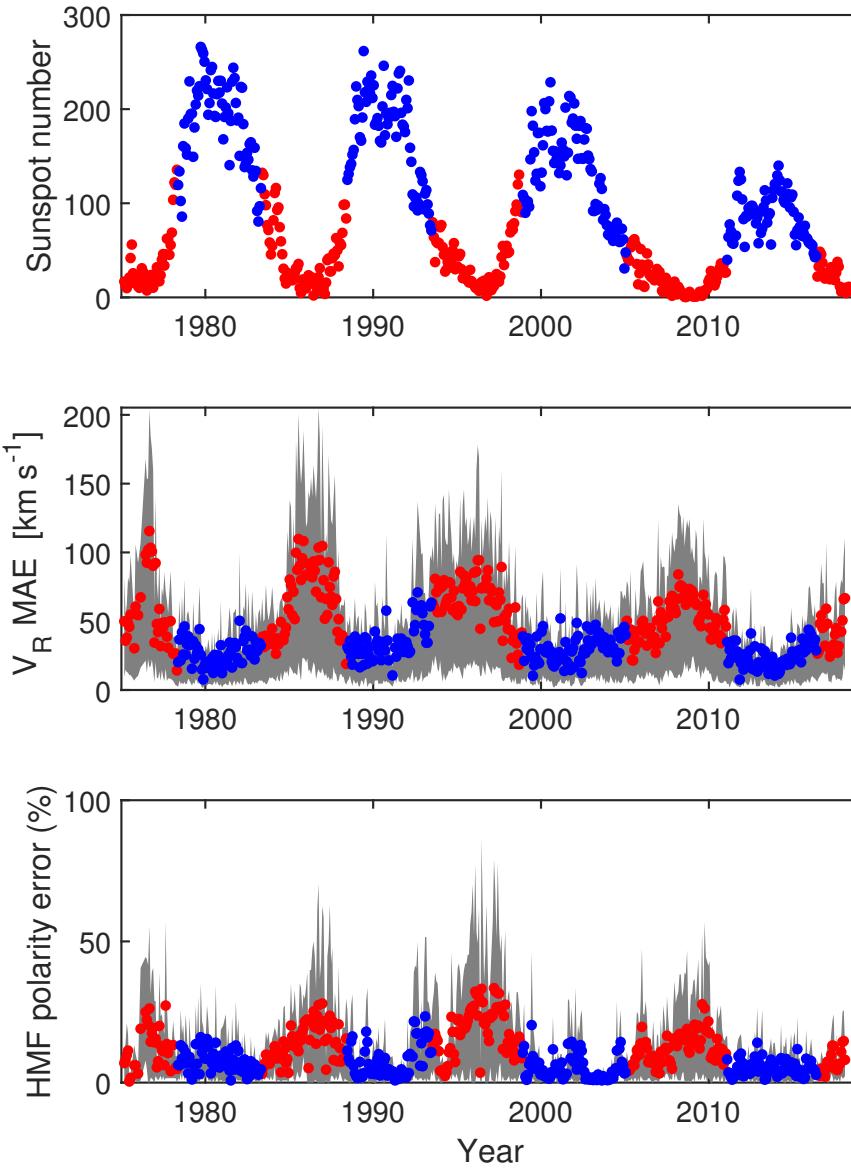
197 We extend this analysis to a statistical study using all Carrington rotations from
 198 early 1975 to mid 2018 (CR 1625 to 2203). This spans four sunspot cycles, as shown by
 199 the upper panel of Figure 2. For each CR, we consider how the L5-corotation forecast
 200 would vary with time of year. The annual variation in heliolatitude of L5 and Earth is
 201 shown in the top panel of Figure 3. It can be seen that monthly sampling of the heli-
 202 olatitudes of Earth and L5 gives full coverage of the latitudinal offsets between the two
 203 positions. Thus for each Carrington rotation, we sample the L5 and Earth latitudes us-
 204 ing the latitudinal positions for each individual calendar month.

205 Figure 2 shows the time series of L5 corotation errors purely from latitudinal off-
 206 set. As expected from the two example CRs shown in Figure 1, on average the V_R MAE
 207 is maximised around solar minimum and minimised at solar maximum (very similar trends
 208 are found for other metrics, such as root-mean-square error). The percentage of incor-
 209 rect magnetic polarity intervals follows the same basic trend. For a given CR, there is
 210 a large spread in the V_R and magnetic polarity error depending on the month at which
 211 the prediction is made. e.g., During 1986, the mean V_R MAE is approximately 100 km
 212 s^{-1} , but the range spans 20 to 190 km s^{-1} . In 1996 the average percentage of incorrect
 213 HMF polarities is approximately 30%, but the range spans 0 to more than 75% (i.e., sig-
 214 nificantly worse than random chance). We also note long-term trends, with both peak
 215 and mean values of V_R MAE reduced in the most recent solar minimum (2008-2010) com-
 216 pared with the three previous minima.

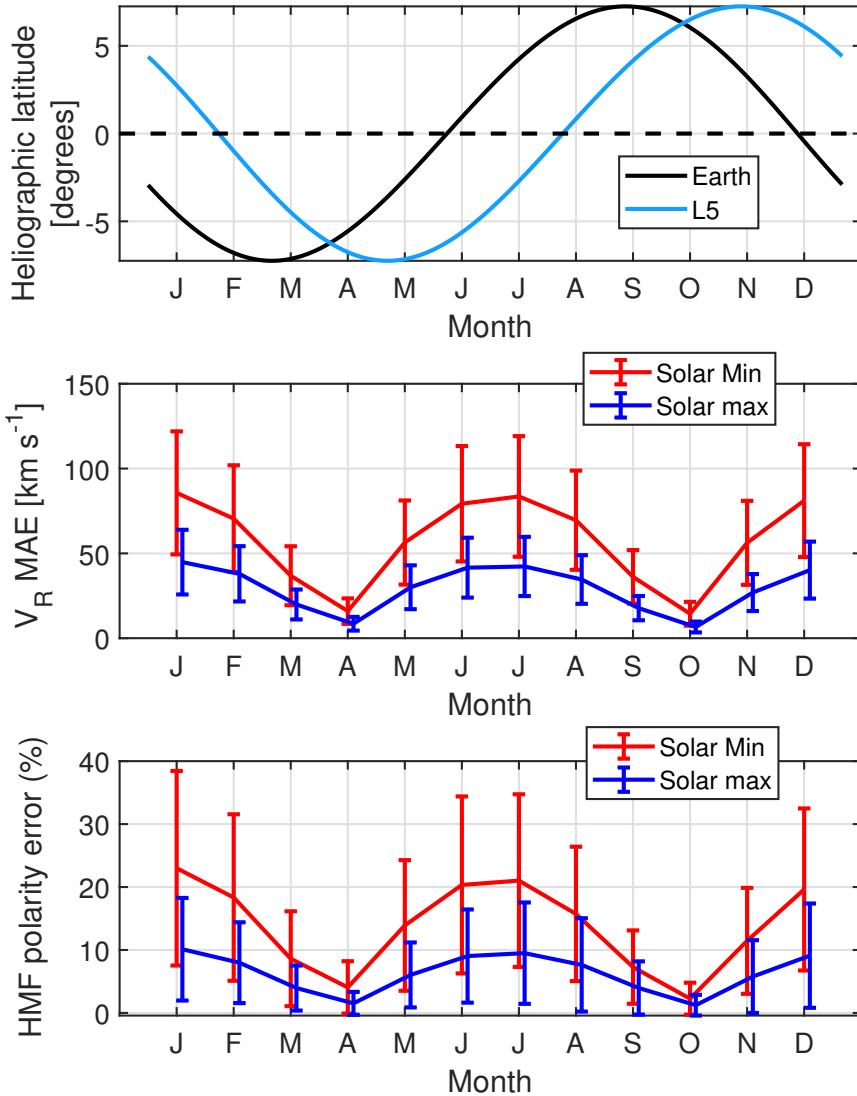
222 Figure 3 shows that the large spread in forecast errors for a given CR is a result
 223 of the large annual variation in heliolatitude offset between L5 and Earth. CRs are split
 224 into solar minimum and maximum periods on the basis of the solar cycle phase. This
 225 is deemed preferable to using a simple sunspot number threshold, as that would select
 226 different proportions of small/large sunspot cycles. We use solar cycle phase limits of 0.17
 227 and 0.67 of the way through the cycle from solar minimum (Owens et al., 2011), which
 228 both splits the dataset in half and selects similar sunspot number gradients in the rise
 229 and declining phase of each cycle (but allows for changing magnitude of cycles). At so-
 230 lar minimum during the solstices, the V_R MAE is $80 \pm 30 \text{ km s}^{-1}$. During the equinoxes,
 231 however, this drops to $20 \pm 11 \text{ km s}^{-1}$. For the HMF polarity, $23 \pm 16\%$ of intervals
 232 are incorrect at solstices, dropping to $5 \pm 4\%$ at equinoxes.

233 4 Conclusions

234 Near-Earth solar wind can be forecast from in-situ L5 observations by assuming
 235 the solar wind is steady-state over the 4.5 days it takes structures to rotate between the
 236 two positions (Kohutova et al., 2016; Thomas et al., 2018; Miyake et al., 2005). Such L5-
 237 corotation forecasting will undoubtedly be a useful additional tool in space-weather pre-



189 **Figure 2.** Time series of L5-corotation prediction error resulting purely from latitudinal
 190 offsets between L5 and Earth. All data are Carrington rotation averages with red and blue in-
 191 dicating solar minimum and maximum, respectively, based on the phase of the solar cycle. Top:
 192 Sunspot number. Middle: MAE (Mean Absolute Error) in near-Earth V_R based on corotation
 193 from L5. We compute MAE for Earth and L5 latitudes during the 12 calendar months. Coloured
 194 dots show the mean value for all months, and the shaded area spans the maximum and minimum
 195 values. Bottom: The percentage of incorrect magnetic polarity intervals, in the same format as
 196 the middle panel.



217 **Figure 3.** The annual variation in L5-corotation prediction error resulting purely from lat-
 218 ititudinal offsets between L5 and Earth. Top: The heliographic latitude of L5 (blue) and Earth
 219 (black) over the year. Middle: V_R MAE as a function of time of year for solar minimum (red)
 220 and solar maximum (blue). The line shows the mean of all CRs, while the error bar is one stan-
 221 dard deviation. Bottom: The percentage of incorrect HMF polarities, in the same format.

238 diction. However, there are inherent limitations to the skill of such forecasts, not only
 239 from the steady-state assumption and existence of transient solar wind structures, but
 240 also from the heliographic latitudinal difference resulting from the inclination of the ecliptic
 241 plane to the solar rotational equator. Isolating and quantifying the contributions of
 242 these effects is difficult with the limited solar wind data from L5 (or L5-like longitudi-
 243 nal spacecraft separations) presently available.

244 In this study we have used magnetogram-constrained simulations of the solar wind
 245 over the last 40+ years to isolate and quantify the latitudinal offset effect. The latitu-
 246 dinal separation of L5 and Earth has the largest influence on corotation forecasting at
 247 solar minimum, when the solar wind is latitudinally structured, and during the summer
 248 and winter solstices (December/January and June/July). At solar maximum, when so-
 249 lar wind features are more longitudinally structured, and during the equinoxes, the ef-
 250 fect of latitudinal offset is reduced. However, at solar maximum the steady-state approx-
 251 imation breaks down and reduces the usefulness of L5-corotation forecasting.

252 Comparing the last four solar minima, the STEREO/ACE L5-analogous periods
 253 from 2008 and late 2009 may not be representative of the skill of corotation forecasting
 254 in general. While the average V_R MAE owing to latitudinal offset in the deep 2008-2010
 255 minimum was reduced compared with previous minima (by approximately 25%). This
 256 is likely to be the result of the 2008-2010 minimum producing a broader slow wind band
 257 than previous minima, meaning the latitudinal gradient in solar wind speed near the he-
 258 lioequator is reduced (Owens et al., 2014).

259 At solar minimum, the expected V_R MAE for L5-corotation forecasts ranges from
 260 approximately 20 km s⁻¹ at the equinoxes to 80 km s⁻¹ at the solstices. This is purely
 261 from the latitudinal offset and, in practice, will be the lower limit with additional error
 262 introduced by time-dependent solar wind structures and transient CMEs which are only
 263 seen (or are different) at either L5 or Earth. As a comparison, steady-state magnetogram-
 264 constrained numerical MHD solar wind models have been shown to produce MAE er-
 265 rrors in V_R of around 70-80 km s⁻¹ (Owens et al., 2008). This estimate, however, is from
 266 comparison with real solar wind observations and thus does include time evolution, tur-
 267 bulence and CMEs. For the heliospheric magnetic field polarity, L5-corotation is expected
 268 to produce incorrect polarities around 25% of the time at the solstices, but correctly pre-
 269 dict the polarity around 95% of the time at the equinoxes. Thus, a hybrid scheme may
 270 be the best forecasting approach, wherein L5-corotation are preferentially weighted at
 271 equinox and numerical solar wind models are preferentially weighted at solstice. Such
 272 a hybrid scheme could be based upon data assimilation (Lang et al., 2017) that combines
 273 the observations, at L5, with a numerical solar wind forecasting model in order to op-
 274 timally estimate the solar wind at Earth.

275 Note that the HMF polarity error in co-rotation predictions from L5 data will have
 276 complex effects on predictions of geoeffectiveness because of the influence of the north-
 277 south component of the field in the Geocentric Solar Magnetospheric (GSM) frame of
 278 reference ($[B_z]_{GSM}$), which predicts the magnetic shear across the dayside boundary of
 279 the magnetosphere and so the degree to which the transfer into the terrestrial magne-
 280 tosphere of solar wind mass, momentum and energy can occur. This is a complex mix-
 281 ture of the “Rosenberg-Coleman” (RC) and “Russell-McPherron” (RM) effects (Rosenberg
 282 & Coleman, 1969; Russell & McPherron, 1973). The RC effect means that the latitude
 283 difference between L5 and Earth can cause a difference in the polarities of $[B_x]_{GSE}$ (to-
 284 ward, radial in the Geocentric Solar Ecliptic frame) component at the two sites and this
 285 would generally give a difference in the $[B_y]_{GSE}$ component because the dominant gar-
 286 denhose HMF orientation of the Parker spiral means that $[B_y]_{GSE}/[B_x]_{GSE} < 0$. The
 287 RM effect is where the $[B_y]_{GSE}$ component of the HMF yields a ($[B_z]_{GSM}$) component
 288 because of the angle between the GSE and GSM frames. This angle has an annual and
 289 a diurnal variation due to, respectively, the angle between the Earth’s rotational axis and
 290 the X direction of the GSE frame and due to the offset between Earth’s rotational and

291 magnetic axes (Lockwood et al., 2016). Specifically, southward HMF in the GSM frame
 292 is generated around the March equinox (peaking at 20 UT) when $[B_y]_{GSE} < 0$ (i.e. $[B_x]_{GSE} >$
 293 0) and around the September equinox (peaking at 10 UT) when $[B_y]_{GSE} > 0$ (i.e. $[B_x]_{GSE} <$
 294 0). Although this RC-RM effect introduces quite a lot of diversity into the difference in
 295 geoeffectiveness between the solar wind/HMF seen at Earth and L5 when the polarities
 296 of the radial HMF at Earth and L5 are different, that diversity is systematic with time
 297 of year and UT and so is relatively easily predicted.

298 In addition to corotation forecasting, in-situ solar wind observations from L5 would
 299 enable improvement of numerical solar wind forecasting via data assimilation (DA). DA
 300 is the merging of model and observational data to ensure an optimal estimate for real-
 301 ity. Forecast skill has been improved via DA using L5-like observations with a two-dimensional
 302 solar wind model (Lang & Owens, 2019). In standard L5-corotation forecasts, the ob-
 303 servation is assumed to be “truth”, in that it contains no errors. Figure 3 shows one ex-
 304 ample of why this is not always the case; L5 can be sampling solar wind from a latitude
 305 that is not representative of the solar wind at Earth. DA allows estimates of the solar
 306 wind to be modified to account for observation errors present as a result of incorrect mod-
 307 ellings/assumptions (e.g. assuming that the observations have no latitudinal offset). The
 308 results presented here also suggest that L5 data assimilation would benefit greatly from
 309 use with fully three-dimensional solar wind models.

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 317 tions used in this study can be visualised at http://www.predsci.com/mhdweb/plot_2d.php and downloaded at http://www.predsci.com/mhdweb/data_access.php. Sunspot
 318 data are available from <http://www.sidc.be/silso/datafiles>.

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