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1 **To be resubmitted to Building and Environment 2019**

2 **The influence of aspect ratios and wall heating conditions on flow**
3 **and passive pollutant exposure in 2D typical street canyons**

4

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20

21 Abstract

22 Deep street canyons and unfavourable meteorological conditions usually induce high
 23 pollutant exposure. Validated by experimental data, this paper employs computational fluid
 24 dynamic simulations with RNG $k-\varepsilon$ model to investigate the flow, and passive pollutant
 25 dispersion within scale-model two-dimensional street canyons($H=3m$). As a novelty, this
 26 paper quantifies the impacts of various wall heating scenarios(bottom, leeward/windward
 27 wall and all-wall heating), ambient velocity($U_{ref}=0.5-2m/s$, Froude numbers $Fr=0.25-4.08$,
 28 Reynolds numbers $Re=95602-382409$) and aspect ratios(building height/street width, $AR=0.5$,
 29 $0.67, 1, 2, 3$) on personal intake fraction for entire streets($\langle P_{IF} \rangle$). The governing equations
 30 are implicitly discretized by a finite volume method (FVM) and the second-order upwind
 31 scheme with Boussinesq model for quantifying buoyancy effects. The SIMPLE scheme is
 32 adopted for the pressure and velocity coupling.

33 In most isothermal cases, one-main-vortex structure exists as $AR=0.5-3$ ($\langle P_{IF} \rangle=0.43-$
 34 $3.96ppm$ and $1.66-27.51ppm$ with $U_{ref}=2$ and $0.5m/s$). For non-isothermal cases with
 35 $Fr=4.08$ ($U_{ref}=2m/s$), wind-driven force dominates urban airflow as $AR=0.5-1$ and four
 36 heating conditions attain similar $\langle P_{IF} \rangle$ ($0.39-0.43ppm, 0.57-0.60ppm, 0.91-0.98ppm$). As
 37 $AR=2$, windward and all-wall heating get two-vortex structures with greater $\langle P_{IF} \rangle$ ($3.18-$
 38 $3.33ppm$) than others($\langle P_{IF} \rangle=2.13-2.21ppm$). As $AR=3$, leeward-wall heating slightly
 39 reduces $\langle P_{IF} \rangle$ ($\sim 3.72-3.96ppm$), but the other three produce two-vortex structures with
 40 greater $\langle P_{IF} \rangle$ ($6.13-10.32ppm$). As $Fr=0.25$ ($U_{ref}=0.5m/s$), leeward-wall heating always
 41 attains smaller $\langle P_{IF} \rangle$ ($1.20-7.10ppm$) than isothermal cases($1.66-27.51ppm$) as $AR=0.5-3$,
 42 however the influence of the other three is complicated which sometimes raises or reduces
 43 $\langle P_{IF} \rangle$. Overall, smaller background wind speed ($Fr=0.25$) with two-vortex structures
 44 attains much larger $\langle P_{IF} \rangle$. Special attention is required at night(all-wall heating),

45 noon(bottom-heating) and cloudy period(no-wall heating) as $AR=2-3$, while it is during
46 windward-wall heating and cloudy period for $AR=0.5-1$.

47

48 **Keywords:** Street canyon, Aspect ratio (AR), Wall heating, Street intake fraction $\langle P_{IF} \rangle$,

49 Froude number, Computational fluid dynamic (CFD) simulations

50

51 1. Introduction

52 With the increase in number of vehicles on the road due to global urbanization, traffic
53 emissions have become one of the major pollutant sources in cities [1, 2]. Critical pollutants
54 emitted from these sources include oxides of nitrogen (NO_x), carbon monoxide (CO),
55 ultrafine particles and fine particulate matter ($PM_{2.5}$) with an aerodynamic diameter of less
56 than $2.5 \mu m$ and volatile organic compounds (VOCs). Heavy traffic flow, deep street canyons
57 and unfavourable meteorological conditions are the main factors that result in poor
58 ventilation capacity, a high pollutant exposure for urban residents and the related adverse
59 impacts on human health [3, 4]. Traffic-related pollutant exposure is determined by three
60 factors: the emission rate of pollutants, as determined by traffic density and types; the
61 pollutant dilution capacity, which is correlated with street layouts and meteorological
62 conditions, and the distance between people and pollutant sources. Vehicular pollutant
63 exposure for residents living in near-road buildings merits special attention because their
64 proximity to emission sources puts them at higher health risk than those living in other urban
65 micro-environments. In addition to reducing vehicular pollutant emissions, improving urban
66 ventilation capacity through sustainable street design is another effective technique to reduce
67 such traffic-related pollutant exposure in cities [5-9].

68 In the past three decades, the relation of street layout and atmospheric conditions to
69 turbulent flow and pollutant dispersion has been widely investigated and modeled using field
70 and wind tunnel experiments and computational fluid dynamic (CFD) simulations [10-20].
71 Street aspect ratios (building height/street width, AR or H/W) [8, 10-12, 21-23], building
72 packing densities and urban porosity (e.g., [24-27]) are the most significant factors
73 influencing urban airflow and traffic-related pollutant dispersion. Other reportedly key
74 parameters are ambient wind directions [33-34], uneven street layouts and building height
75 variability [30-32], street vegetation [33] and lift-up building designs [e.g., 34-35]. The
76 thermal dynamics of street canyons related to solar shading and the thermal storage of
77 buildings also affect the flow of pollutants through the urban environment. Field
78 measurement have shown that air-wall temperature differences can reach 12-14°C or more
79 [50-53]. Models of three-dimensional (3D) urban-like environments [37-40] and two-
80 dimensional (2D) street canyons [40-47] have shown that if the Richardson number (Ri) is
81 large or the Froude number (Fr) is small, thermal stratification and buoyancy force, or
82 thermal forcing, can influence or dominate the flow regime and pollutant dispersion . In the
83 2D street canyon models in most of the aforementioned studies, various uniformly heated
84 walls are considered with arbitrary air-wall temperature differences corresponding to solar
85 angle or building heat release within a day [40-47], i.e., no-wall heating (periods of strong
86 wind or cloudy days with small temperature differences), leeward-wall heating, windward-
87 wall heating, ground or bottom heating (at noon on a sunny day), and all-wall heating (i.e., at
88 nighttime with the urban heat island heating all wall surfaces).

89 In 2D street canyon models, four isothermal flow regimes have been reported [8-12, 20-
90 23, 41-54], those being the isolated roughness flow regime (IRF, AR or $H/W < 0.3$), the wake
91 interference flow regime (WIF, $0.3 < AR < 0.67$), the skimming flow regime with single main
92 vortex (SF, $0.67 < AR < 1.5$), and the multi-vortex flow regime in deep street canyons. The

93 literature is generally consistent with regard to the first three flow regimes but differs on the
94 fourth multi-vortex flow in deep street canyons, in which the flow and vehicular pollutant
95 dispersion capacity are usually weak. For instance, Xie et al. [52] and Li et al. [23] reported
96 two contra-rotative vortices where $AR=2$, and three to five vertically aligned vortices where
97 $AR=3-5$ for wind-tunnel-scale 2D street canyons with a building height of $H=0.6$ m and the
98 reference Reynolds number (Re) of 12000. Other research has contradicted these findings.
99 Zhang et al. [53] found a single-main-vortex structure in a full-scale street canyon where
100 $AR=2.7$ and $Re=5 \times 10^6$. Later validated by wind tunnel and scale-model outdoor
101 experimental data, He et al. [54] numerically confirmed a single-main-vortex structure as
102 $AR=1-4$ and two main vortices as $AR=5-6$ for full-scale 2D street canyons ($W=24$ m,
103 $Re \sim 10^6-10^7$). That study [54] reported that the reason for this difference was that Re must be
104 much greater than 11000 to ensure Reynolds number independence in urban airflow [55], and
105 that full-scale models usually satisfy this requirement [53-54] but wind-tunnel-scale models
106 sometimes cannot (e.g. [22-23, 49-52]). Recently, Chew et al. [56] further confirmed this issue
107 by conducting water channel experiments with $Re \sim 10^4-10^5$ at three aspect ratios ($AR=1, 1.5$
108 and 2) and pointed out that the widely adopted $Re=11,000$ is not applicable for the Re
109 independence of street canyons with an aspect ratio greater than 1.5.

110 In recent years, experimental and numerical studies of wind-tunnel-scale models
111 ($H \sim 0.1$ m, $Re \sim 10^4$) [48-52], scaled models ($H \sim 1$ m, $Re \sim 10^5$) [44-47] and full-scale models
112 ($H \sim 10$ m-100 m, $Re \sim 10^6 \sim 10^7$) [41-43] have been performed to investigate the relative flow and
113 temperature distribution and pollutant dispersion in 2D street canyons by coupling dynamic
114 and thermal effects. Chew et al. [57] reported differing findings between wind-tunnel-scale
115 experiments and full-scale field measurements with heated windward walls, even with similar
116 Fr or Ri numbers. Such contradictory buoyancy effects are present mainly because wind-
117 tunnel-scale experiments with heated windward walls do not satisfy the the requirement of

118 Reynolds number independence [56]. Wind-tunnel-scale results for canyon flows with
119 thermally induced buoyancy should not be assumed to represent full-scale street canyons,
120 unless the flow is verified to be independent of both Reynolds number and a similar Grashof
121 number (or Ri and Fr). In particular, for wind-tunnel-scale models, it is relatively difficult to
122 measure or simulate non-isothermal urban airflow with significant thermal effects because it
123 is a challenge to simultaneously attain a sufficiently large Reynolds number and relatively
124 small Fr (or large Ri) because this usually requires a large temperature difference ($\sim 100^\circ\text{C}$)
125 [48-50]. Thus, full-scale models ($H\sim 10\text{m}-100\text{m}$, $Re\sim 10^6\sim 10^7$) [38-43] and scaled models
126 ($H\sim 1\text{m}$, $Re\sim 10^5$) [44-47] are proposed to study urban airflow coupling dynamic and thermal
127 effects. A scale-model CFD simulation ($H\sim 1\text{m}$, $Re\sim 10^5$) was selected in this study,
128 considering that scaled models [44-47] make it easier to satisfy the Re independence
129 requirement and get similar Fr (or Ri) as full-scale models [38-43], and CFD simulations of
130 full-scale 2D or 3D streets with heated walls usually require enormous computational
131 resources because fine grids are required to solve the viscous sub-layer and heat transfer near
132 wall surfaces [38-40].

133 Considering the differing findings in the literature, further investigations are still
134 necessary to verify the non-isothermal flow mechanisms in high-rise deep street canyons
135 ($AR>1.5$) with a sufficiently large Reynolds number and various Froude (or Richardson)
136 numbers. In addition, most previous studies only investigated the flow and spatial distribution
137 of temperature and pollutant concentration. Few researchers have estimated the impact of
138 street layouts and wall heating conditions on personal exposure to air pollutants within micro-
139 scale street canyons. For instance, Memon et al. [44] studied the impact of street aspect ratios
140 ($AR=0.5-8$), four wall-heating conditions and ambient wind speeds ($0.5-4\text{ m s}^{-1}$) on air
141 temperature in 2D scale-model street canyons ($H\sim 0.5\text{m}-8\text{m}$) with $Re\sim 16000-270000$ and a
142 bulk Richardson number ($Ri\sim 0.01-17.1$). Tong and Leung [43] later modeled the reactive

143 pollutant dispersion within full-scale urban street canyons ($AR=0.5-8$, $H=20-80m$) with
144 various wall heating and ambient wind conditions. Yet few studies have considered the
145 impacts of these factors on the detailed flow structure and the related pollutant exposure on
146 street level.

147 Finally, vehicular intake fraction (IF) in urban areas was used to represent the fraction of
148 total pollutant emissions by vehicles that is inhaled by a population [58-60]. An IF of 1 ppm
149 (one per million or 10^{-6}) indicates 1 g of air pollutants is inhaled by an exposed population for
150 every one ton of pollutants emitted by the vehicles in that city and its street canyons. Personal
151 intake fraction (P_IF), which is independent of population size and density, has also been
152 adopted by the literature [8-9, 61-63] to compare the fraction of pollutants inhaled by each
153 person in a population on average to the total emitted vehicular pollutants in nearby streets or
154 neighborhoods. So far these investigations on pollutant exposure emphasize 2D or 3D
155 idealized urban models under neutral atmospheric conditions [8-9, 58-60, 61-63].

156 This paper couples the personal intake fraction (P_IF) with CFD simulations to
157 quantify the impacts of street aspect ratios ($AR=H/W=0.5-0.67$ (avenue canyon), 1 (regular
158 canyon), 2-3 (deep canyon) [10]) and four kinds of wall heating conditions (at leeward,
159 windward, ground and all walls) on the detailed flow structure, CO dispersion and personal
160 exposure in 2D scale-model street canyons. As a novelty, the interaction of wind-driven
161 airflow and buoyancy force with a sufficiently large Reynolds number and various Froude
162 numbers and the detailed flow structure and related street-scale CO exposure are emphasized,
163 as this interaction is still unclear and requires further investigation.

164 The structure of the remainder of this paper is as follows: Section 2 describes the
165 concept of the personal intake fraction. Section 3 depicts the cases investigated and the
166 numerical set-up. Results and discussions are given in Section 4. Conclusions are provided in

167 Section 5. The Appendix presents CFD model validation using wind tunnel data [22, 49] as
168 well as the scale-model outdoor field measurement in Zhang et al. [9].

169

170 **2. Population intake fraction (IF) and Personal intake fraction (P_IF)**

171 Intake fraction (IF) has been extensively applied to evaluate the population exposure to
172 vehicular emissions in streets or cities, some examples being the ~270 ppm value derived for
173 the high-rise compact city of Hong Kong [58], the street-scale vehicular IF of 371 ppm in a
174 street ($AR=H/W=1.5$) in central Athens in Greece [59] and the overall IF of 3000 ppm for a
175 typical street canyon in midtown Manhattan, New York [60]. For idealized 2D street canyon
176 models, Hang et al. [8] reported vehicular CO IF s of 230-913 ppm where $AR=1-0.5$. Later,
177 He et al. [54] further clarified that IF could reach $\sim 10^5$ ppm in extremely deep 2D street
178 canyons with two main vortexes ($H/W=5-6$). It is therefore apparent that vehicular intake
179 fraction for a population (IF) is independent of the pollutant emission rate and depends on
180 several factors, such as the street layout, meteorological conditions, distance to pollutant
181 sources and local population size and density.

182 The literature [8-9, 61-63] has also adopted personal intake fraction (P_IF) to quantify
183 the average pollutant exposure for each person in a population, which is independent of
184 population size and density and can emphasize the influence of urban morphology and
185 atmospheric conditions. Similarly, the spatial mean values of a building or entire street were
186 named as building intake fraction or street intake fraction, respectively ($\langle P_IF \rangle$) [61-63].
187 One study numerically estimated the $\langle P_IF \rangle$ in 2D street canyons as ~1-5 ppm when
188 $AR=0.5-1$ [8] and ~100-1000 ppm when $AR=5-6$ [9, 54]. Other studies [61-63] further
189 evaluated $\langle P_IF \rangle$ in 3D urban district models ($AR=0.5-1$, ~0.1 ppm) to be one-order smaller
190 than that in 2D street canyons with similar aspect ratios (~1 ppm).

191 The intake fraction (IF) for the emission of a specific pollutant is defined as:

$$192 \quad IF = \sum_i^N \sum_j^M P_i \times Br_{i,j} \times \Delta t_{i,j} \times Ce_j / q \quad (1)$$

193 where N is the number of population groups and M is the number of different
 194 microenvironments considered, P_i is the total number of people exposed in the i^{th} population
 195 group; Δt_{ij} (s) is the time spent in the microenvironment j for population group of i ; $Br_{i,j}$ is
 196 the average volumetric breathing rate for individuals in the i^{th} population group (m^3/s) in the
 197 microenvironment j ; Ce_j is the pollutant concentration attributed to urban traffic emissions in
 198 microenvironment j (kg/m^3); and q is the total vehicular emission rate over the period (kg).

199 As shown in Table 1 [58], three age groups were defined: children, adults, and the
 200 elderly, which means that $N=3$ in this study. As depicted in Fig. 1 [64-65], the time-activity
 201 patterns were divided into four micro-environmental types ($M=4$) for the three age groups,
 202 including indoors at home ($j=1$), other indoor locations ($j=2$), in or near vehicles ($j=3$), and
 203 other outdoor locations away from vehicles ($j=4$). It was assumed that the near-road buildings
 204 were naturally ventilated residential buildings, and two microenvironments $j=1$ (indoor at
 205 home) and $j=3$ (in or near a vehicle, i.e., pedestrian level) were considered. The values for the
 206 breathing rates from previous studies [64-65] were adopted for the current study. Furthermore,
 207 As the indoor/outdoor (I/O) ratio in naturally ventilated buildings is close to one [3-4, 66-67],
 208 it is reasonable to use the pollutant concentration, originating from vehicle emissions, at
 209 building wall surfaces as the indoor concentration in naturally ventilated buildings.

210 The overall IF value increases linearly as the population density rises. To normalize this
 211 value, the personal intake fraction (P_IF) was applied for the average intake fraction of each
 212 person in a specific population. The definition of P_IF is expressed in Eq. (2) [8-9].

$$P_IF = IF / \sum_j^M P_i \quad (2)$$

where IF is the total population intake fraction, and P_i is the total number of people exposed in the i^{th} population group.

The spatial mean P_IF for an entire street is defined as the street intake fraction $\langle P_IF \rangle$ to evaluate the average P_IF for a population on the entire street.

218

219 3. Methodology

220 Ansys Fluent [68] with the Renormalization Group (RNG) $k-\varepsilon$ model [69] was adopted
221 to perform CFD simulations and numerically investigate the effects of typical aspect
222 ratios ($AR=0.5-3$) and thermal buoyancy force induced by various types of wall heating on
223 turbulent structures, passive pollutant dispersion and its exposure in two-dimensional (2D)
224 street canyons.

225

226 3.1 Consideration of 2D street geometry and selection of turbulence model

227 This study first considers idealized 2D street canyons with a simplified urban geometry
228 where the street is infinitely long (e.g., street length $L > 8H$) and surrounded by buildings, with
229 a wind approaching perpendicular to the street axis [10-12, 20-22]. Modelling urban street
230 canyons in 2D may simplify the 3D recirculation flows that lead to the removal of pollutants
231 and mass-momentum exchange through the lateral boundaries of 3D streets (Madalozzo et al.,
232 [70]). 2D streets usually experience worse ventilation and higher pollutant concentrations
233 than 3D cases with similar aspect ratios and atmospheric conditions. For instance, studies
234 have reported a street intake fraction of 1-5 ppm in 2D street canyons where $AR=1$ whereas
235 the intake in 3D cubic building arrays was in the order of 0.1ppm. Despite corresponding

236 with the worst urban ventilation performance, 2D street canyon models are still commonly
237 employed to study and clarify the basic governing mechanisms in urban areas (e.g., [7-11,20-
238 23, 41-54, 56-57]). By simplifying the urban geometry as 2D, this study aims to build on the
239 existing literature and investigate the influence of various wall heating types and typical
240 aspect ratios on the fine details of flow pattern and pollutant exposure.

241 Large-eddy simulations (LES) [23, 30, 38-39, 42, 70] are known to outperform
242 Reynolds-Averaged Navier-Stokes (RANS) models [16, 57, 70-72] in predicting turbulence
243 and simulating urban flow and pollutant dispersion. Remaining challenges to the applications
244 of LES include a longer computational time, difficulty in specifying appropriate boundary
245 conditions at wall surfaces and a time-dependent inlet. Despite their limitations, RANS
246 approaches are still widely used [7-9, 14-20, 24-29, 31-35, 43-46, 50-52]. Among the RANS
247 models, the RNG $k-\varepsilon$ model has been one of the most widely adopted and has been
248 successfully validated in predicting flow and dispersion of gaseous pollutants [43-46, 50-52,
249 57, 73]. Chew et al. [57] reported that, RANS approaches perform well at reduced scales but
250 over-predict the thermal effects of heated windward walls at full scale, while LES predictions
251 agree closely with measurements at both scales. Considering both numerical accuracy and
252 computational time, the RNG $k-\varepsilon$ model was selected to solve the steady-state flow field and
253 pollutant dispersion in scale-model street canyons [43-46, 50-52,57].

254

255 **3.2 Model descriptions in the CFD test cases**

256 Fig. 2a shows the few 2D street canyon models that were built for numerical simulations.
257 The scale ratio of the simulated model to the full-scale model is 1:10. The building height (H)
258 of the CFD models is a constant of 3 m corresponding to the 30 m height of full-scale
259 buildings (10 floors). The width (W) of the target street canyon is set as 1m, 1.5 m, 3 m, 4.5

260 m or 6 m, which produces various aspect ratios: $AR=3$ and 2 (deep canyon); 1 (regular
261 canyon); 0.67 and 0.5 (avenue canyon), according to Vardoulakis et al. [10]. This selection of
262 the street width is to cover the aspect ratios from 0.5 to 3, which refer to regular and deep
263 street canyons, respectively. In the upstream and downstream of the target street canyon,
264 there are five identical street canyons to explicitly reproduce roughness elements at both sides
265 [e.g., 51, 54] (Fig. 2a).

266 In addition to the different aspect ratios, this study also investigated five different wall
267 heating scenarios of the target street canyon: an isothermal case as a controlled base case,
268 bottom heating, leeward wall heating, windward wall heating, and all wall heating. All of
269 these cases have the same temperature difference ($\Delta T=10$ K) between air and wall and
270 denote various thermal effects induced by solar radiation and wall heating. The model
271 descriptions of all of the test cases are listed in Table 2. Two mean wind speeds of 0.5 m s^{-1}
272 and 2 m s^{-1} were selected as the reference velocity (U_{ref}) at H of the domain inlet boundary
273 condition. The two U_{ref} values represent wind conditions with different Reynolds numbers,
274 and the case name follows the coding system: Case Heating type [AR , U_{ref}]. The heating types
275 of N, B, L, W and A represent no wall heating, bottom heating, leeward heating, windward
276 heating, and all wall heating. For example, N [0.5, 0.5] refers to the isothermal target street
277 canyon with an aspect ratio of 0.5 under a 0.5 m s^{-1} mean wind speed condition.

278 The reference Reynolds numbers ($Re = \rho U_{ref} H / \mu$, $H=3\text{ m}$) are 95602 at $U_{ref}=0.5\text{ m s}^{-1}$ and
279 382409 at $U_{ref}=2\text{ m s}^{-1}$ (Table 3) which are in the order of 10^5 to ensure Reynolds number
280 independence [55-56].

281 To characterize the effect of buoyancy force on turbulent airflow, the Froude number is
282 defined as:

$$F_r = \frac{U_{ref}^2}{\beta g H (T_w - T_{ref})} = \frac{U_{ref}^2}{g H \Delta T / T_{ref}} \quad (3)$$

where T_w is the surface temperature of the heated wall, $T_{ref} = 300\text{K}$ is the reference air temperature in the free stream and at the domain inlet and $\Delta T = 10\text{K}$ is a constant for all cases with wall heating. $\beta = 1/T_{ref}$ is the thermal expansion coefficient and g is the gravitational acceleration. The Froude number ranges from 0.25 to 4.08 (Table 3).

288

289 3.3 CFD setups for flow modelling

As shown in Fig. 2a, the 2D computational domain was built to be $23H$ in length and $6H$ in height. A total number of approximately 0.4 million cells were used. To capture the viscous sub-layer and heat transfer near the wall surfaces, the grid was refined toward the wall surfaces with a minimum grid size of 0.6mm (i.e., $2 \times 10^{-4}H$, see Fig. 2b). This grid arrangement is confirmed to be sufficiently refined by our CFD validation study with grid independence tests in subsection 4.1.

Table 4 summarizes boundary conditions and solver settings for the CFD simulations.

At the domain inlet, a power-law velocity profile $(U(y) = U_{ref} \times \left(\frac{y-y_{ref}}{y_{ref}}\right)^a)$, $U_{ref} = 0.5$ or 2.0m s^{-1} , $a=0.22$) and the profiles of turbulent kinetic energy and its dissipation rate were used as displayed in Table 3 [45-46]. No-slip wall boundary conditions with enhanced wall functions (EWF) were applied for near-wall treatment. The present grid is sufficiently refined close to the wall surfaces with a minimum grid size equal to 1mm (see Fig. 2b) to ensure that the dimensionless wall distance y^+ near walls is in order of 1 and satisfy the requirement of enhanced wall functions [38-40, 42-46]. This solves the viscous sub-layer near wall surfaces and heat transfer within it. Zero normal gradient conditions were used at the domain top (symmetry boundary) and domain outlet (fully developed outflow boundary).

306 The Boussinesq model was employed to assess the buoyancy effect [38-46, 49-51], in
 307 which the air density is treated as a constant except in the momentum equation of vertical
 308 velocity. The governing equations for the flow and turbulent quantities were implicitly
 309 discretized by a finite volume method (FVM) with the second-order upwind scheme to
 310 guarantee the numerical accuracy. The SIMPLE scheme was used for the pressure and
 311 velocity coupling.

312 The under-relaxation factors for the pressure term, momentum term, turbulent kinetic
 313 energy k , its dissipation rate ε and energy are 0.3, 0.7, 0.8, 0.8 and 1, respectively. CFD
 314 simulations do not stop until all of the residuals become constant. Typical residuals at
 315 convergence are 1×10^{-6} and 1×10^{-7} for U_x and U_y , respectively, 1×10^{-7} for continuity, $1 \times$
 316 10^{-6} for turbulent kinetic energy k , 1×10^{-6} for dissipation rate ε and 1×10^{-13} for the energy.

317

318 **3.4 CFD setups for pollutant dispersion modelling**

319 Apart from the solver setting shown above, the gaseous pollutant carbon monoxide (CO)
 320 was released from a pollutant line source with streamwise width of $W_x=0.038H=0.115\text{m}$
 321 which is positioned in the middle of the target street canyon at a height of 0.04m (Fig. 2a).
 322 Carbon monoxide (CO) was released with a small pollutant emission rate ($S_c=10^{-7} \text{ kg/m}^3\text{s}$) to
 323 ensure that the source release produced little disturbance to the flow field [25-27,61-63]. The
 324 geometry size and the pollutant emission rate were the same in all test cases. The sidewalks
 325 on both windward and leeward side represent the pedestrian regions with a height of 0.2 m,
 326 corresponding to 2 m height in full-scale.

327 The steady-state governing equation of CO concentration is:

$$328 \quad \bar{u}_j \frac{\partial \bar{c}}{\partial x_j} - \frac{\partial}{\partial x_j} \left(K_c \frac{\partial \bar{c}}{\partial x_j} \right) = S_c \quad (4)$$

329 where, \bar{c} is the pollutant concentration (kg/m^3), K_c is the turbulent eddy diffusivity of
330 pollutants, the S_c is the item of the pollutant source. According to $K_c = \nu_t / Sc_t$, ν_t is the
331 kinematic eddy viscosity and Sc_t is the turbulent Schmidt number, which is treated as a
332 constant of 0.7 [25-27, 31, 61-63] .

333 In Eq. (4), the zero normal flux condition was used at the wall surfaces, and a zero
334 normal gradient condition was applied at the domain outlet and domain roof. At the domain
335 inlet, the concentration was null.

336 To quantify the pollutant dispersion, the CO concentration was normalized as follows:

$$337 \quad K = \bar{c} U_{ref} L W_s / Q \quad (5)$$

338 where U_{ref} is taken as a constant of 0.5 m s^{-1} or 2.0 m s^{-1} for all test cases, L and W_s are the
339 source length and source width and Q is the total mass release rate (kg/s).

340 In summary, the personal intake fraction (P_{IF}) was analyzed based on the results for
341 pollutant concentration. As described and assumed in Section 2, only one microenvironment
342 (i.e. indoor at home) was considered, and the concentrations along the windward-side wall
343 and leeward-side wall of near-road buildings were emphasized for quantifying the
344 concentration of the microenvironment ("indoor at home") originating from outdoor vehicular
345 pollutant emissions (see Fig. 2a and Table 1).

346

347 **4. Results and Discussions**

348 **4.1 Validation of flow and pollutant dispersion modelling**

349 For flow modelling in regular street canyons ($AR=H/W=1$) with wall heating conditions,
350 CFD simulations were validated by the wind tunnel data reported by Allegrini et al. [49]. The
351 details of the validation procedure can be found in our previous research (Lin et al. [45]) and

352 Appendix A1. To validate the finding of passive pollutant dispersion in the regular and deep
353 street canyons ($AR=1$ and 2), we compared the CFD simulation results with the concentration
354 distribution in the wind tunnel data from Meroney et al. [22]. Further detailed description can
355 be found in He et al. [54] and Appendix A2. Moreover, scale-model outdoor field
356 experiments ($H=1.2\text{m}$) were carried out to confirm that only one main vortex existed in the
357 deep street canyon where $AR=2$ and 3 as the background wind speed was sufficiently high; in
358 other words, where the wind-driven dynamic force is dominant and buoyancy force is
359 relatively weak. Descriptions of the scale-model experiments have been introduced in detail
360 in Zhang et al. [9] and Appendix A2. Finally, in Appendix A4, CFD validation and grid
361 independence study of flow modelling in scaled deep street canyon ($AR=2.4$, $H=1.2\text{m}$,
362 $U_{ref}=13\text{m s}^{-1}$, $W=B=0.5\text{m}$, $Re\sim 10^6$) are conducted under the estimation by wind tunnel data
363 ($AR=2.4$, $H=12\text{cm}$, $U_{ref}=13\text{m s}^{-1}$, $W=B=5\text{cm}$, $Re\sim 10^5$). The results of the above validation
364 tests indicated that the CFD simulations presented in this study have a satisfactory
365 performance and agree well with the experimental data.

366 In the following sections, firstly, we discuss the effects of street aspect ratios ($AR=0.5-3$),
367 wall heating and Froude numbers ($U_{ref}=2$ and 0.5m s^{-1} , $\Delta T=10\text{K}$, $Fr=0.25$ and 4.08) on
368 urban airflow, pollutant distribution and personal exposure (i.e. street intake fraction) within
369 street canyons.

370 **4.2 *Re*-number independence evaluation**

371 The airflow characteristics within the reduced-scale street canyons are different from
372 that within full-scale models if *Re*-number-independence cannot be satisfied. To verify this
373 issue, we performed additional scaled CFD simulations with various background wind speeds
374 ($U_{ref}=0.5$, 2 and 4m s^{-1} , $H=3\text{m}$) and reference Reynolds number ($Re=95602$, 382409 ,
375 764818) to verify whether the Reynolds numbers independence is satisfied or not in the
376 isothermal cases where $AR=0.5-3$. Fig. 3a shows that, background wind speed of 0.5m s^{-1}

377 and 2 m s^{-1} ($Re=95602$ and 382409) are sufficient to ensure Reynolds numbers independence
 378 as $AR=0.5, 0.67, 1,$ and $2,$ however as $AR=3$ (Fig. 3b-c), the flow with $Re=382409$ ($U_{ref}=2\text{ m}$
 379 s^{-1}) and 764818 ($U_{ref}=4\text{ m s}^{-1}$) are Re -number independent but that with $Re=95602$ ($U_{ref}=0.5$
 380 m s^{-1}) is not. Therefore, in the following sections, the analyses mainly focus on the cases as
 381 $U_{ref}=0.5$ and 2 m s^{-1} .

382 **4.3 Flow and pollutant dispersion as $U_{ref}=2\text{ m s}^{-1}$ and $Fr=4.08$**

383 All of the CFD simulations reported in this subsection were conducted under the
 384 condition of a high wind speed condition ($U_{ref}= 2.0\text{ m s}^{-1}$ and $Fr=4.08$) with a relatively weak
 385 buoyancy force ($Fr=4.08$), which satisfied the Re -number independence.

386 **4.3.1 In deep street canyons with aspect ratio of $AR =3$ (Re -number independence)**

387 Fig. 4a shows the distribution of the mean wind speed and normalized pollutant
 388 concentration (K) of the deepest street canyon ($AR=3$) under five different heating scenarios.
 389 A single clockwise vortex formed in Cases N[3, 2] (No heating) and L[3, 2] (Leeward wall
 390 heating). The flow patterns of the remaining cases formed a multi-vortex structure with a
 391 worse pollutant dilution capacity and higher concentrations near the ground (three vortexes
 392 for Case W[3, 2], and Case A[3, 2] and two vortexes for Case B[3, 2]). Due to the high AR of
 393 the street canyon, the approaching wind has difficulty entering the space inside the canyon,
 394 especially at the street level. Therefore, the mean wind speeds at the pedestrian level of the
 395 street canyons are relatively small ($<0.2\text{ m s}^{-1}$). Fig.4b-4d summarize the wind speed
 396 distribution (u_x and u_y) at the windward line, leeward line and the line near the bottom. Fig.4e
 397 shows the normalized CO concentration (K) along the windward and leeward wall.

398 Due to the different airflow patterns within street canyons, the distributions of K also
 399 vary under different heating scenarios. As shown in Fig. 4a and 4e, the single clockwise
 400 vortex in Case N[3, 2] and L[3, 2] results in a higher K at the leeward wall than the windward

401 wall. The mean wind speeds u_y below $y/H=0.6$ of these two cases are relatively higher, which
 402 tends to reduce K within the street canyons more effectively (Fig. 4e). Unlike with the single
 403 vortex pattern, the formation of multiple vortices significantly aggravated the air pollutant
 404 dispersion near the ground, as presented in Case W[3, 2], B[3, 2] and A[3, 2] (Fig. 4a and 4e).
 405 In other words, the vertical buoyancy force even worsens the air pollution near the ground
 406 under high wind speed conditions ($Fr=4.08$). For example, the concentration K near the
 407 leeward-side ground ($y/H<0.2$) rapidly increases to 1500 in Case B[3,2] and 1750 in Case
 408 A[3, 2] (Fig. 4e).

409 **4.3.2 In deep street canyons with aspect ratio of $AR = 2$ (Re -number independence)**

410 Fig. 5 shows the distributions of wind speed and K in the street canyon where $AR=2$. In
 411 contrast to the street canyon where $AR=3$, the wind-driven force becomes a more dominant
 412 factor in flow pattern formation and pollutant dispersion in the wider canyons where $AR=2$.
 413 As shown in Fig. 5a, only one clockwise vortex can be observed in Case N[2, 2], B[2, 2] and
 414 L[2, 2] and two vortices with opposite directions were formed in Case W[2, 2] and A[2, 2].
 415 In cases with one main vortex, the leeward-side K is much higher than that near windward-
 416 side. Both u_x and u_y of Case N[2, 2] are slightly smaller than those of Case B[2, 2] and L [2,
 417 2] (Fig.5b-5d). This phenomenon indicates that bottom heating and leeward wall heating can
 418 slightly strengthen the turbulent flow and pollutant dilution capacity. In Case A[2, 2] and
 419 W[2, 2], there is a stronger main vortex at the upper levels and a much weaker one at low
 420 levels (Fig. 5a), which produces a higher K value near the windward side (below $y/H=0.5$),
 421 but a smaller K value at the upper level. The peak of the K values in the two-vortex cases is
 422 much higher (maximum 400) than those of one-vortex cases (maximum 230) (Fig. 5e).
 423 Overall, the K values where $AR=2$, varied between 80 to 400 are much lower than that in the
 424 deeper street canyons where $AR=3$ (Fig. 4e).

425

426 **4.3.3 In regular street canyons with an aspect ratio of $AR = 1$ (Re -number independence)**

427 When the street aspect ratio was further reduced to $AR=1$ with $Fr=4.08$ ($U_{ref}=2\text{m s}^{-1}$), as
428 shown in Fig. 6, only one clockwise vortex was formed in all five cases and the leeward-side
429 K was much higher than for the windward side. The velocity and K profiles were close in all
430 cases with different heating scenarios (Fig. 6c-6e), confirming that the buoyancy force hardly
431 changed the flow and dispersion pattern. The K values were around 35 on the windward side
432 and 90 on the leeward side (Fig. 6f), which are much lower than those with $AR=2$ (Fig. 5f)
433 and 3 (Fig. 4f).

434 Overall, when $Fr = 4.08$, and AR was reduced from 3 and 2 to 1, the wind-driven force
435 became more dominant than the buoyancy force. Therefore, the flow pattern and pollutant
436 dispersion in street canyons where $AR=0.67$ and 0.5 is not shown here. However, the personal
437 intake fraction is analyzed in sub-section 3.4

438

439 **4.4 Pollutant dispersion under low wind speed condition ($U_{ref}=0.5\text{m s}^{-1}$ and $Fr=0.25$)**

440 **4.4.1 In deep street canyons with an aspect ratio of $AR = 3$**

441 As we discussed in the section 4.2, the cases with the AR of 3, are not satisfied the Re -
442 number independence under the condition of a relatively weak background wind ($U_{ref}=0.5\text{m}$
443 s^{-1}). The flow patterns and pollutant distributions are more complicated in street canyons
444 under the low wind speed conditions than under high wind speed conditions ($U_{ref}=2\text{m s}^{-1}$ and
445 $Fr=4.08$). Due to the relatively low Reynolds number ($Re=95602$), two vortexes moving in
446 the opposite directions were formed in the isothermal Case N [3, 0.5], and the pollutants near
447 the ground were transported to the windward side. In Case B [3, 0.5]), a stronger, narrower
448 vortex was formed at the leeward bottom side and the upper vortex was smaller compared
449 with that in Case N [3, 0.5]. There was a similar phenomenon in Case W [3, 0.5] but with an

450 even stronger vortex at the bottom side, which resulted in pollutant transportation to the
451 windward side. Heating of the leeward wall (L [3, 0.5]) modified the flow pattern within the
452 street canyon significantly, where the flow pattern was dominated by a single clockwise
453 vortex, as shown in Fig. 7a. When the walls and ground were all heated (A[3, 0.5]), three
454 main vortexes were formed, with two clockwise vortexes located at the left side and one
455 counter-clockwise vortex located at the right side. However, the core of the vortex was
456 mostly located at the upper level of street canyons, so although the formations of multiple
457 vortexes may assist the dilution of pollutants in the entire street canyons, the pollutants would
458 accumulate in the middle and lower levels because of the lower wind speeds there.

459 Fig. 7b-7c show that the vertical wind speed u_y in Case L and W [3, 0.5] is higher than
460 that in other cases. Wind circulation with higher wind speed results in the lower
461 concentration of pollutants, as shown in Fig. 7a. Compared to Case N [3, 2], the insufficient
462 wind speed u_y in Case N [3, 0.5] leads to the accumulation of air pollutants. The K value at
463 the bottom level of Case N [3, 0.5] reached 1500, as shown in Fig. 7e. The buoyancy effect is
464 the dominant force driving wind flow and the pollutant dispersion within the street canyon
465 under a low wind speed. This resulted in pollutant accumulation at street level in Case N [3,
466 0.5]. This can also be observed in Case N [2, 0.5], in which the wind force became more
467 important with the decrease of AR and the pollutants had the potential to be spread along
468 the leeward wall. Case N [2, 0.5] is further discussed in the following sections.

469 **4.4.2 In deep street canyons with an aspect ratio of $AR = 2$ (Re -number independence)**

470 The effects of wind flow became more distinct when the AR value decreased to 2 and the
471 case tends to satisfy the Re -number independence with the decrease of AR (Fig. 3) in the
472 presence of background wind of 0.5 m s^{-1} . As shown in Fig. 8a, the wind speed contour maps
473 of Case N [2, 0.5] showed the single clockwise vortex pattern within the canyons. This
474 pattern significantly improved pollution levels at street level. Compared to Case N [3, 0.5],

475 the maximum K value of Case N [2, 0.5] decreased from 1500 to 900, as shown in Fig. 8e.
476 The flow patterns of Case B [2, 0.5], W [2, 0.5], L [2, 0.5] and A [2, 0.5] were similar to the
477 patterns of the corresponding cases with $AR=3$. Specifically, the heated wall at the leeward
478 side enhanced the wind circulation of the single clockwise vortex shown in Fig. 8a, and the
479 heating windward wall is the main reason for the formation of the large counter-clockwise
480 vortex observed in Case W [2, 0.5].

481 The modified flow patterns affect the CO dispersion within the canyon. The pollutant
482 level at the leeward side was higher than that at windward side in Case N [2, 0.5] and L [2,
483 0.5]. However, the large vortex moving in the opposite direction from the one in Case W [2,
484 0.5] drove the pollutant accumulation at the windward side. The two large vortexes of Case B
485 [2, 0.5] caused the K value to decrease sharply at the height of $y/H=0.5$, shown in Fig. 8e.
486 This is due to the variation in the u_y value at the windward and leeward sides shown in Fig.
487 8b-8c. Comparing Case B [3, 0.5] and N [2, 0.5], Case B[2, 0.5] shows the strengthened
488 upper level vortex partly compresses the development of the bottom vortex as the aspect
489 ratios decrease,, resulting in a higher concentration K below $y/H=0.5$.

490 **4.4.3 In regular street canyons with an aspect ratio of $AR= 1$ (Re -number independence)**

491 For the street canyon with an AR of 1, the forces with the highest impact are the wind
492 force and buoyancy effect, as shown in Fig. 9a. As in Case N [2, 0.5], a single vortex formed
493 in the canyon of Case N [1, 0.5], which was affected by wind force only. The pollutants
494 assembled at the leeward side in Case N [1, 0.5] as occurred in Case L [1, 0.5]. The
495 concentration differences between the two cases were due to the buoyancy force near the
496 leeward wall in Case L [1, 0.5], which enhanced the circulation in the street canyons and
497 removed the pollutants from the street level (Fig. 9b and 9f). For Case B [1, 0.5], the bottom
498 heating generated the buoyancy force, which formed another vortex to the right of the center.
499 The two opposite vortexes brought about the dramatically change of K values at $y/H = 0.7$, as

500 shown in Fig. 9f. The vertical buoyancy force is another reason for the right vortex formation
501 observed in Case W [1, 0.5].

502 **4.4.4 In avenue street canyons with an aspect ratio of $AR=0.67$ (Re -number** 503 **independence)**

504 As the aspect ratio of the street canyons decreased further, the impacts of wind force
505 increased. As shown in Fig. 10a, Case B [0.67, 0.5], N [0.67, 0.5], and L [0.67, 0.5] all have
506 the single vortex flow pattern. The wind circulation of Case B [0.67, 0.5] is even stronger
507 than that of Case L [0.67, 0.5], and the u_x and u_y values of the two cases are higher than the
508 mean wind speeds in Case N [0.67, 0.5], as shown in Fig. 10c-10e. A right vortex appeared in
509 both Case W [0.67, 0.5] and Case A [0.67, 0.5] but differed in sizes. The smaller vortex of
510 Case A [0.67, 0.5] indicated that the buoyancy forces near the ground compressed this vortex
511 into a smaller one, as the buoyancy near the windward side and ground drove the vortex in
512 the opposite direction.

513 From Fig. 10f, we can see that the vertical distribution pattern of the K values for the
514 five cases. Case A [0.67, 0.5] and B [0.67, 0.5] had the lower pollutant concentration near the
515 buildings, and Case N [0.67, 0.5] had the worst air quality within the canyon. Compared to
516 Case N [0.67, 0.5], Case W [0.67, 0.5] had a weaker wind circulation but a lower pollutant
517 concentration at the leeward side. This is because the pollutants were not be dispersed
518 throughout the entire canyon, especially at the windward side, under the low wind speed
519 conditions ($<0.5\text{ m s}^{-1}$) with $Fr=0.25$. This caused a lower pollutant concentration at the
520 leeward side in Case W [0.67, 0.5], shown in Fig. 10b and 10f.

521 **4.4.5 In avenue street canyons with an aspect ratio of $AR=0.5$ (Re -number**
522 **independence)**

523 For the wider street canyon with $AR=0.5$, the flow patterns and the pollutant
524 concentration distributions for the five cases are shown in Fig. 11a-11b. The detailed vertical
525 and horizontal mean wind speed in various locations and the vertical K values near the
526 buildings are shown in Fig. 11f. The flow patterns and the K distributions are quite similar to
527 those in the previous cases with an aspect ratio of 0.67. In the previous studies [21,70] the
528 flow patterns within the wide street canyons were clarified, with $AR<0.66$ as reference, as the
529 wake interference flow, which was not observed in this study. This is because the Re number
530 of this study is different from the previous number. With a higher Re number, an even smaller
531 AR value is required to transform the flow pattern from a skimming flow to a wake
532 interference flow.

533 As discussed above, the buoyancy effect can be very effective in removing the air
534 pollutants within the street canyons under the low wind speed conditions. Greater street width,
535 corresponding with a smaller AR value, allows wind into the street level. That wind flow can
536 further modify the pollutant dispersion inside the canyon by, for example, decreasing the
537 pollutant concentration. The aspect ratio, background wind speed, and the heating scenario
538 are the three important factors that should be carefully considered by the urban planners and
539 engineers when designing urban environments.

540

541 **4.5 Effects of different heating conditions on the personal intake fraction of CO**

542 The patterns of pollutant dispersion within the street canyons are mainly determined by
543 the street aspect ratio, heating scenario, and background wind speed conditions. Under high
544 wind speed conditions ($Fr>1$, $Re=382409$), the wind force almost acts as the dominant factor

545 forming the flow patterns within the canyon with lower AR ($AR=1-0.5$). As the wind speed
546 decreases ($Fr<1$, $Re=95602$), the dominant force is switched to a buoyancy force within the
547 canyons with a higher AR ($AR=3-1$). These characteristics lead to the non-uniform
548 distribution of pollutants within the street canyons, as shown in the previous analysis.
549 Considering that factors such as different human activities, various durations of stay and
550 breathing rates in different microenvironments would affect the amount of pollutants inhaled
551 by urban residents, one average value of pollutant concentration cannot represent the real
552 pollutant exposure in local streets or districts. This is the reason why P_{IF} (personal intake
553 fraction) was applied in this study to evaluate the effect of the heating conditions and the AR
554 on personal exposure. A higher P_{IF} value refers to a higher amount of pollutants inhaled by
555 pedestrians. Fig. 12 gives the detailed variation in P_{IF} value under high wind speed
556 conditions (Fig. 12a-12b) and low wind speed conditions (Fig. 12c-12d). The P_{IF} value is
557 higher in the street canyons with a higher aspect ratio under both high and low wind speed
558 conditions. The following will discuss separately discuss the P_{IF} values calculated under
559 the high and low background wind conditions respectively.

560 Within the narrow street canyon ($AR=3$), the approaching wind had difficulty
561 penetrating the street, even under high wind speed conditions. This means that the buoyancy
562 effect is one of the key factors impacting the pollutant dispersion. For an isothermal case
563 where $AR=3$, the P_{IF} is 3.96 ppm, as shown in Fig. 12a. The P_{IF} value was the lowest in
564 Case L [3, 2] at $P_{IF}=3.72$ ppm. This result that the heating leeward side would enhance the
565 single vortex, which carries the pollutant to the upper level of the street, is consistent with the
566 results shown in Fig. 4a. The wind and buoyance forces both present in the flow patterns of
567 Case B [3, 2], W [3, 2] and A [3, 2] competed, resulting in the multiple vortexes within the
568 canyon. The variation in the P_{IF} value follows the flow features discussed in Fig. 4a, where
569 the P_{IF} s are 6.13 ppm, 8.63 ppm and 10.32 ppm for Case W [3, 2], B [3, 2] and A [3, 2],

570 respectively. For the cases with an AR value of 2, where more wind can blow at street level,
571 the overall P_{IF} value decreases by 40% compared to the cases with an AR of 3. The flow
572 patterns of Case W [2, 2] and Case A [2, 2] contained two vortexes, which brought about the
573 higher P_{IF} values ($P_{IF}=3.33$ and 3.18 ppm, respectively) compared to the close range for
574 P_{IF} in Case N [2, 2], B [2, 2] and L [2, 2] (P_{IF} ranges from 2.13-2.21 ppm). For the wider
575 street canyons with AR s of 1, 0.67 and 0.5, the influence of wind was more dominant, so the
576 P_{IF} values were lower and relatively constant under different heating scenarios (P_{IF} s were
577 around 0.91-0.98 ppm, 0.57-0.60 ppm and 0.39-0.43 ppm where $AR=1, 0.67$ and 0.5), as
578 shown in Fig. 12b.

579 Under low wind speed conditions ($U_{ref}=0.5\text{m s}^{-1}$ and $Fr=0.25$), the overall P_{IF} values
580 of all cases increased by 30% to 50%, compared to the case under the same heating
581 conditions but a higher wind speed (Fig. 12). For cases with an AR value of 3, the P_{IF}
582 values were mainly affected by the buoyancy effects, but the P_{IF} values of the rest of the
583 cases were determined by both wind and buoyancy force, as shown in Fig. 12c-12d. The
584 maximum P_{IF} value was 27.51 ppm, appearing in the Case N [3, 0.5]. For the narrow
585 streets, the low approaching wind and lack of heating conditions made the air movement very
586 weak inside the street canyon. The almost static airflow stopped the pollutant dilution and
587 resulted in the highest P_{IF} value. The buoyancy effect from heating increased the strength
588 of the bottom vortex in Case L [3, 0.5] and Case W [3, 0.5] (Fig. 10), and therefore the P_{IF}
589 values decreased to 7.10 and 9.56 ppm, respectively. The vortex structures of Case A [3, 0.5]
590 and B [3, 0.5] were multi-vortex patterns due to the heating walls and the vortexes at the
591 pedestrian level were relatively weak. This means that a multiple-vortex situation deteriorates
592 the air quality, resulting in a higher P_{IF} values of 15.39 and 14.77 ppm, respectively. For all
593 the cases with an AR value of 2, the P_{IF} values decreased due to more wind entering the
594 canyon. When the AR value further decreased to 1, 0.67 and 0.5, the wind became the main

595 force to modify the flow features within the canyons, and the P_{IF} values were further
596 decreased to approximately 0.71-1.66 ppm.

597 Above all, the variation in the overall P_{IF} value was determined by the flow patterns
598 within the street canyons. The single-vortex pattern was more efficient in removing pollutants
599 at street level than the multi-vortex flow. Whether the buoyancy effect enhances or worsens
600 the pollutant dilution capacity depends on the aspect ratio and wall-heating types of the street
601 canyons. A lower background wind speed usually results in higher pollutant exposure. As
602 $U_{ref}=0.5\text{m s}^{-1}$ ($Re=95602$, $Fr=0.25$), deep street canyons ($AR=2$ and 3) with no wall heating
603 (cloudy day), all wall heating (nighttime with urban heat island effects) and bottom heating
604 (at noon) experienced a larger street intake fraction than leeward or windward heating.
605 Regular and avenue street canyons ($AR=1$ and $0.5-0.67$), windward wall heating and no wall
606 heating produced greater pollutant exposure. Leeward wall heating always improved
607 pollutant dilution and reduced the street intake fraction. Overall, if the background wind
608 speed is relatively lower and the buoyancy force has significant effects, attention should be
609 paid to deep street canyons (e.g. $AR=2-3$) at nighttime (all-wall heating), noon (bottom
610 heating) and cloudy weather (no wall heating). Similarly, regular and avenue street canyons
611 with $AR=0.5-1$ are of particular concern during windward-wall heating and periods of cloudy
612 weather.

613

614 **4.6 Limitations and future researches**

615 As pointed out by Chew et al. [56-57], the reduced-scale street canyons may experience
616 different findings from full-scale models if Re-number-independence cannot be satisfied.
617 Therefore, in the near future, we will conduct CFD simulations from the wind-tunnel scale
618 ($H\sim 0.1$ m, $Re\sim 10^4-10^5$) to scaled model ($H\sim 1$ m, $Re\sim 10^5-10^6$) and full-scale models ($H\sim 10$
619 m-100 m, $Re\sim 10^6\sim 10^7$) under the validation using experimental data. The critical Re number

620 (Re_c) required for Re number independence will be quantified and compared in cases with
621 various aspect ratios (e.g. $AR=0.5-6$) with the coupling effect of dynamic force and thermal
622 buoyancy force on the flow and pollutant exposure in street canyons.

623 In the heating scenarios, the heating of the building roof is not considered. Although the
624 wind speed above the building roof is considerable and the buoyancy force induced by roof
625 heating is not significant when the background wind speed is relatively large, such impacts
626 cannot be disregarded in calm weather condition. In addition, more realistic urban heating
627 scenarios have been used in several CFD studies by the literature [42, 73-77] in which the
628 integrated impacts of urban turbulence and radiation processes with partially-heated walls
629 determined by solar angles are considered. Furthermore, the scaled outdoor experiments in
630 Appendix 3 verify that the realistic background wind speed and direction may vary with time,
631 thus the influence of such unsteady boundary conditions and indoor-outdoor interactions
632 coupling with radiation processes and/or wall heating scenarios on urban turbulence,
633 pollutant dispersion and pollutant exposure should be further investigated under the high-
634 quality scaled outdoor experimental data (e.g., Fig. A3e, measured data from Appendix A3).

635

636 5. Conclusion

637 Deep street canyons and unfavourable meteorological conditions (e.g., a weak
638 background wind) are the main factors producing poor ventilation capacity, a high pollutant
639 exposure of urban residents and the related adverse impacts on human health. This study
640 focuses on the impact of aspect ratios ($AR= 3, 2, 1, 0.67, 0.5$; $H=3m$), background wind
641 speeds ($U_{ref}=2m\ s^{-1}$ and $0.5m\ s^{-1}$) and various wall-heating scenarios ($T_{wall}-T_{air}=10K$) on air
642 flow, pollutant dispersion and the related human exposure in scale-model street canyon
643 models, which has not received significant research attention. Various Froude numbers
644 ($Fr=0.25$ and 4.08) and reference Reynolds numbers ($Re=95602$ and 382409) were

645 considered. The use of CFD methodologies combined with a RNG $k-\varepsilon$ model has been
646 validated by wind tunnel data and scale-model outdoor field experiments. The personal intake
647 fraction (P_{IF}) and its spatial mean value for an entire street, or street intake fraction $\langle P_{IF} \rangle$,
648 are used to quantify personal exposure in near-road buildings.

649 In most isothermal cases, only one-main-vortex structure exists when $AR=0.5-3$, but two
650 vortexes appear for $AR=3$ and $Re=95602$, confirming that $Re=95602$ cannot satisfy the Re
651 independence requirement when $AR=3$.

652 In non-isothermal cases with $Fr=4.08$ and $U_{ref}=2\text{m s}^{-1}$ ($Re=382409$), the most salient
653 features is that the formation of a single vortex removes the pollutant efficiently; however,
654 the formation of a multi-vortex structure due to different heating scenarios increases $\langle P_{IF} \rangle$
655 to a certain extent, where $AR=2-3$. As $AR=0.5-1$, the wind dynamic force dominates the flow
656 patterns in street canyons and the buoyancy effect is less important. The four heating
657 conditions attain similar $\langle P_{IF} \rangle$ in isothermal cases (0.91-0.98 ppm, 0.57-0.60 ppm, 0.39-
658 0.43 ppm for $AR=1, 0.67, 0.5$ respectively).

659 In contrast to the isothermal case as $AR=3$, leeward-wall-heating slightly enhances the
660 single-main-vortex structure and slightly reduces $\langle P_{IF} \rangle$ (3.96 ppm to 3.72 ppm), but other
661 heating scenarios induce a multi-vortex structure that significantly increases pollutant
662 exposure ($\langle P_{IF} \rangle=3.96$ ppm to 6.13-10.32 ppm). When $AR=2$, bottom or leeward wall
663 heating only slightly affects the single vortex, resulting in a similar $\langle P_{IF} \rangle$ (2.13-2.21 ppm)
664 but windward and all-wall heating creates multi-vortex structures, resulting in an increased
665 $\langle P_{IF} \rangle$ (3.18-3.33 ppm).

666 When $Fr=0.25$ and $U_{ref}=0.5\text{m s}^{-1}$, the isothermal case where $AR=3$ experiences the
667 highest $\langle P_{IF} \rangle$ (27.51 ppm), and $\langle P_{IF} \rangle$ decreases with the decrease of AR (7.85 ppm, 3.47
668 ppm, 2.30 ppm and 1.66ppm where $AR=2, 1, 0.67, 0.5$). The four heating condition all
669 significantly influence vortex structure. Leeward wall heating always enhances pollutant

670 dilution and results in a lower $\langle P_{IF} \rangle$ than in the isothermal case (i.e. 7.10 ppm, 4.41 ppm,
671 2.29 ppm, 1.57 ppm, 1.20 ppm where $AR=3, 2, 1, 0.67, 0.5$), but the influence of the other
672 three heating conditions is complicated. Where $AR=0.67$ and 0.5 , the other three heating
673 conditions will improve the air quality (Fig. 12). However, where $AR=2$, the bottom wall
674 heating results in a higher $\langle P_{IF} \rangle$ (10.07 ppm) compared to the isothermal case where $AR=2$.
675 Where $AR=1$, both the bottom and windward heating will increase the $\langle P_{IF} \rangle$ to 3.51 and
676 4.52 ppm, respectively. The flow patterns and pollutant dispersion under weak conditions
677 also depend on the competition between the wind-driven dynamic force and buoyant force.

678 In general, a single vortex pattern is more efficient in removing the pollutants at the
679 street level for both high and low wind speeds. Leeward wall heating always enhances the
680 circulation in street canyons where $AR=0.5-3$. The buoyancy effect induced by other wall
681 heating scenarios can sometimes raise or reduce pollutant exposure, depending on the aspect
682 ratios, ambient wind speed and wall-heating types. Lower background wind speeds merit
683 more attention, since they usually result in a higher pollutant exposure. Certain other
684 conditions require particular attention: $U_{ref}=0.5\text{ m s}^{-1}$ ($Re=95602, Fr=0.25$), and deep street
685 canyons (e.g., $AR=2-3$) at nighttime (all-wall heating), at noon (bottom heating) and in cloudy
686 weather periods (no wall heating); while regular and avenue street canyons with $AR=0.5-1$
687 need more attention during windward-wall heating and cloudy weather periods.

688 Further investigations are still required before providing guidelines for design purposes,
689 but this study serves as one of the first attempts to evaluate the influence of various wall
690 heating and aspect ratios on pollutant exposure in urban street. The methods adopted in this
691 study can be used to assess the street intake fraction in more complicated urban streets or
692 neighborhoods under a variety of atmospheric conditions.

693

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704

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Journal Pre-proof

904 **Figure list**

905 Fig. 1. Breathing rate and time patterns for various age groups and microenvironments [64-
906 65].

907 Fig. 2. (a) Dimensions of the simulated street canyon model in CFD. (b) The grid
908 arrangement of 2D CFD simulations.

909 Fig. 3. Normalized stream-wise velocity along the street centerline in isothermal cases with
910 (a) background wind speed of $U_{ref}=0.5\text{ m s}^{-1}$ and 2 m s^{-1} ($Re=95602$ and 382409) where
911 $AR=0.5, 0.67, 1,$ and $2,$ (b) $U_{ref}=0.5\text{ m s}^{-1}, 2\text{ m s}^{-1}, 4\text{ m s}^{-1}$ where $AR=3$ ($Re=95602,$
912 $382409, 764818$). (c) Normalized stream-wise velocity and streamline in deep street
913 ($AR=3$) in isothermal case.

914 Fig. 4. In cases where $AR=3, U_{ref}=2.0\text{ m s}^{-1}$: (a) Contour of the mean wind speed (m s^{-1}) and
915 normalized concentration K , Vertical velocity u_y , along (b) the windward line and (c) the
916 leeward line, (d) streamwise velocity u_x along the bottom line. (e) Spatial average $\langle K \rangle$
917 along the windward wall and leeward wall.

918 Fig. 5. In cases where $AR=2, U_{ref}=2.0\text{ m s}^{-1}$: (a) Contour of the mean wind speed (m s^{-1}) and K ;
919 u_y along (b) the windward line and (c) the leeward line; (d) u_x along the bottom line; (e)
920 $\langle K \rangle$ at the windward and leeward walls.

921 Fig. 6. In cases where $AR=1$ and $U_{ref}=2\text{ m s}^{-1}$: Contour of (a) the mean wind speed (m s^{-1})
922 and K . u_y along (c) the windward line and (d) the leeward line; (e) u_x along the bottom
923 line; (f) $\langle K \rangle$ at the windward and leeward walls.

924 Fig. 7. In cases where $AR=3$ and $U_{ref}=0.5\text{ m s}^{-1}$: (a) Contour of the mean wind speed (m s^{-1})
925 and K . u_y along (b) the windward line and (c) the leeward line; (d) u_x along the bottom
926 line; (e) $\langle K \rangle$ at the windward and leeward walls.

927 Fig. 8. In cases where $AR=2$ and $U_{ref}=0.5\text{m s}^{-1}$: (a) Contour maps of the mean wind speed (m
 928 s^{-1}) and $K. u_y$ along (b) the windward line and (c) the leeward line; (d) u_x along the
 929 bottom line; (e) $\langle K \rangle$ at the windward and leeward walls.

930 Fig. 9. In cases where $AR=1$ and $U_{ref}=0.5\text{m s}^{-1}$: Contour of (a) the mean wind speed (m s^{-1})
 931 and (b) $K. u_y$ along (c) the windward line and (d) the leeward line; (e) u_x along the
 932 bottom line; (f) $\langle K \rangle$ at the windward and leeward walls.

933 Fig. 10. In cases where $AR=0.67$ and $U_{ref}=0.5\text{m s}^{-1}$: Contour of (a) the mean wind speed (m s^{-1})
 934 and (b) $K. u_y$ along (c) the windward line and (d) the leeward line; (e) u_x along the
 935 bottom line; (f) $\langle K \rangle$ at the windward and leeward walls.

936 Fig. 11. In cases where $AR=0.5$ and $U_{ref}=0.5\text{m s}^{-1}$. Contour maps of (a) the mean wind speed
 937 (m s^{-1}) and (b) $K. u_y$ along (c) the windward line and (d) the leeward line; (e) u_x along
 938 the bottom line; (f) $\langle K \rangle$ at the windward wall and leeward wall.

939 Fig. 12. Spatial mean value of the personal intake fraction of a local street with different
 940 heating conditions as $AR=0.5-3$ with (a) $U_{ref}=2.0\text{m s}^{-1}$, and (b) $U_{ref}=0.5\text{m s}^{-1}$.

941

942

943 **Table list**

944 Table 1 Breathing rate and time patterns for indoor at home for each age group[58, 64-65].

945 Table 2. Model descriptions of the simulated test cases.

946 Table 3. Boundary conditions and solver settings for the CFD simulations

947 Table 4. Reynolds and Froude numbers investigated in all test cases with wall heating (T_{wall} -

948 $T_{\text{air}}=10\text{K}$)

949

950 **Appendix: CFD validation by experimental data**

951

952 **Appendix A1. Flow validation for 2D street canyon with wall heating by**

953 **wind tunnel experiment ($AR=1$)**

954 The CFD simulations were first evaluated using wind tunnel data from the work of Allegrini
955 et al. [49], which studied the flow and turbulence characteristics within a street canyon

956 ($W=H=0.2\text{m}$) under isothermal and non-isothermal conditions (Fig. A1). Four situations,
957 including leeward wall heating, windward wall heating, ground heating and all wall heating,

958 were investigated. The Froude number ($F_r = \frac{U_{ref}^2}{\beta g H (T_w - T_{ref})}$) ranges from 0.65 to 17.29,

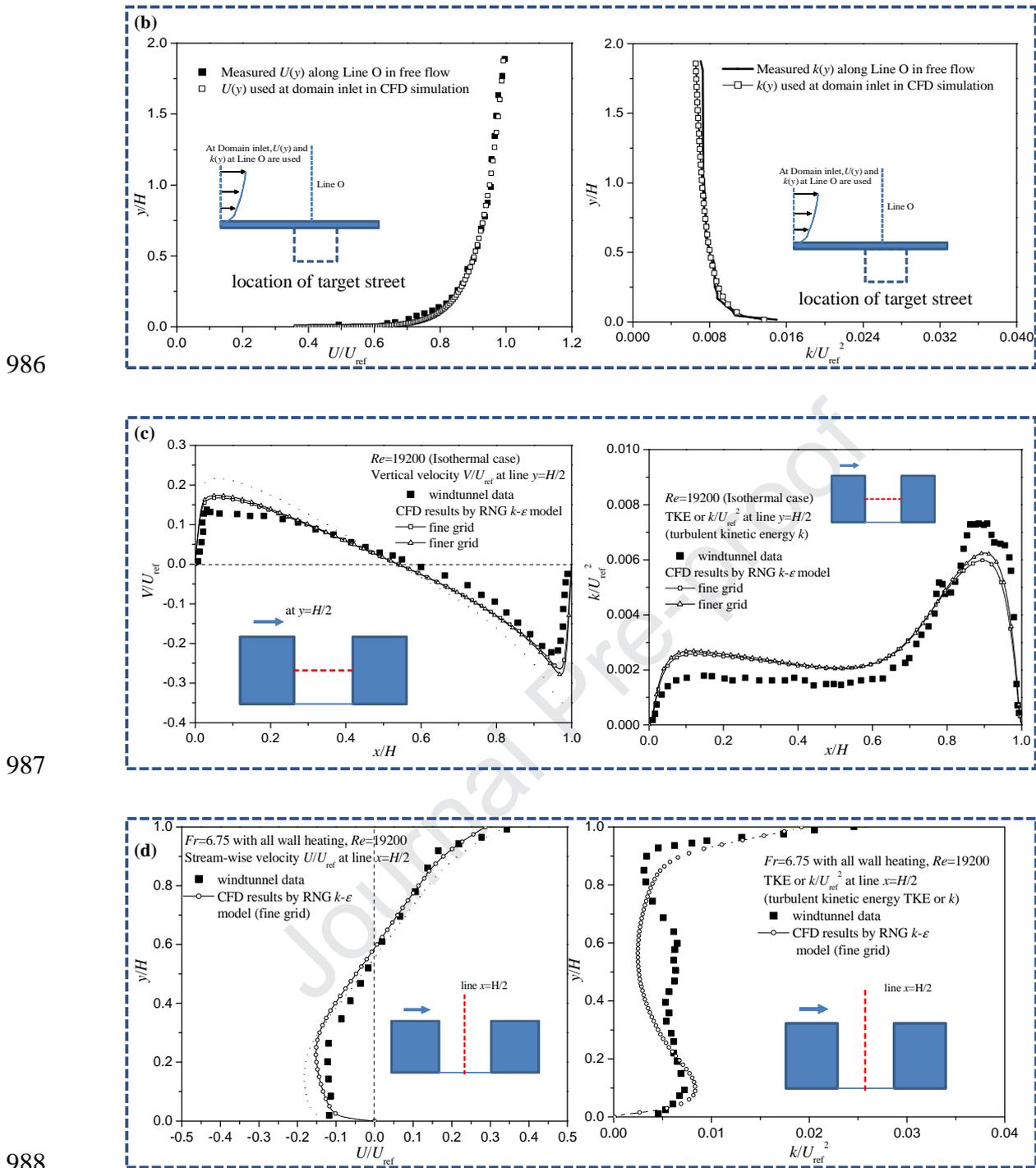
959 where U_{ref} and T_{ref} are the reference velocity and air temperature in the upstream free stream,

960 respectively ($T_{ref}=23^\circ\text{C}$, U_{ref} ranges from 2.32 to 0.68 m s^{-1}); T_w is the surface temperature of

961 the heated building wall or ground, ranging from 70 $^\circ\text{C}$ to 130 $^\circ\text{C}$.

962 In this CFD validation case, the computational domain has the same dimension as the
963 wind tunnel (Fig. A1a) and the CFD setup is similar to that described in subsection 2.2. A
964 fine grid with enhanced wall function (EWF) near wall surfaces is used to resolve the viscous
965 sub-layer, in which the order of magnitude of y^+ is 1 and the grid number is 58085 with a
966 minimum cell size of 1mm. To verify the grid independence, we also compared the results
967 with results from a finer grid arrangement with 190.016 cells and a minimum cell size of
968 0.5mm. Fig. A1b shows the vertical profiles of stream-wise velocity $U(z)$ and turbulent
969 kinetic energy (TKE) $k(z)$ measured in the free flow ($U_{ref}=1.45 \text{ m s}^{-1}$) of the wind tunnel.
970 They are used as the domain inlet boundary conditions in the CFD simulations.

971 The RNG $k-\varepsilon$ model is used with a predefined x -Component “wall shear stress” at
972 upstream and downstream points on the ground to minimize the stream-wise TKE decay and



989 Fig. A1. CFD validation study with reference to the literature [45-46]: (a) Wind tunnel model
 990 from Allegrini et al. [49] and CFD set-up. (b) Measured inlet profiles for the domain inlet
 991 boundary conditions in the CFD simulations. Validation profiles obtained from CFD
 992 simulations and wind tunnel data in (c) the isothermal case and (d) case with all wall heating.
 993 Here $U_{ref}=1.45 \text{ m s}^{-1}$ and $Re=19200$.

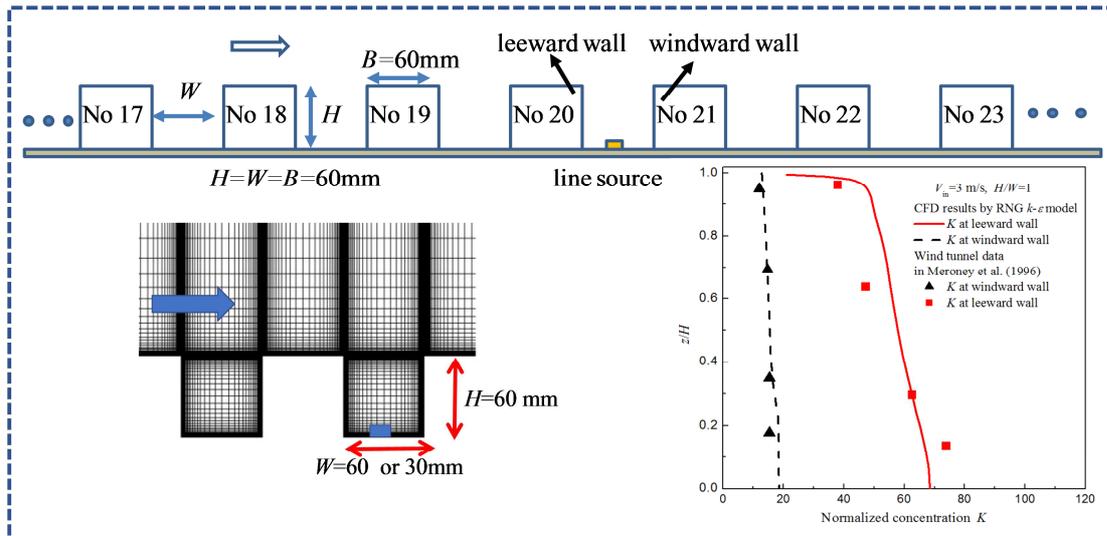
994

995 **Appendix A2. Validation of pollutant dispersion in a 2D street canyon ($AR=1$)**

996 Fig. A2 shows a sketch of the geometry and grid used for the validation of the pollutant
997 dispersion in street canyons. The CFD results are compared with the wind tunnel data from
998 Meroney et al. [22] which were performed in isothermal conditions. The experiments are
999 conducted with 28 parallel 2D street canyons of a uniform building height H
1000 ($H=W=B=60\text{mm}$), considering the street canyons that completely spanning the width of the
1001 tunnel and are perpendicular to the wind direction. There are 20 street canyons upstream of
1002 the target street canyon and 8 downstream. A steady line source (also lying entirely across the
1003 width of the wind tunnel) is located in the target street canyon. Measurements are taken of the
1004 vertical profiles of tracer gas (ethane) concentration along the windward and leeward wall
1005 surfaces. Here the concentration is presented in dimensionless form as $K = CUHL/Q$, where C
1006 is the measured ethane concentration, U is wind velocity measured in the free stream at
1007 0.50m above the tunnel floor, and L is line source length and Q is the source emission rate.

1008 In this CFD validation case, the 2D computational domain, the size of the street canyon
1009 and the boundary conditions are the same as in wind tunnel experiments. The total number of
1010 cells is 372.889 with a minimum grid size of 0.025mm at the wall surfaces (Fig. A2). To
1011 validate the numerical simulations, Fig. A2 shows vertical profiles of K at the leeward-side
1012 and windward-side walls of the target street canyon with $V_{in}=3\text{m s}^{-1}$ (at domain inlet). As
1013 expected K at the leeward wall is much higher than that at the windward wall; and K along
1014 the windward wall is almost constant, while that along the leeward wall decreases with
1015 increasing height. Overall, the results show that calculated K is in good agreement with the
1016 wind tunnel data, even though slightly over-estimated.

1017



1018

1019 Fig. A2. CFD set-up, grid arrangement and validation profiles of the normalized
 1020 concentration K along the windward wall and leeward wall evaluated using wind tunnel data
 1021 from Meroney et al. [22].

1022

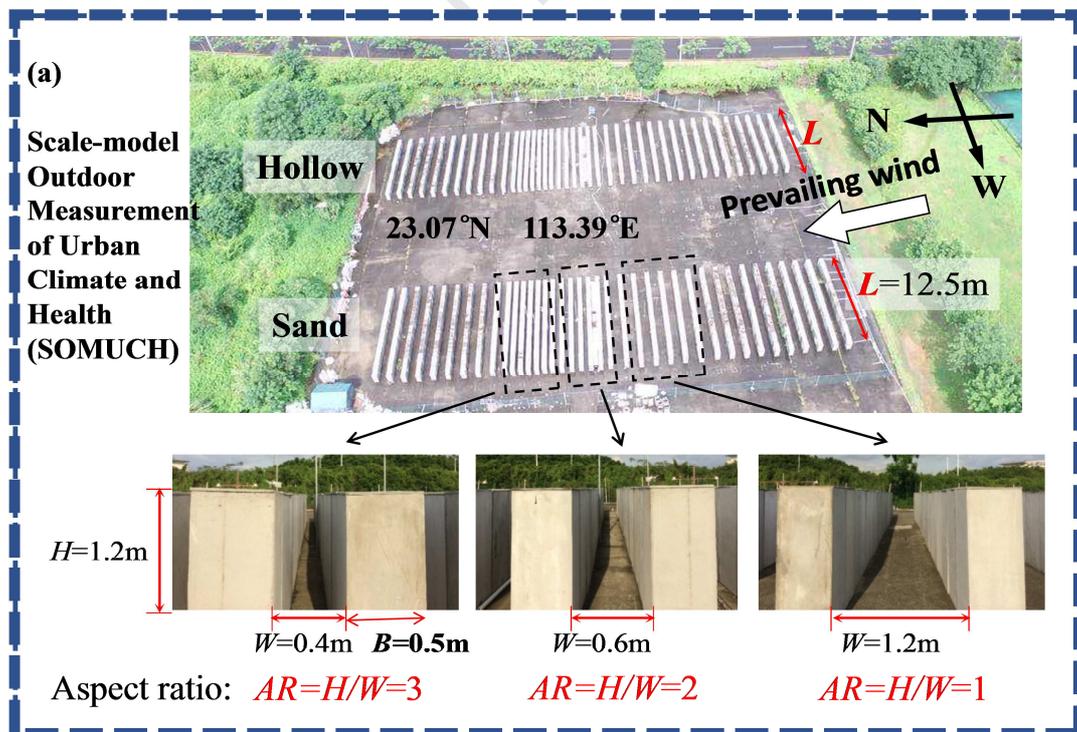
1023 Appendix A3. Flow pattern validation for a 2D deep street canyon by scaled outdoor 1024 experiments ($AR=2$ and 3 , $H=1.2m$)

1025 As displayed in Fig. A3a, Zhang et al. [9] carried out the scale-model outdoor field
 1026 experiments to study the flow patterns in a 2D street canyon with various street aspect ratios
 1027 (building height $H=1.2$ m; $AR=1,2,3$; street length $L=12.5m > 10H$). The velocity and
 1028 turbulence distribution, radiation fluxes, and the wall and air temperature in and above the
 1029 idealized street canyons were measured by 3D ultrasonic anemometers (Gill windmaster,
 1030 UK), four component radiometers (CRN4), thermal couples (K type) and temperature sensors
 1031 (iButton thermochron data logger).

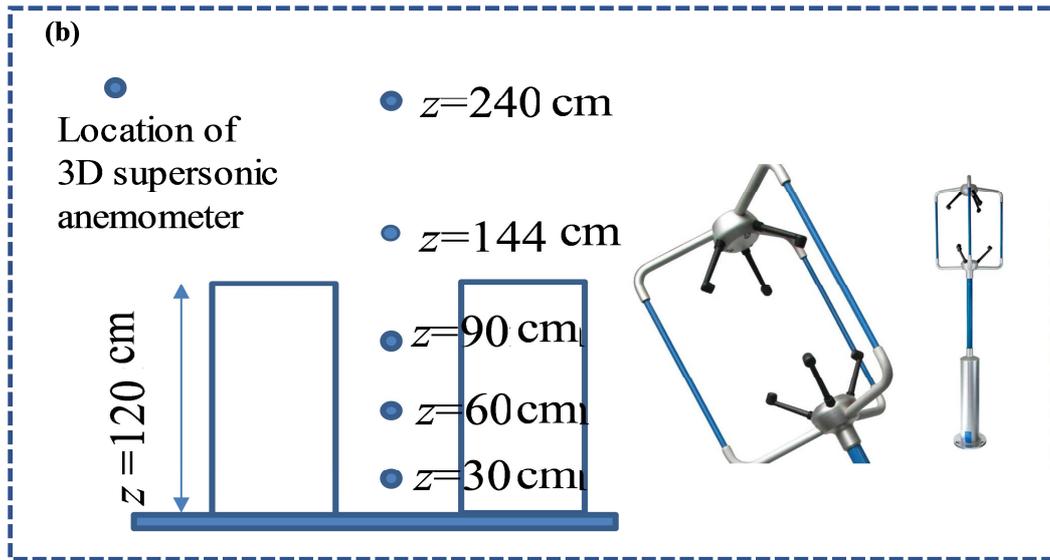
1032 For each type of streets canyon ($AR=1, 2, 3$), five 3D ultrasonic anemometers were used
 1033 to measure the temporal profiles of velocity components (U_x , U_y and U_z) and turbulence at
 1034 five heights ($z=0.3, 0.6, 0.9, 1.44, 2.4$ m)(Fig. A3b). The sampling rate of the ultrasonic

1035 anemometer was 20 Hz. Here, Fig. A3c and A3d only presents some examples of the
 1036 experimental profiles of stream-wise velocity (U_x , i.e. perpendicular to the street axis) in
 1037 street canyon with $AR=3$ when the Reynolds number is large ($Re \sim 1.5 \times 10^5 \gg 11000$ as $U_{ref} \sim 2.0$
 1038 m s^{-1}) and the buoyancy force is relatively weak (i.e. Froude number $F_r = \frac{U_{ref}^2}{gH(\Delta T / T_{ref})}$
 1039 ~ 10.2 as $\Delta T = 10 \text{ K}$ and $U_{ref} = 2.0 \text{ m s}^{-1}$). Thus, the wind-driven dynamic force dominates
 1040 urban airflows and the Reynolds number independence requirement is fully satisfied.

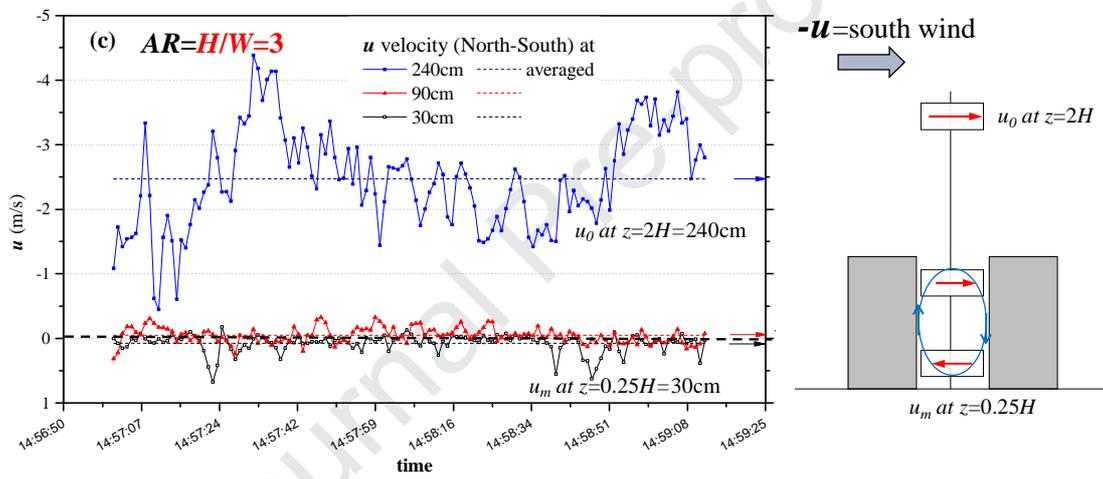
1041 Obviously, Fig. A3c and A3d show that, regardless of the aspect ratio being 3 or 2, the
 1042 stream-wise velocities at $z=0.25H$ in the field measurements are positive while those at
 1043 $z=0.75H$ and $z=2H$ are negative, confirming that there is only one main vortex in such 2D
 1044 deep street canyons ($AR=2$ and 3). This is consistent with the flow patterns of the CFD results
 1045 in subsection 3.1. The more detailed experimental setups can be found in Zhang et al. [9].



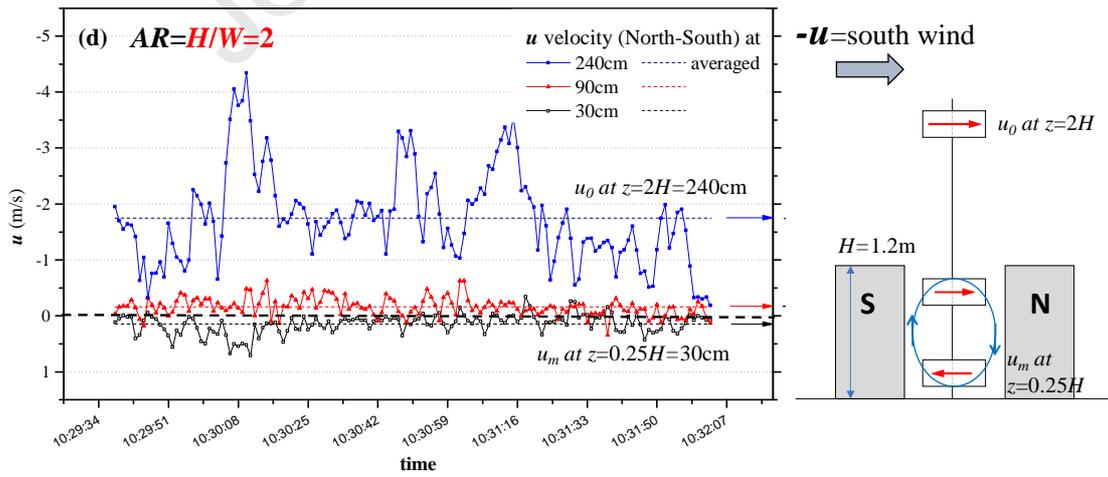
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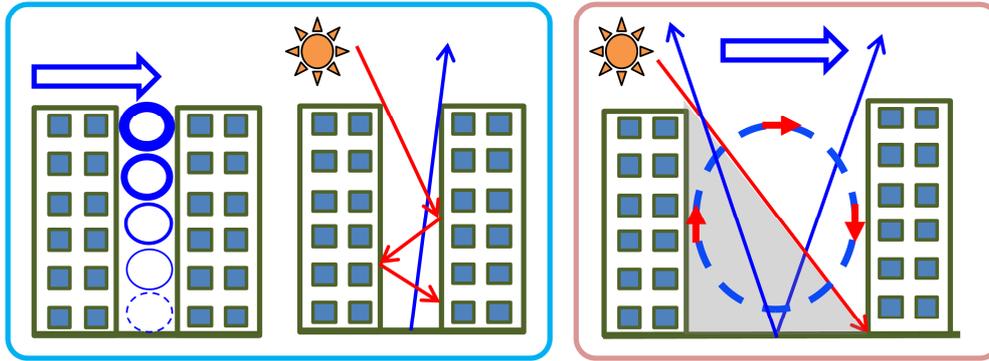


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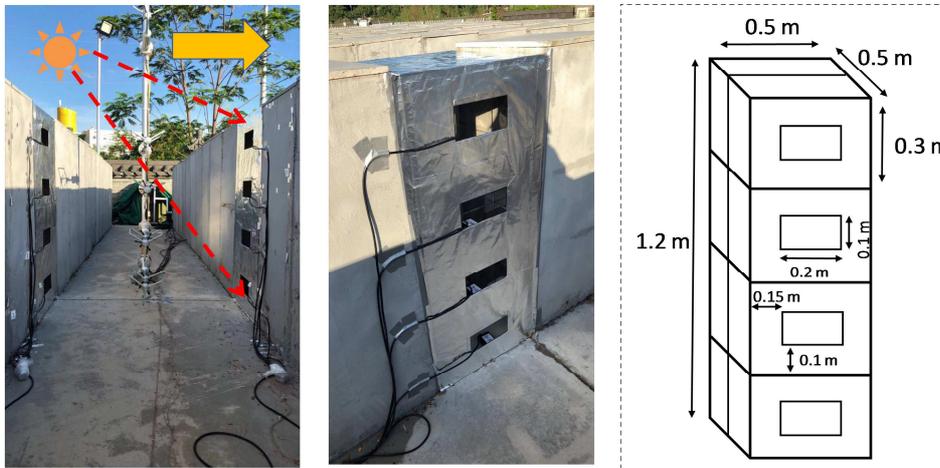


1049

(e) Ventilation and turbulence, solar shading and trapping



Scaled outdoor experiments by coupling turbulence and radiation, or integrating indoor and outdoor environment



1050

1051 Fig. A3. (a) View of the scale-model outdoor experiment on street canyon models with $AR=1$,

1052 2 and 3. (b) Schematic of the 3D ultrasonic anemometer locations. Example profiles of the

1053 stream-wise velocity (U_x , m s^{-1}) in a street canyon with (c) $AR=3$ and (d) $AR=2$. (e) Future

1054 studies of coupling urban turbulence and radiation processes, or integrating indoor and

1055 outdoor.

1056

1057 **Appendix 4 Scaled CFD flow validation estimated by the wind tunnel data ($AR=2.4$,**

1058 **$H=1.2\text{m}$)**

1059 To further evaluate numerical accuracy and grid independence of CFD simulations, we

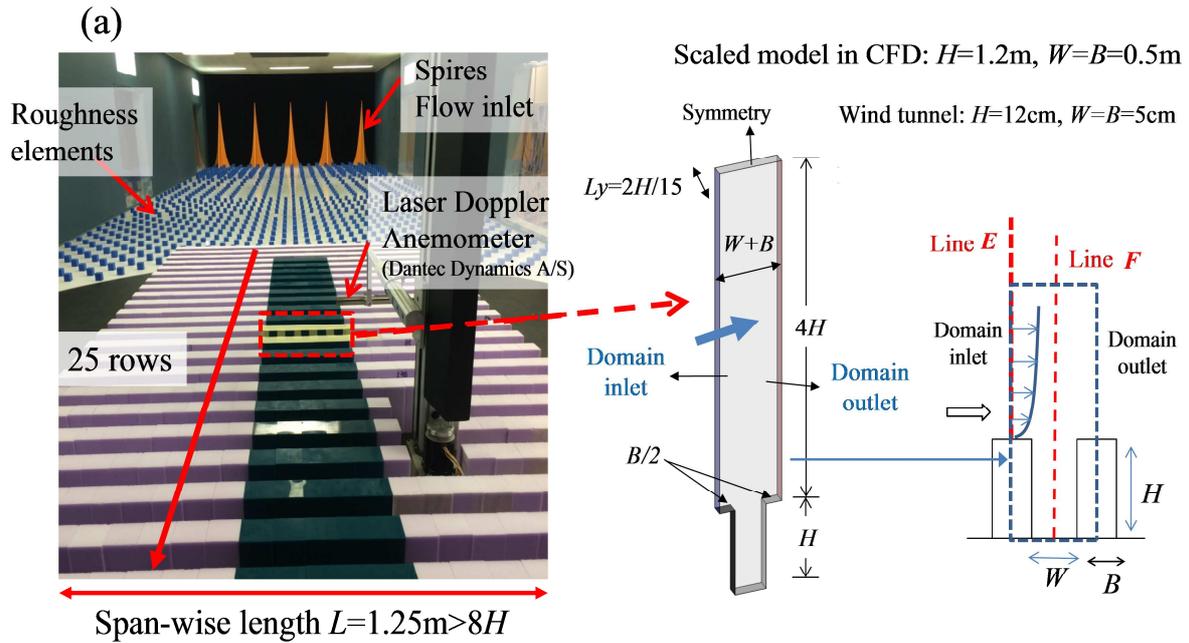
1060 carried out wind tunnel experiments in University of Gavle, Sweden in 2017 and 2018 (Fig.

1061 A4a) with the working section of 11m long, 3m wide, 1.5m tall. There are 25 rows of
1062 building models and 24 street canyons from upstream toward downstream with a
1063 perpendicular approaching wind to street axis. The key parameters of wind-tunnel-scale
1064 street canyon models (Fig. A4a) include building height ($H=12\text{cm}$), building width ($B=5\text{cm}$),
1065 street width ($W=5\text{cm}$), i.e. street aspect ratio is $AR=2.4$. The span-wise length is
1066 $L=1.25\text{m}>10H$ which ensure the 2D flow characteristics in street canyons. The measured
1067 vertical profiles along street centreline in the 12th and 13th street canyon are almost the same
1068 (not shown here), verifying that the flow in the 12th street canyon is fully-developed. The
1069 background wind speed at the boundary-layer height in far upstream free flow is $U_{ref}=13\text{ m}$
1070 s^{-1} , attaining the reference Re numbers of 10^5 ($Re=U_{ref}H/\nu$). Stream-wise (U_x) and vertical
1071 (U_z) velocity components along the street centreline of Line F in the 12th street canyon are
1072 measured by Laser Doppler Anemometry (LDA) System (Fig. A4a). The measured vertical
1073 profiles of stream-wise velocity (U_x) and turbulent kinetic energy (k) along the centreline
1074 above building roof center (Line E) are displayed in Fig. A4b which will be adopted to
1075 provide boundary condition at the domain inlet of CFD simulations.

1076 In the CFD validation case, the scaled street canyon models ($H=1.2\text{m}$, $W=B=0.5\text{m}$) are
1077 investigated with the scale ratio of 10:1 to the wind-tunnel-scale models ($H=12\text{cm}$). Ansys
1078 Fluent with the RNG k - ϵ model is used to perform CFD simulations. The domain inlet
1079 boundary condition is provided by the vertical profiles of stream-wise velocity and turbulent
1080 quantities measured at Line E (Fig. A4b) with a spatial scale ratio of 10:1. To perform a grid
1081 independence study, two kinds of grid arrangements are tested with the minimum grid sizes
1082 of 0.5mm (fine grid) and 1mm (medium grid) in which grid numbers are 1383668 and
1083 807024 respectively (Fig. A4c).

1084 Then Fig. A4d compares CFD results and wind tunnel data by the stream-wise velocity
1085 (U_x) profiles at Line F. Obviously, the predicted wind profiles with the fine and medium

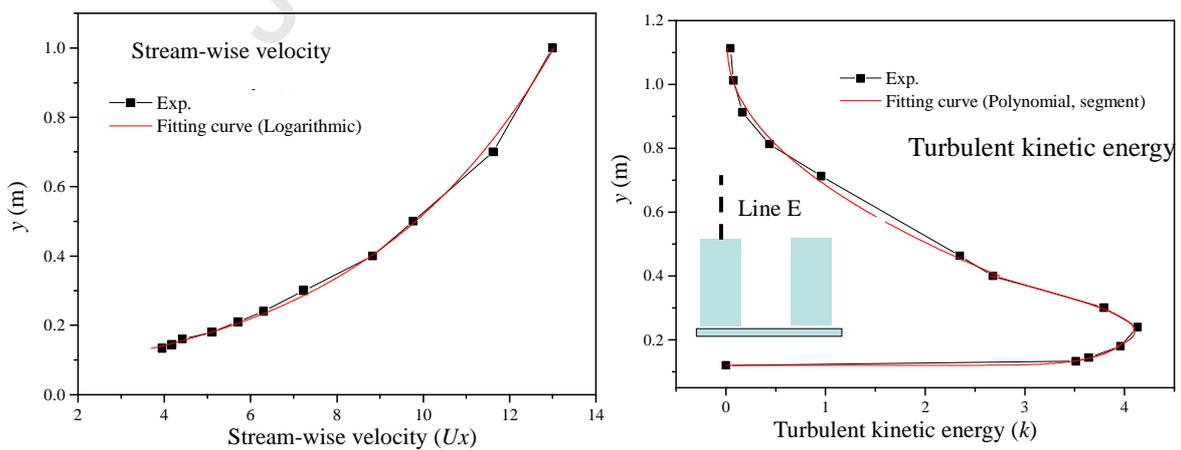
1086 grids are nearly the same and both agree well with wind tunnel data. As a result, the RNG $k-\epsilon$
 1087 model is reliable in simulating flow in 2D idealized street canyons with $AR=2.4$ ($Re \gg 11000$)
 1088 and the medium grid arrangement is recommended in the case studies. The CFD validation
 1089 study also confirms that there is only one main vortex as $AR=2.4$ (Fig. A4d) if the Re-
 1090 number-independence requirement is satisfied.



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 1092

(a)

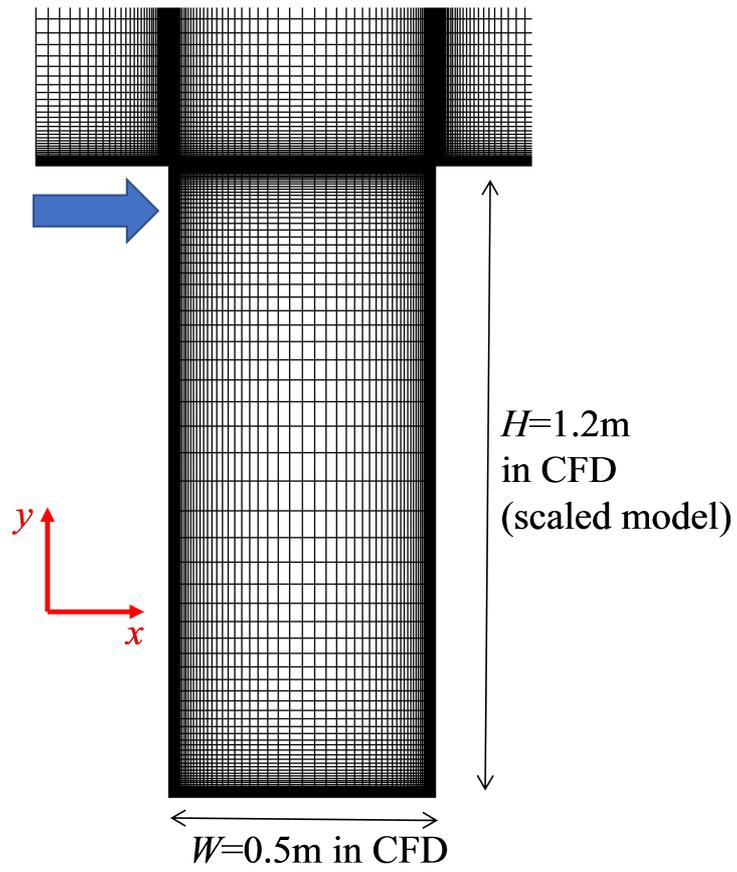
(b) U_x and k along Line E above building roof ($y > H=0.12\text{m}$) in wind tunnel experiment



1093
 1094

(b)

(c) Medium grid arrangement with minimum grid size of 1mm

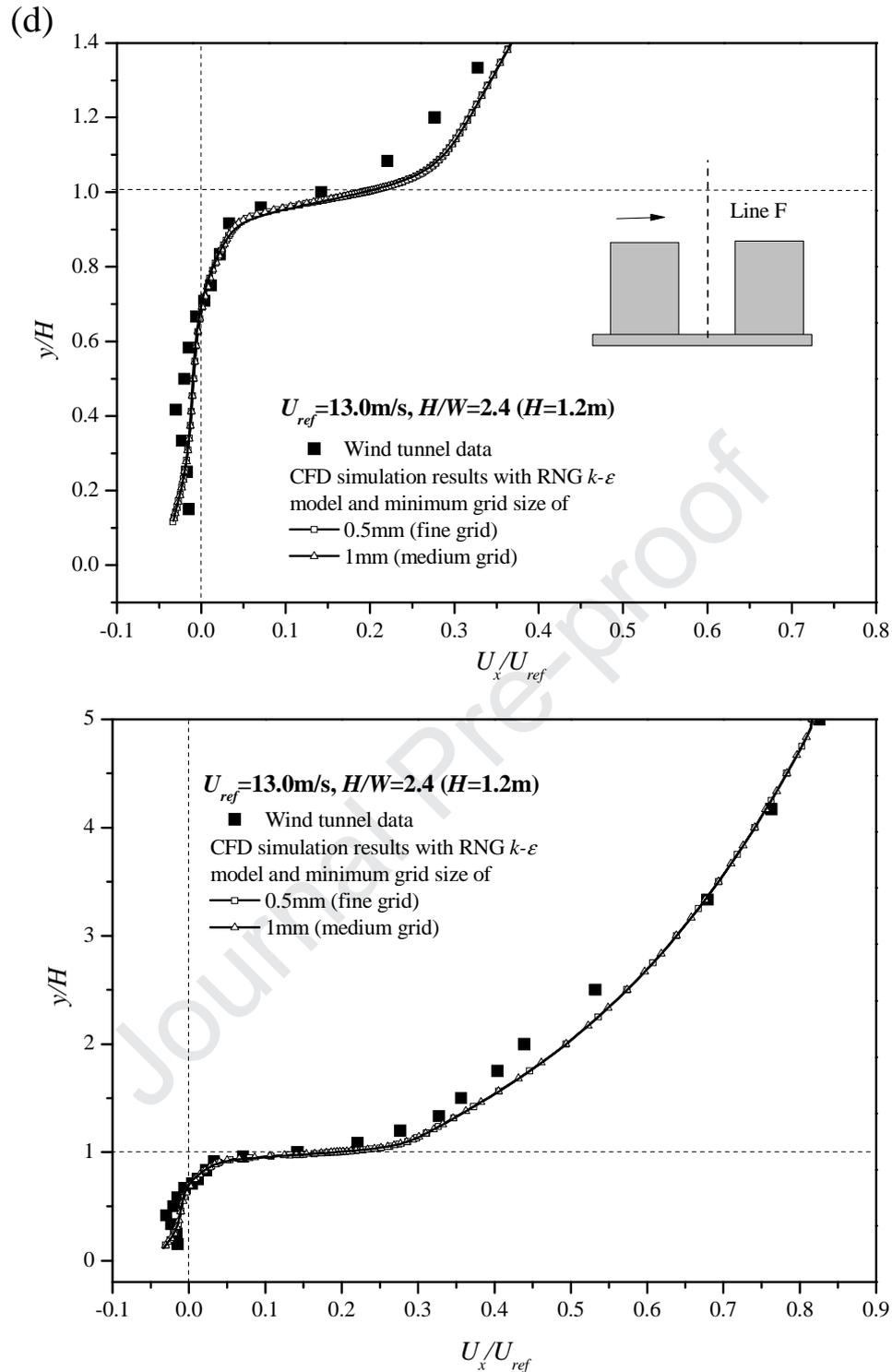


1095

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(c)



1098

1099 Fig. A4. (a) Wind tunnel experiments in 2D street canyon with ($AR=2.4, H=12\text{cm}, U_{ref}=13\text{ m}$
 1100 $\text{s}^{-1}, W=B=5\text{cm}, Re\sim 10^5$) and the CFD setups in scaled model ($AR=2.4, H=1.2\text{m}, U_{ref}=13\text{m s}^{-1},$
 1101 $W=B=0.5\text{m}, Re\sim 10^6$); (b) The measured vertical profiles of the stream-wise velocity (U_x) and
 1102 the turbulent kinetic energy (k) along the centerline above building roof center (Line E).

1103 (c)The medium grid arrangements in the scaled CFD simulations. (d) The CFD validation and
1104 grid-independence study in the scaled CFD simulation with the stream-wise velocity profiles
1105 along the Line F ($H=1.2\text{m}$, $Re\sim 10^6$).

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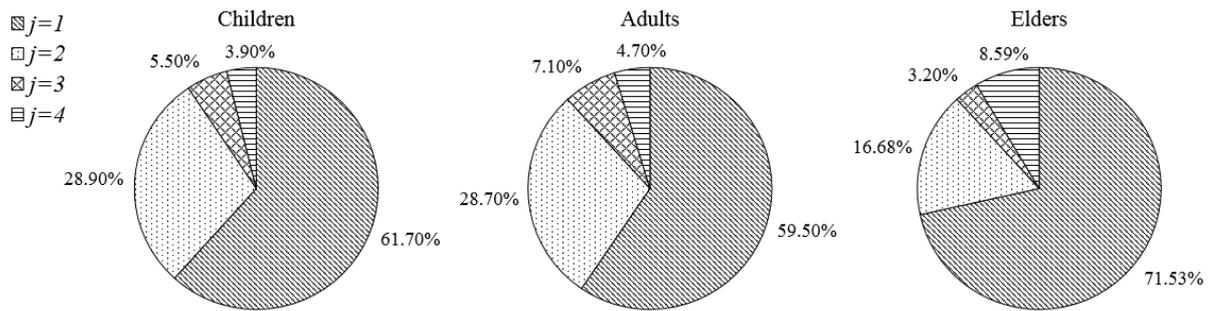


Figure 1. Breathing rate and time patterns for various age groups and microenvironments [64-65].

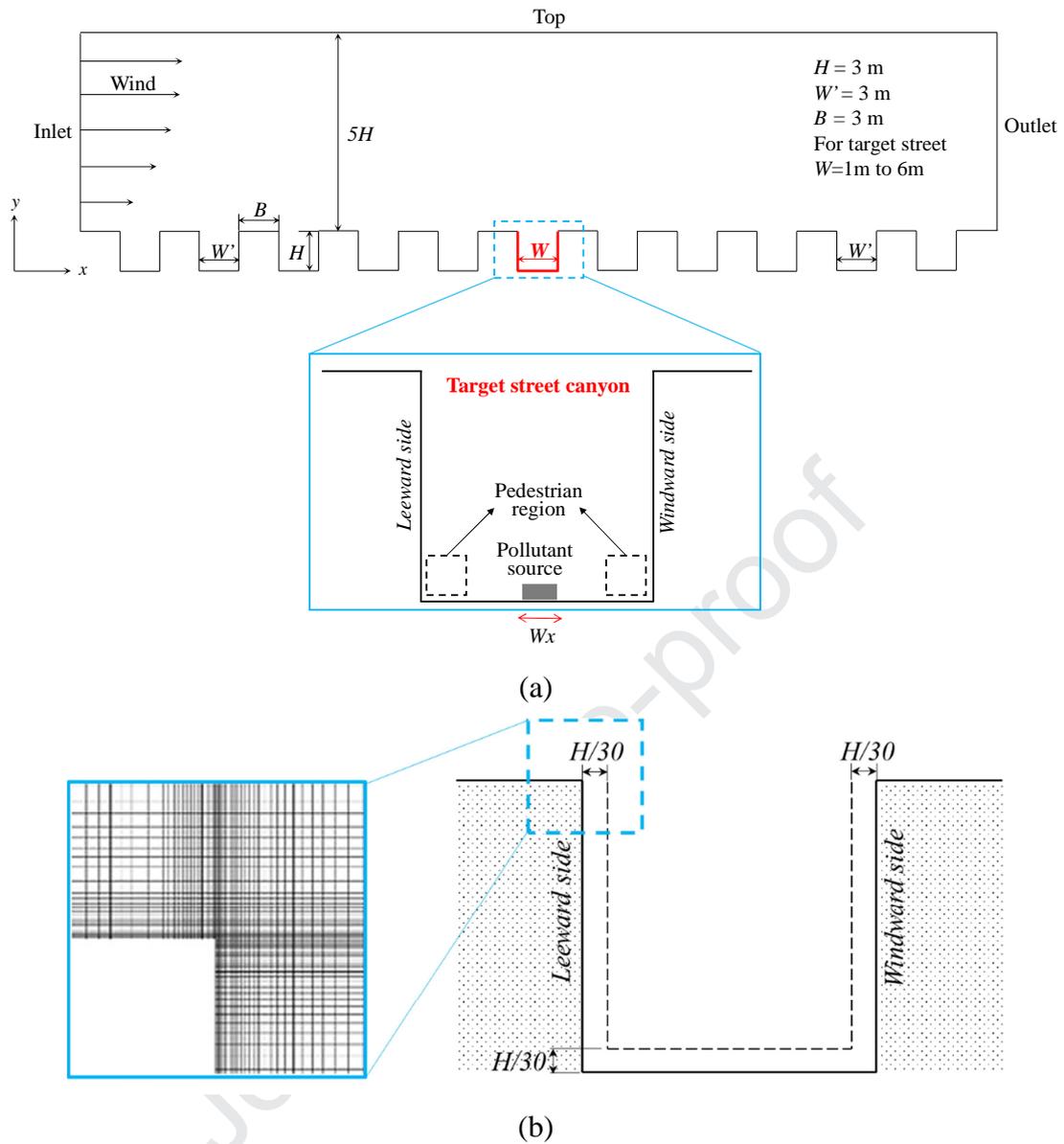
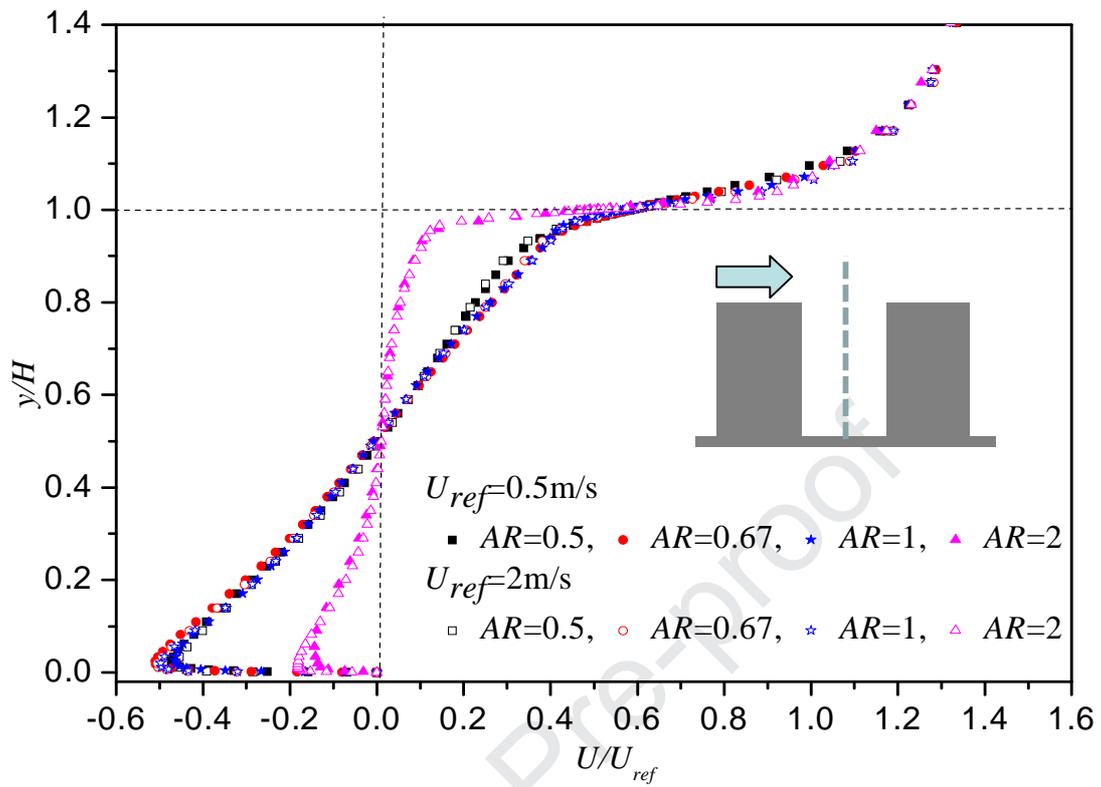
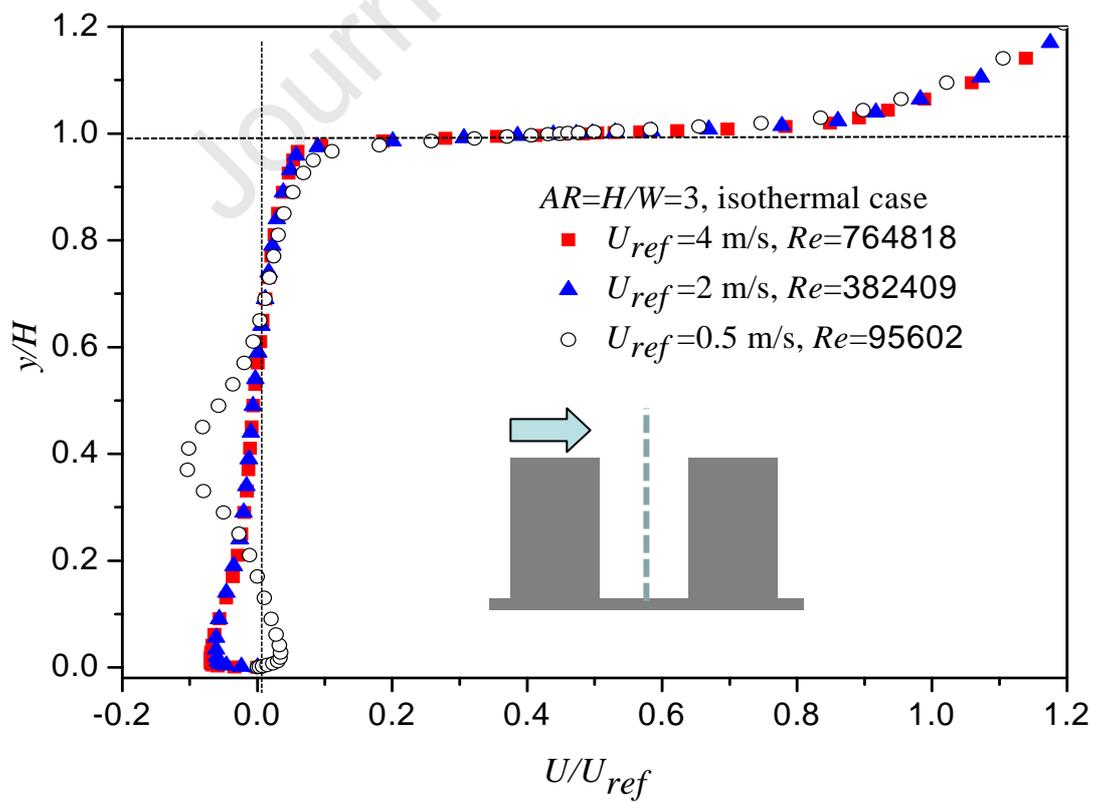


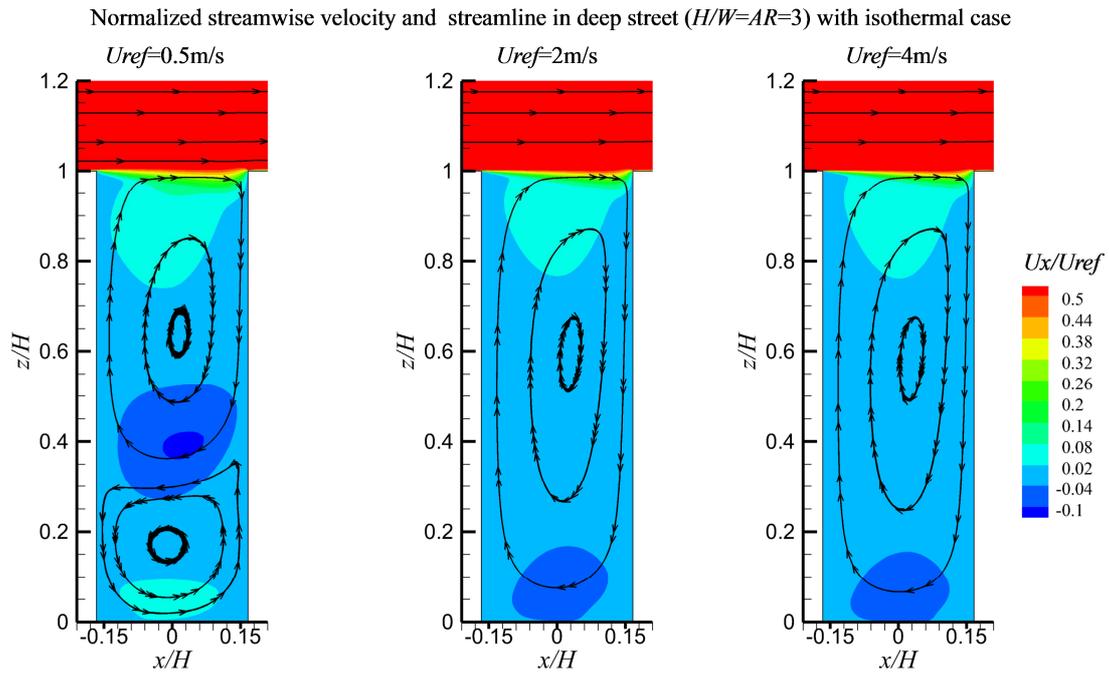
Figure 2. (a) Dimensions of the simulated street canyon model in CFD. (b) The grid arrangement of 2D CFD simulations.



(a)



(b)



(c)

Figure 3 Normalized stream-wise velocity along the street centerline in isothermal cases with (a) background wind speed of $U_{ref} = 0.5 \text{ m s}^{-1}$ and 2 m s^{-1} ($Re = 95602$ and 382409) as $AR = H/W = 0.5, 0.67, 1, \text{ and } 2$, (b) $U_{ref} = 0.5 \text{ m s}^{-1}, 2 \text{ m s}^{-1}, 4 \text{ m s}^{-1}$ where $AR = 3$ ($Re = 95602, 382409, 764818$). (c) Normalized streamwise velocity and streamline in deep street ($H/W = AR = 3$) in isothermal case.

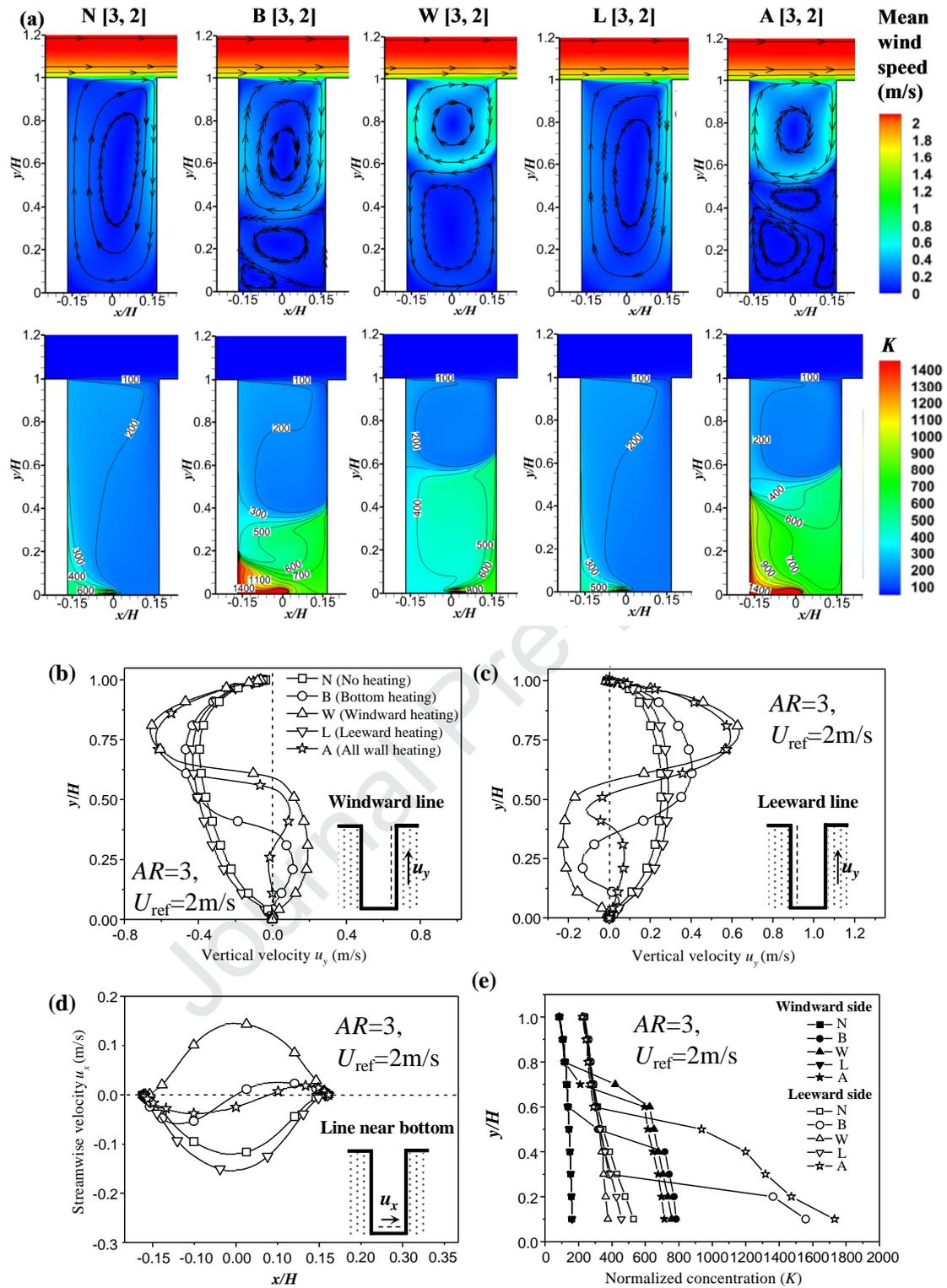


Figure 4. In cases where $AR=3$, $U_{ref}=2.0 \text{ m/s}$: (a) Contour of the mean wind speed (m s^{-1}) and normalized concentration K , Vertical velocity u_y along (b) the windward line and (c) the leeward line, (d) streamwise velocity u_x along the bottom line. (e) Spatial average $\langle K \rangle$ along windward wall and leeward wall.

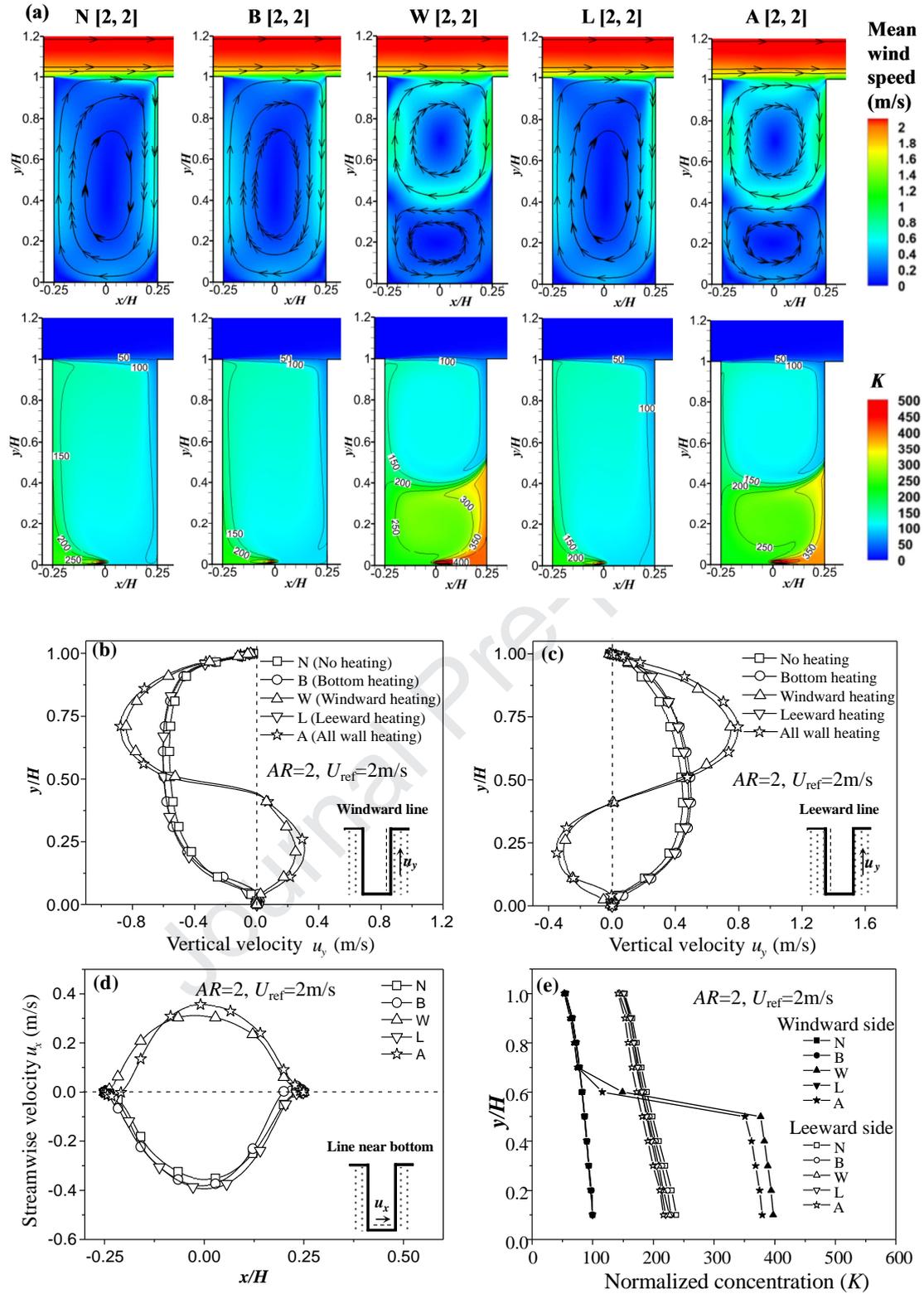
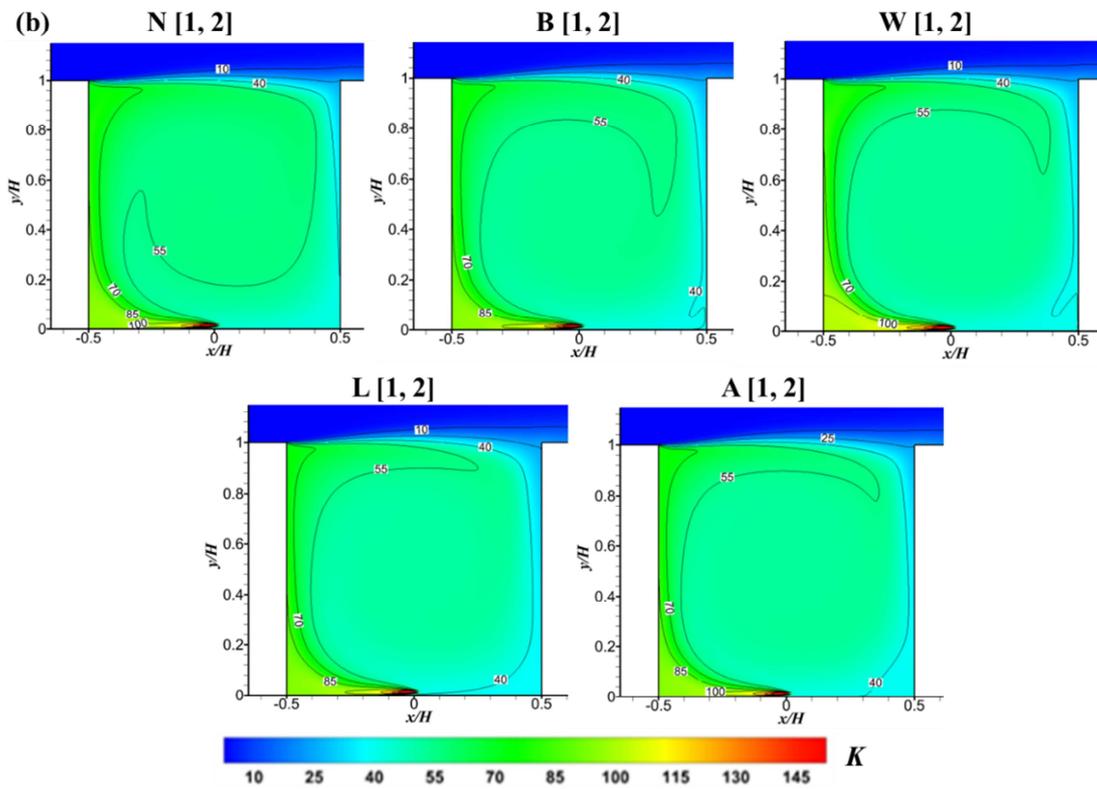
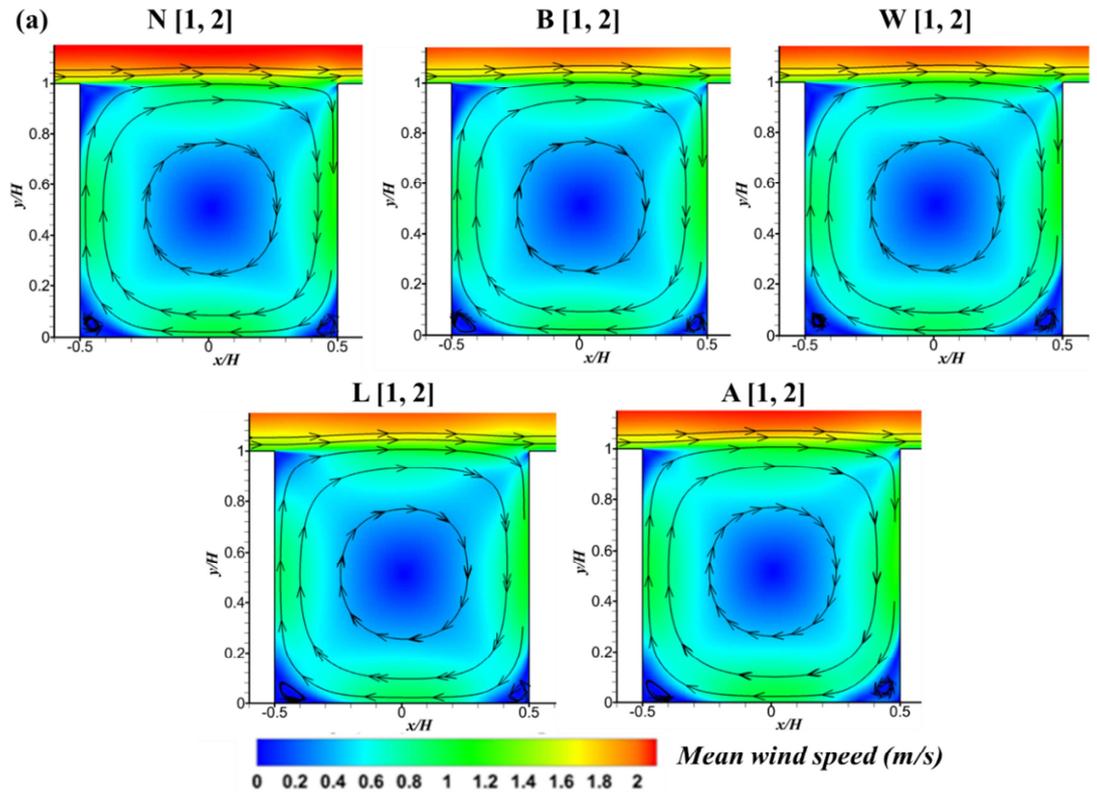


Figure 5. In cases where $AR=2$, $U_{ref}=2.0\text{ m/s}$: (a) Contour of the mean wind speed (m s^{-1}) and K ; u_y along (b) the windward line and (c) the leeward line; (d) u_x along the bottom line; (e) $\langle K \rangle$ at the windward and leeward walls.



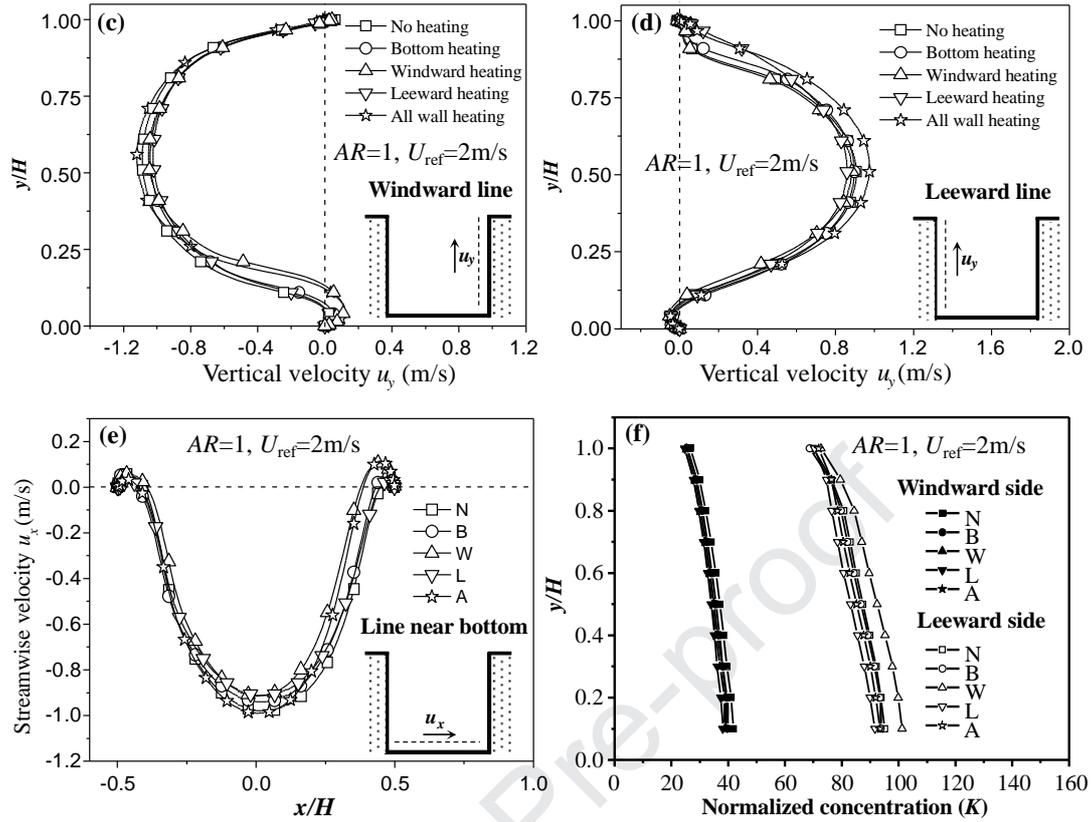


Figure 6. In cases where $AR=1$ and $U_{ref}=2\text{ m s}^{-1}$: Contour of (a) the mean wind speed (m s^{-1}) and K , u_y along (c) the windward line and (d) the leeward line; (e) u_x along the bottom line; (f) $\langle K \rangle$ at the windward and leeward walls.

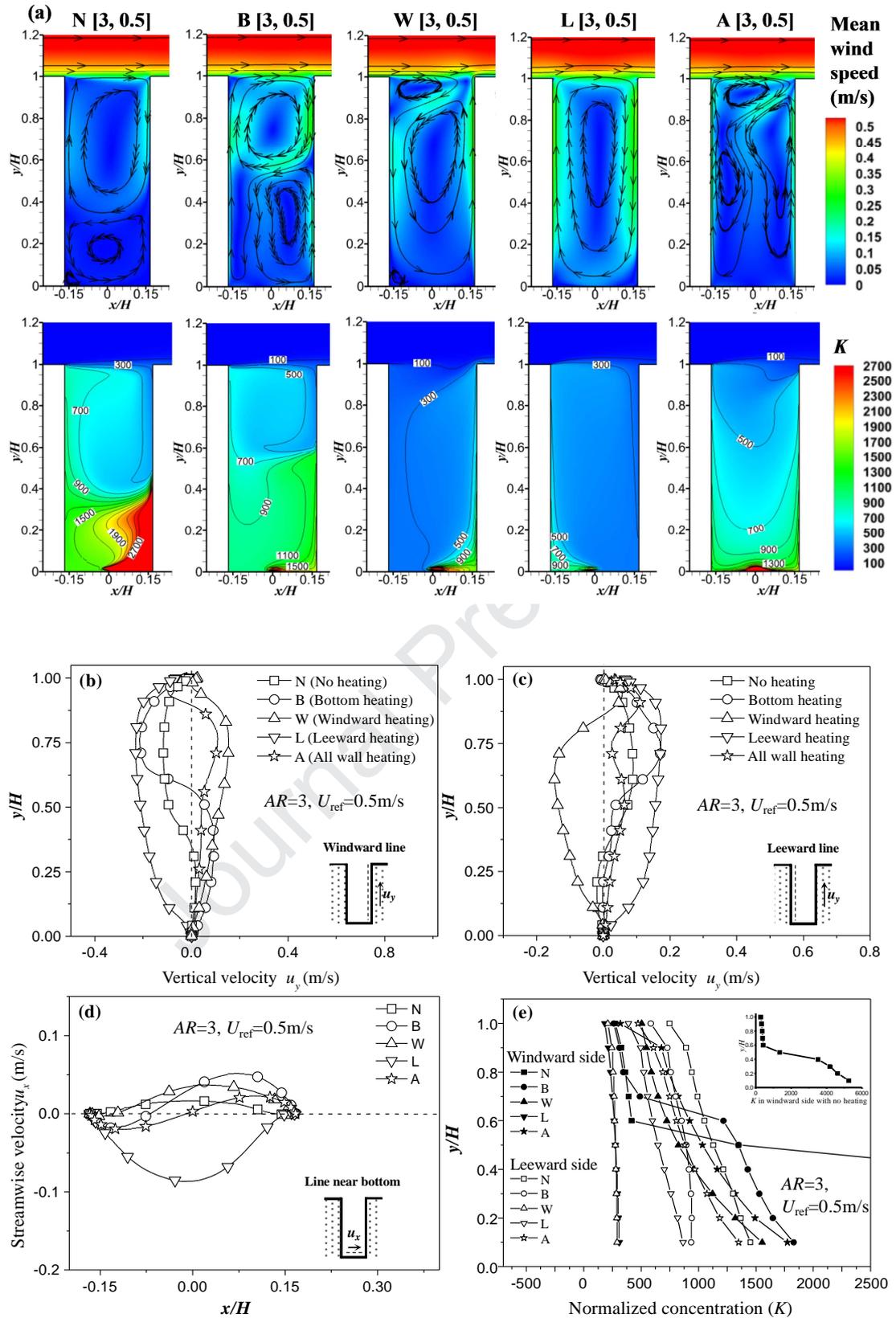


Figure 7. In cases where $AR=3$ and $U_{ref}=0.5 \text{ m s}^{-1}$: (a) Contour of the mean wind speed (m s^{-1}) and K . u_y along (b) the windward line and (c) the leeward line; (d) u_x along the bottom line; (e) $\langle K \rangle$ at the windward and leeward walls.

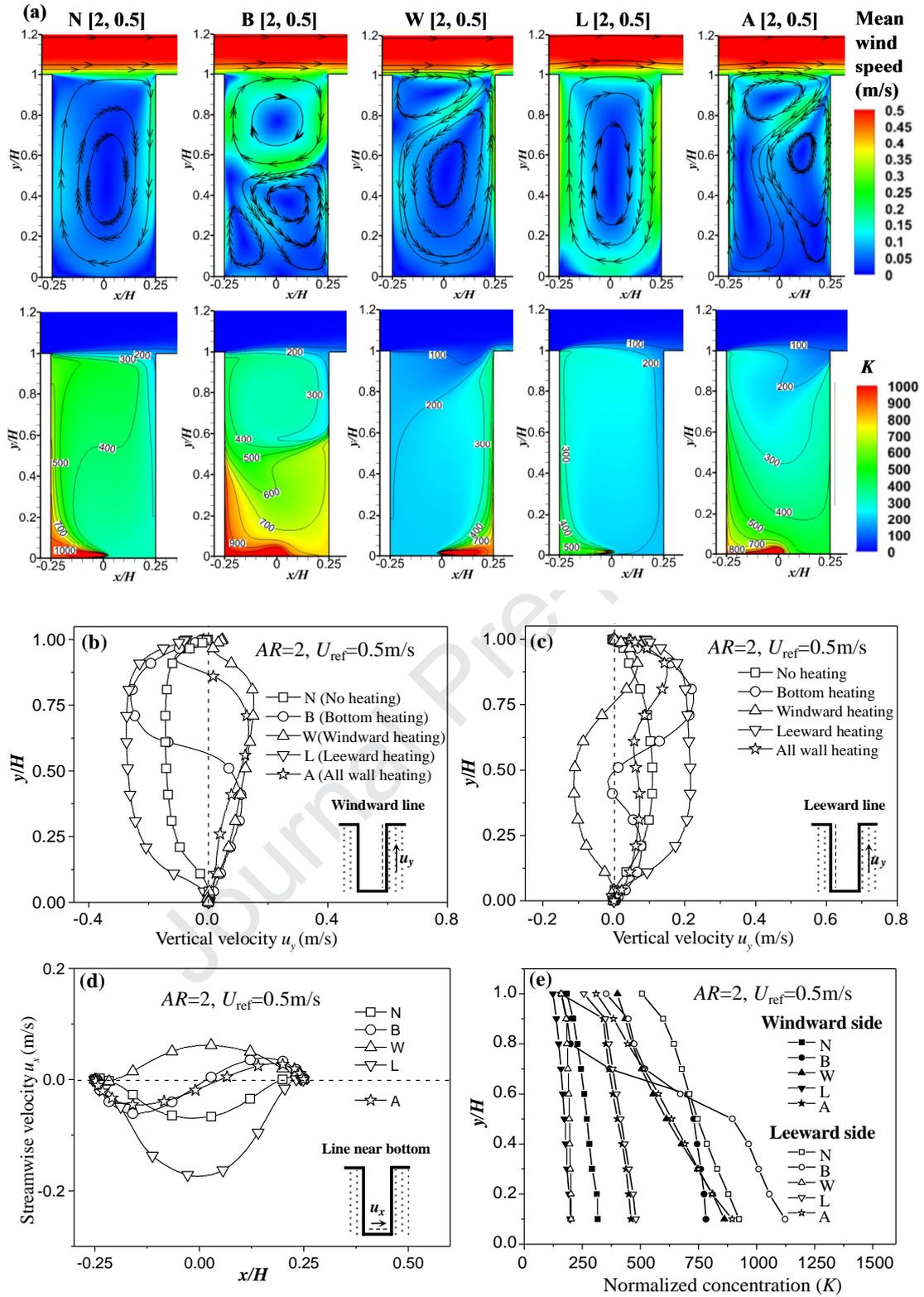
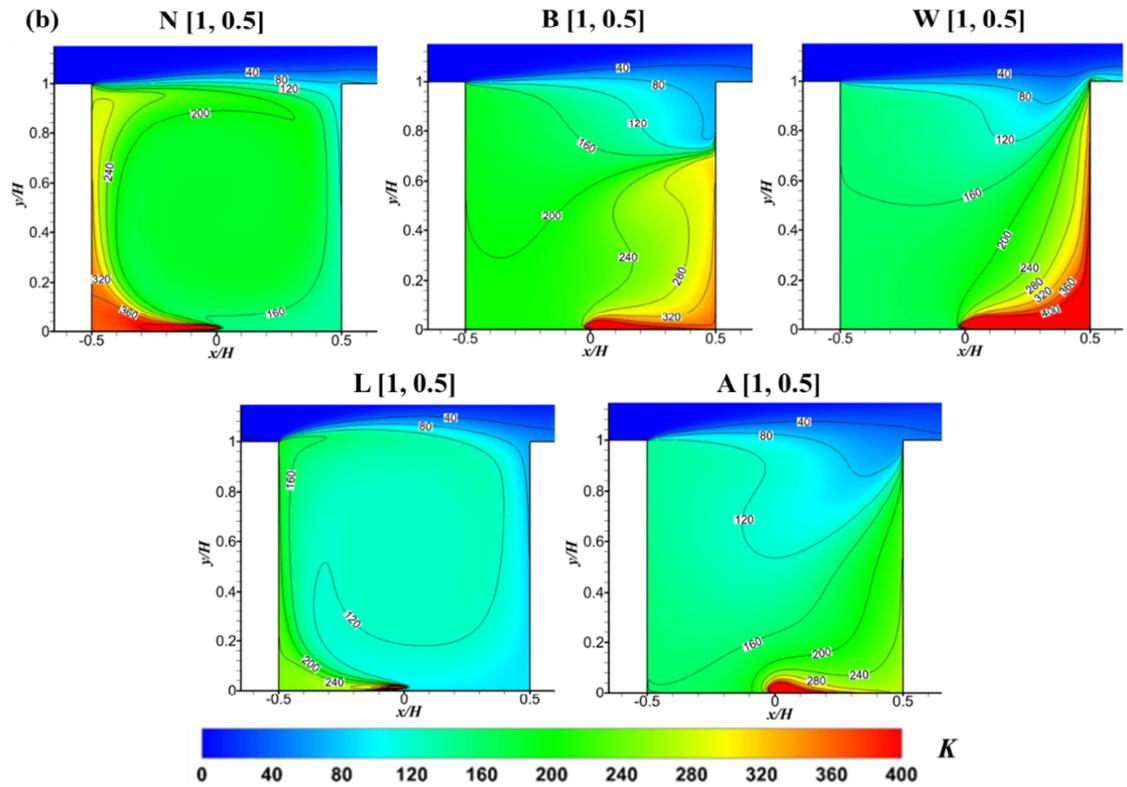
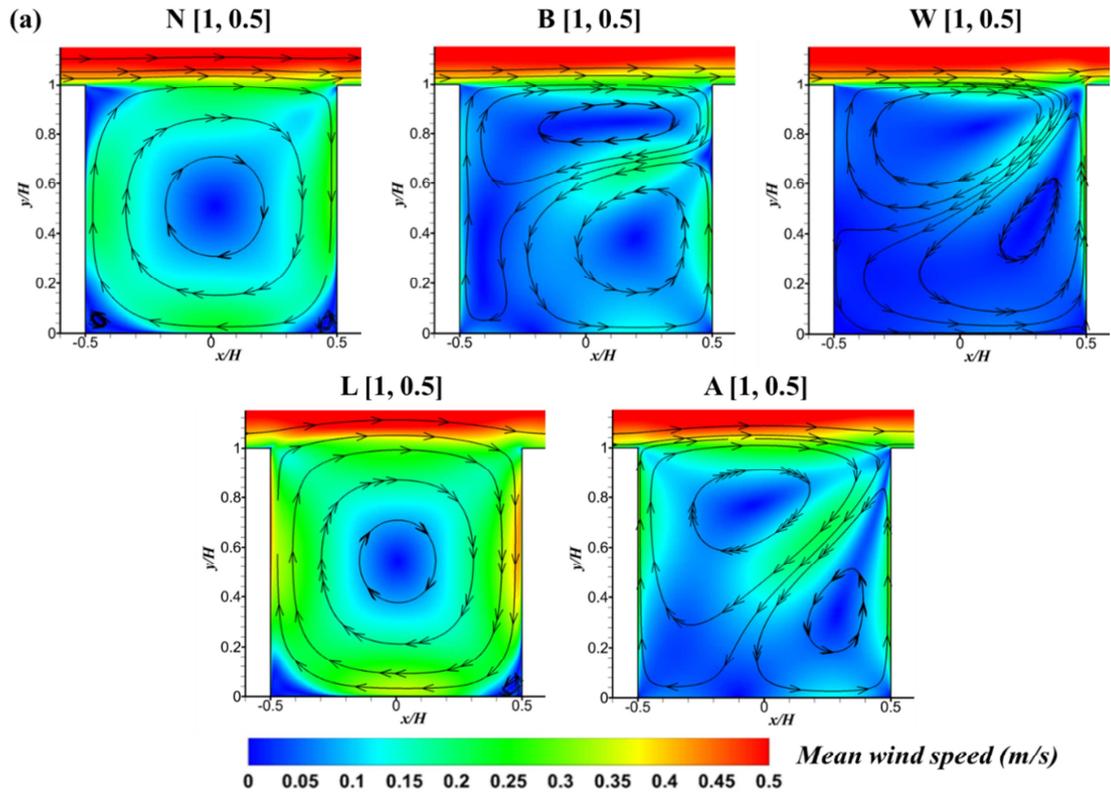


Figure 8. In cases where $AR=2$ and $U_{ref}=0.5 \text{ m s}^{-1}$: (a) Contour maps of the mean wind speed (m s^{-1}) and K . u_y along (b) the windward line and (c) the leeward line; (d) u_x along the bottom line; (e) $\langle K \rangle$ at the windward and leeward walls.



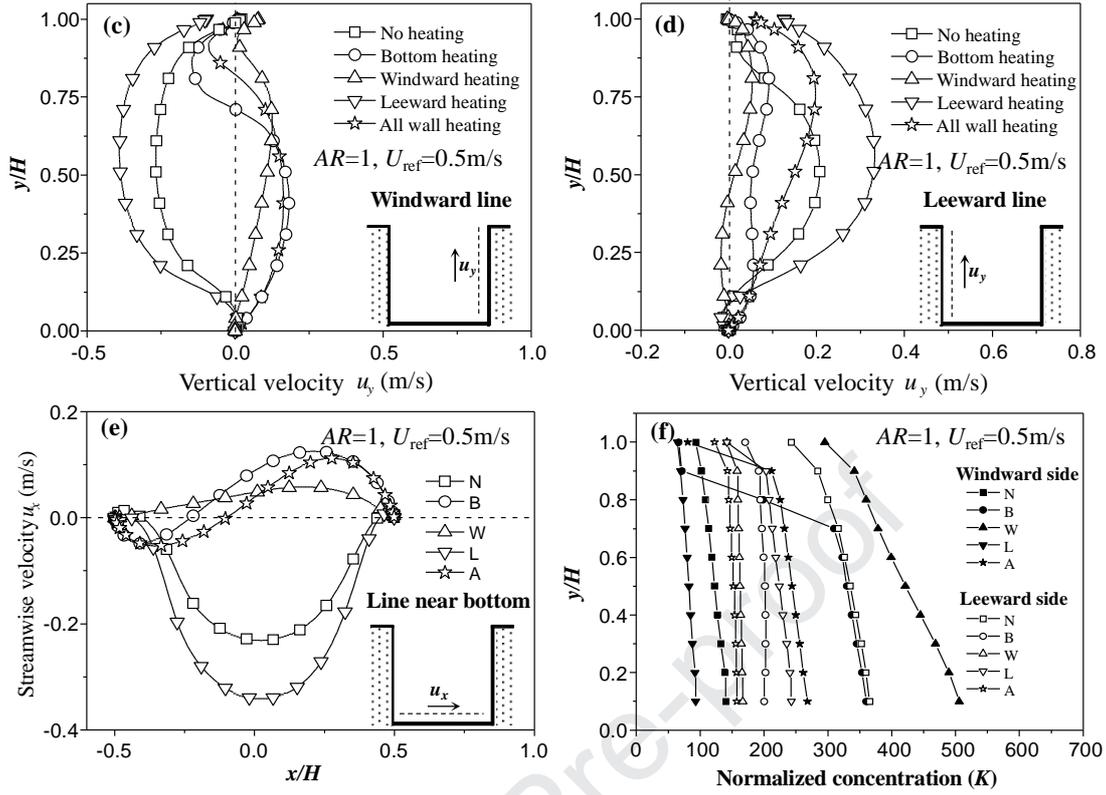
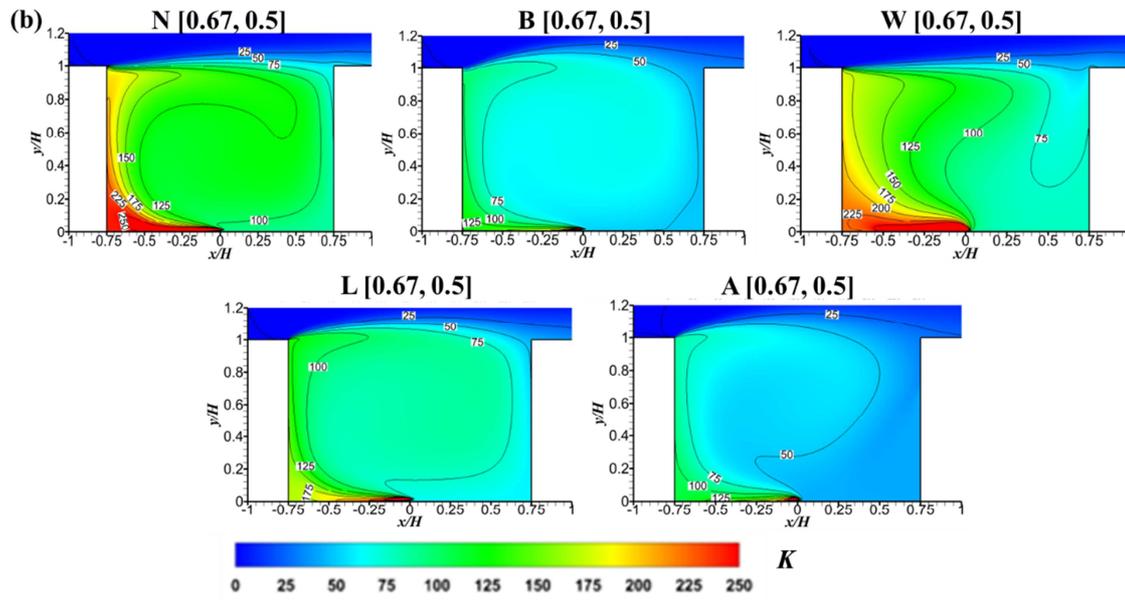
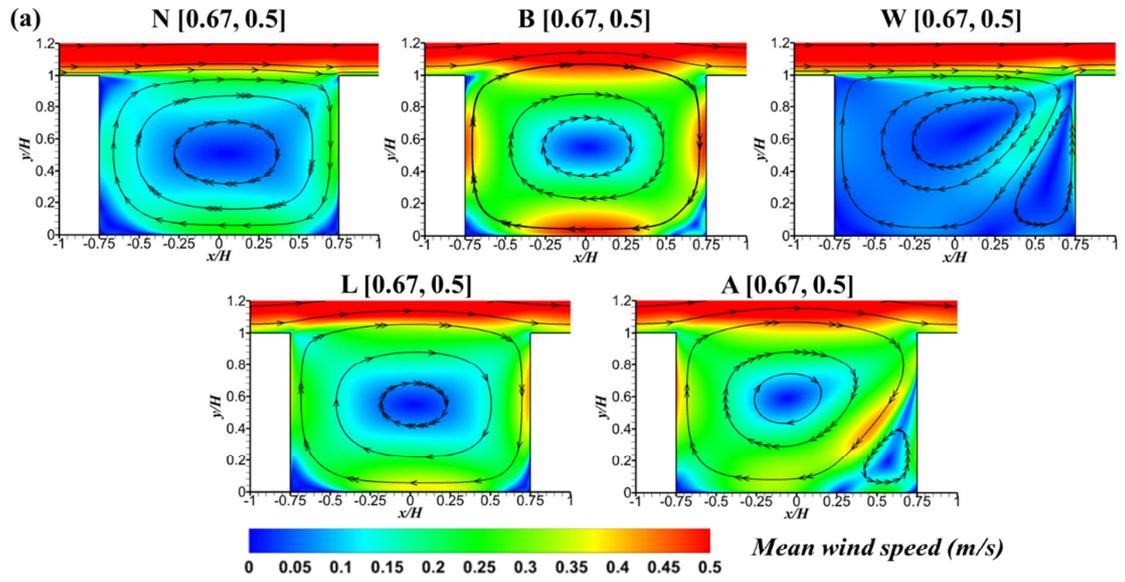


Figure 9. In cases where $AR=1$ and $U_{ref}=0.5 \text{ m s}^{-1}$: Contour of (a) the mean wind speed (m s^{-1}) and (b) K . u_y along (c) the windward line and (d) the leeward line; (e) u_x along the bottom line; (f) $\langle K \rangle$ at the windward and leeward walls.



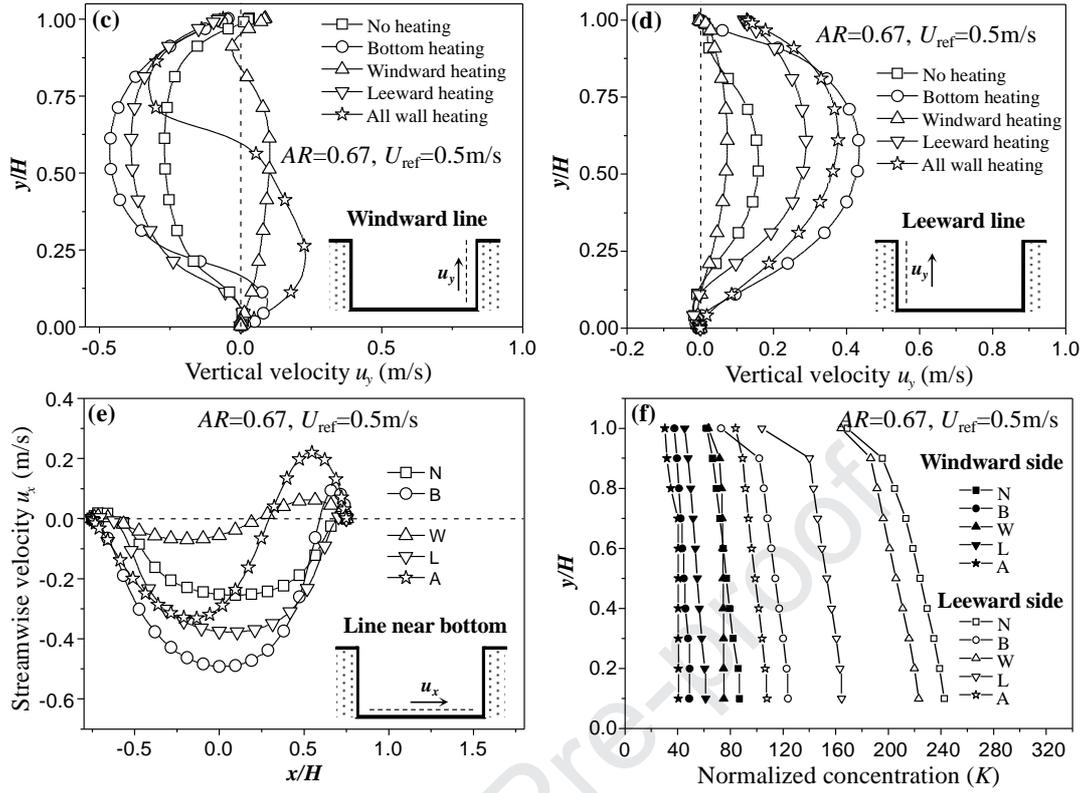
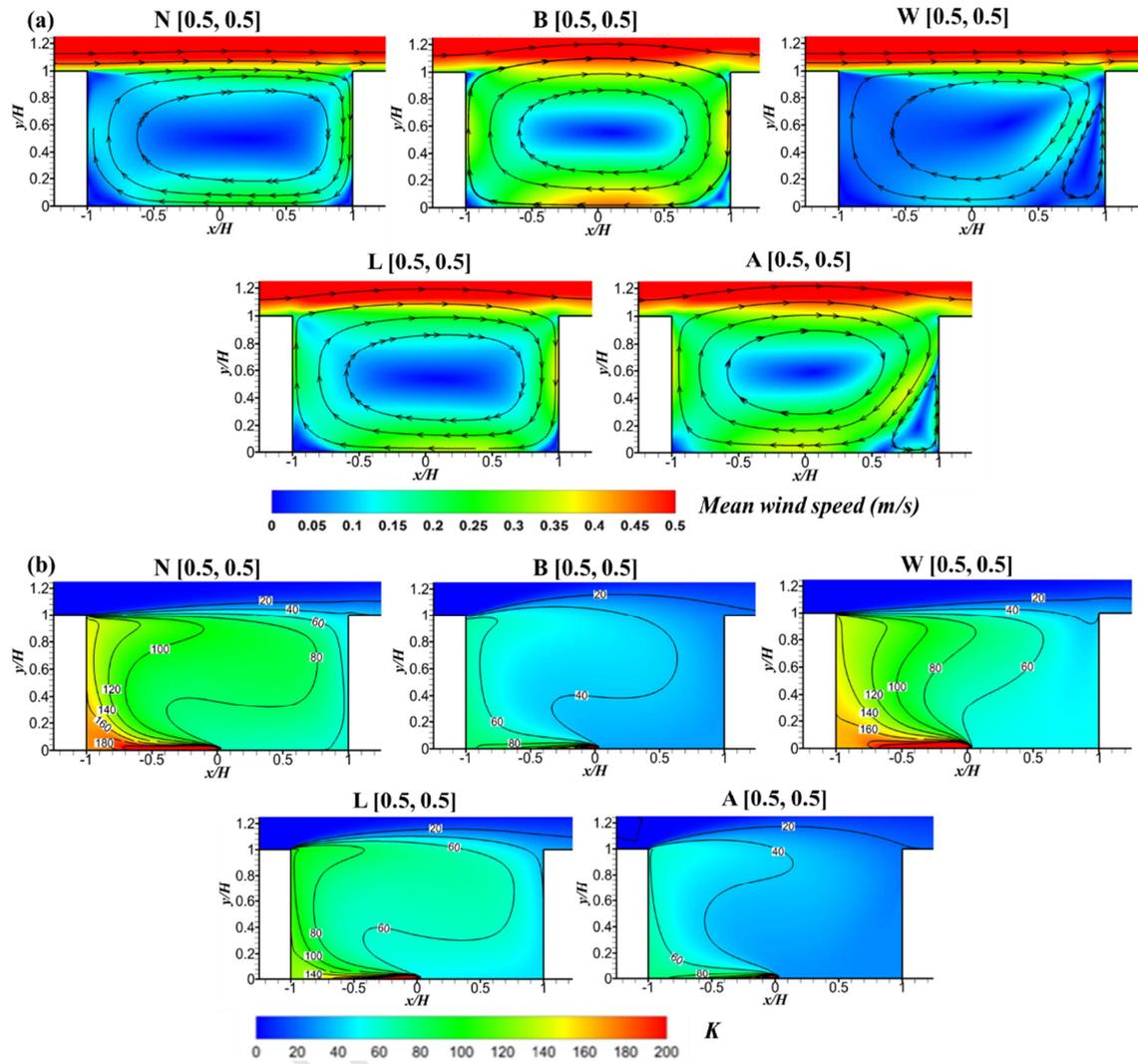


Figure 10. In cases where $AR=0.67$ and $U_{ref}=0.5 \text{ m s}^{-1}$: Contour of (a) the mean wind speed (m s^{-1}) and (b) K . u_y along (c) the windward line and (d) the leeward line; (e) u_x along the bottom line; (f) $\langle K \rangle$ at the windward and leeward walls.



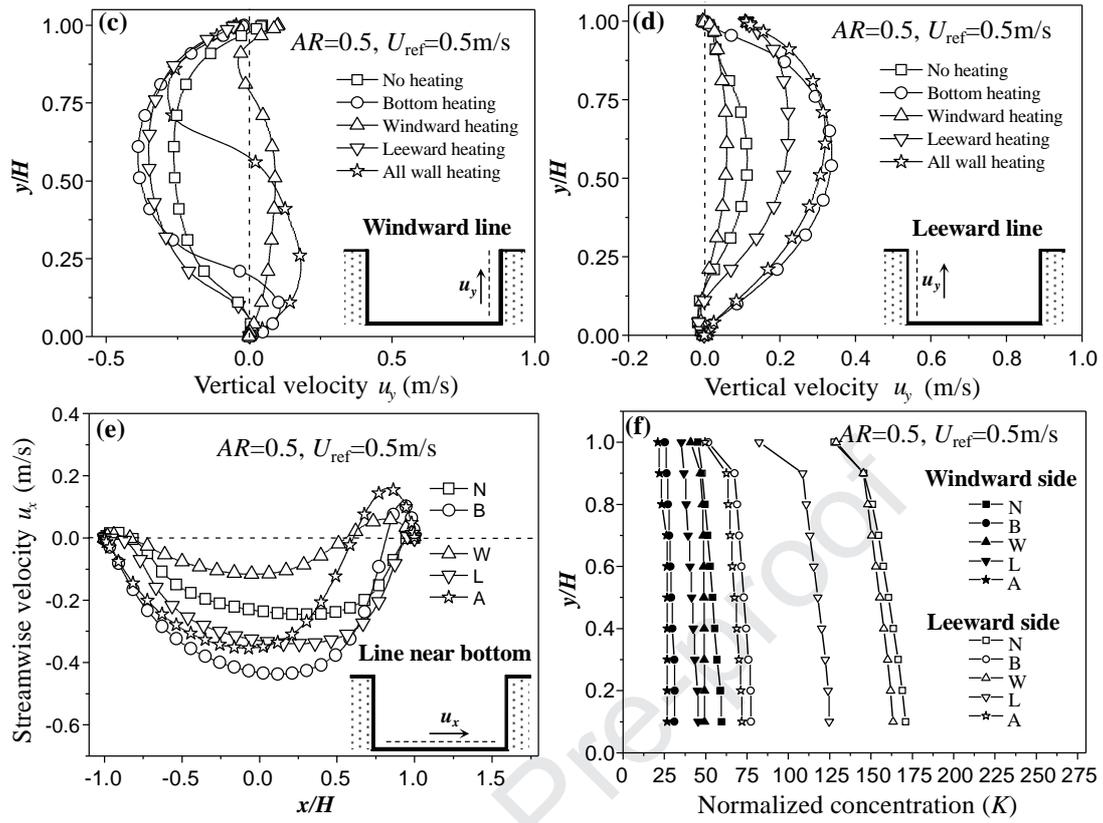
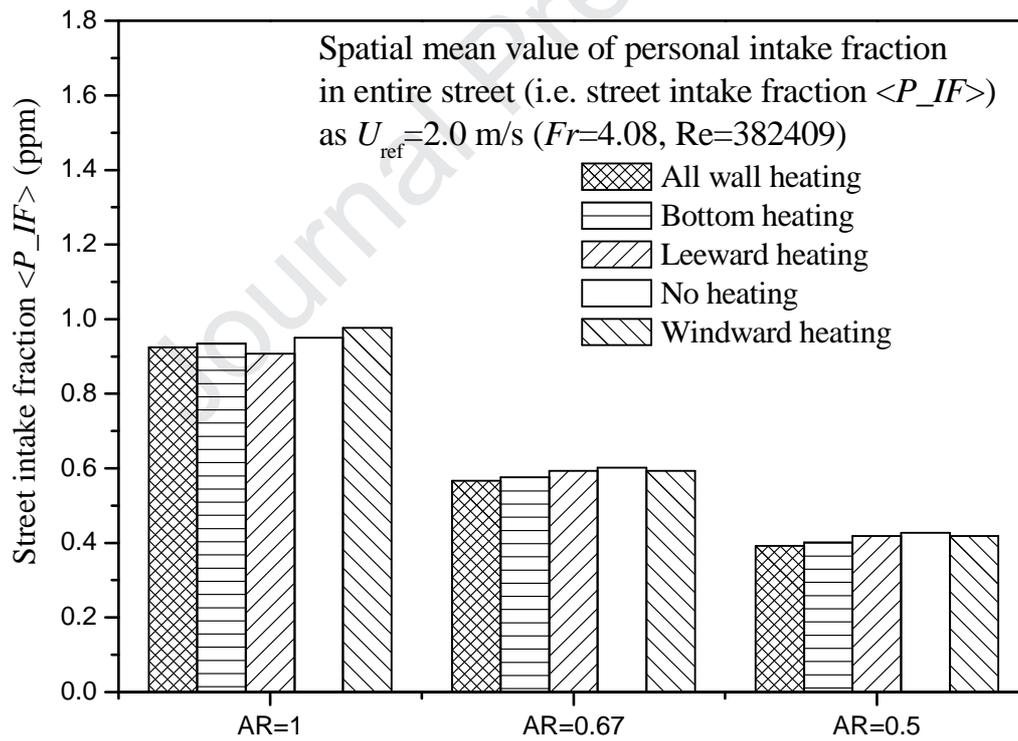
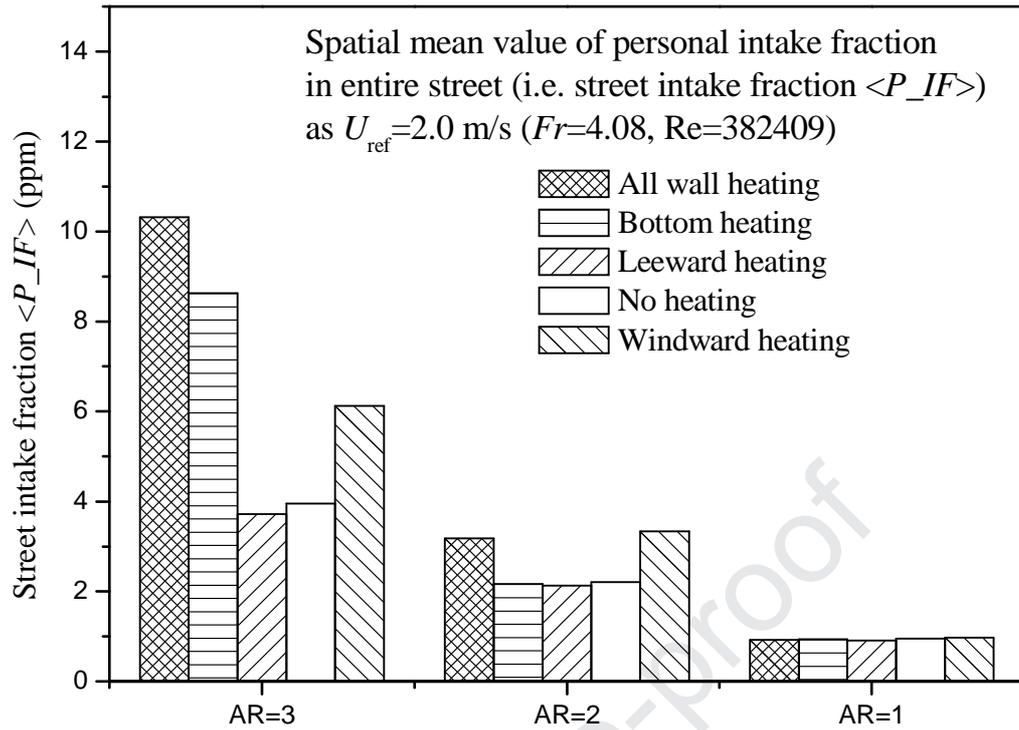


Figure 11. In cases where $AR=0.5$ and $U_{ref} = 0.5\text{ m s}^{-1}$. Contour maps of (a) the mean wind speed (m s^{-1}) and (b) K . u_y along (c) the windward line and (d) the leeward line; (e) u_x along the bottom line; (f) $\langle K \rangle$ at the windward wall and leeward wall.



(a)

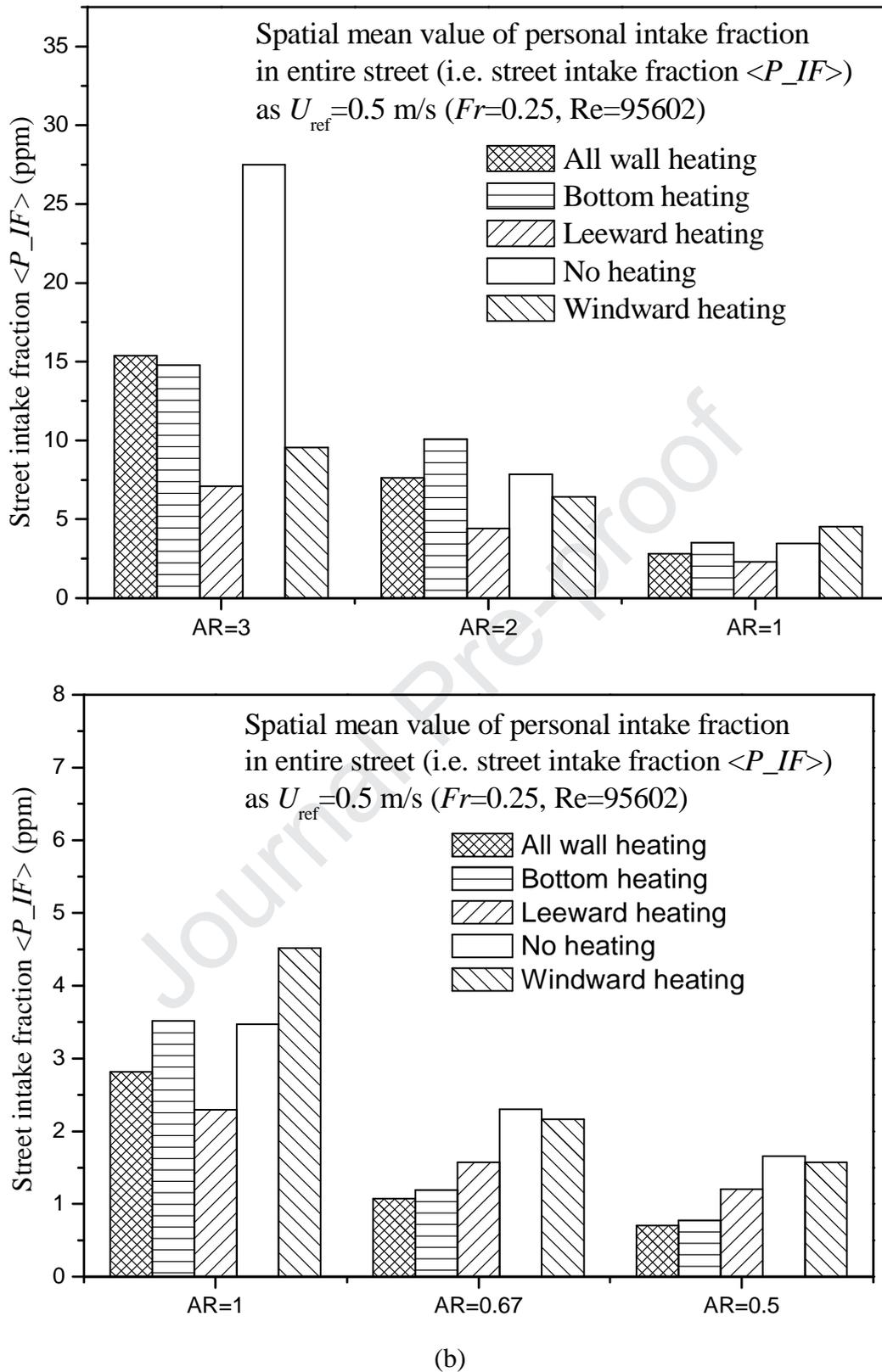


Figure 12. Spatial mean value of the personal intake fraction of a local street with different heating conditions as $H/W=AR=0.5-3$ with (a) $U_{ref}=2.0$ m s⁻¹, and (b) $U_{ref}=0.5$ m s⁻¹.

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Table 1 Breathing rate and time patterns for indoor at home for each age group [58,64-65]

Age groups	Population ratio	Breathing rate Br (m^3/day)	Time patterns
Children (<18)	21.2%	12.5	61.7%
Adults (18 - 60)	63.3%	13.8	59.5%
Elderly (> 60)	15.5%	13.1	71.6%

Table 2. Model descriptions of the simulated test cases.

Case Name: Heating type [AR, U_{ref}]			
Aspect ratio: H/W ($H = 3$ m)	ΔT (K)	The reference mean wind speed (U_{ref})	Heating type
$AR = 3$	10	0.5 m/s or 2.0 m/s	N (no heating)
$AR = 2$			B (bottom heating)
$AR = 1$			L (leeward heating)
$AR = 0.67$			W (windward heating)
$AR = 0.5$			A (all wall heating)

Table 3. Reynolds and Froude numbers investigated in all test cases with wall heating ($T_{wall}-T_{air}=10K, H=3m$)

Aspect ratio $AR=H/W$	Velocity in upstream free flow U_{ref}	Reynolds number ($Re=U_{ref}H/\nu$)	Froude number $F_r = \frac{U_{ref}^2}{gH(\Delta T / T_{ref})}$
3, 2, 1,	2.0m/s	382409	4.08
0.67, 0.5	0.5m/s	95602	0.25

Table 4. Boundary conditions and solver settings for the CFD simulations

Location	Type	Profiles/conditions
Inlet	Velocity inlet	$U_{in}(y) = U_{ref} \left(\frac{y - y_{ref}}{y_{ref}} \right)^\alpha, \quad V_{in}(y) = 0$ $k_{in}(y) = (U_{in}(y) \times I_{in})^2, \quad \varepsilon_{in}(y) = \frac{C_\mu^{3/4} k_i^{3/2}}{\kappa y}$ <p>Here $\alpha=0.22$, $C=0.09$, $I_{in}=0.1$, $\kappa=0.41$, $y_{ref}=H=3$ m</p>
Outlet	Outflow	Zero normal gradients of all flow variables
Top	Symmetry	Zero normal gradients of all flow variables
Street canyon wall	No slip	$v = 0 \quad \frac{\partial}{\partial y}(u, v, k, \varepsilon) = 0$
Solver settings		
Pressure-velocity coupling		SIMPLE algorithm
Discretization scheme		Second-order upwind scheme, implicit solver

Declaration of interests

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

The authors declare the following financial interests/personal relationships which may be considered as potential competing interests:

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1. As $Fr=4.08$, wind-driven force dominates the urban airflow as $AR=0.5-1$.
2. As $Fr=0.25$, most heating conditions would lead to a lower $\langle P_{IF} \rangle$.
3. Formation of single main vortex is the most efficient way to decrease the $\langle P_{IF} \rangle$.
4. Leeward heating condition always decreases the $\langle P_{IF} \rangle$ as $Fr=0.25$ and 4.08 .

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