# Toeplitz operators on the unit ball with locally integrable symbols 

Article
Published Version

Creative Commons: Attribution 4.0 (CC-BY)
Open Access
Hagger, R., Liu, C., Taskinen, J. and Virtanen, J. A. (2022) Toeplitz operators on the unit ball with locally integrable symbols. Integral Equations and Operator Theory, 94 (2). 17. ISSN 1420-8989 doi: https://doi.org/10.1007/s00020-022-02695-3 Available at https://centaur.reading.ac.uk/96053/

It is advisable to refer to the publisher's version if you intend to cite from the work. See Guidance on citing.

To link to this article DOI: http://dx.doi.org/10.1007/s00020-022-02695-3
Publisher: Springer

All outputs in CentAUR are protected by Intellectual Property Rights law, including copyright law. Copyright and IPR is retained by the creators or other copyright holders. Terms and conditions for use of this material are defined in the End User Agreement.

## www.reading.ac.uk/centaur

## CentAUR

Central Archive at the University of Reading
Reading's research outputs online

# Toeplitz Operators on the Unit Ball with Locally Integrable Symbols 

R. Hagger, C. Liu, J. Taskinen and J. A. Virtanen


#### Abstract

We study the boundedness of Toeplitz operators $T_{\psi}$ with locally integrable symbols on weighted harmonic Bergman spaces over the unit ball of $\mathbb{R}^{n}$. Generalizing earlier results for analytic function spaces, we derive a general sufficient condition for the boundedness of $T_{\psi}$ in terms of suitable averages of its symbol. We also obtain a similar "vanishing" condition for compactness. Finally, we show how these results can be transferred to the setting of the standard weighted Bergman spaces of analytic functions.


Mathematics Subject Classification. Primary 47B35.
Keywords. Toeplitz operator, Harmonic Bergman space, Bergman projection, Boundedness, Compactness.

## 1. Introduction

Denote by $\mathrm{d} V$ the normalized $n$-dimensional Lebesgue measure on the unit ball $\mathbb{B}_{n}$ of $\mathbb{R}^{n}$ with $n \geq 2$. For $\lambda>-1$ and $p \geq 1$, let $b_{\lambda}^{p}=b_{\lambda}^{p}\left(\mathbb{B}_{n}\right)$ be the harmonic Bergman space consisting of all harmonic complex-valued functions in $L_{\lambda}^{p}=L^{p}\left(\mathbb{B}_{n}, \mathrm{~d} V_{\lambda}\right)$, where $\mathrm{d} V_{\lambda}=c(n, \lambda)\left(1-|x|^{2}\right)^{\lambda} \mathrm{d} V$ and $c(n, \lambda)$ is a normalization constant (see (3.2)). We further write $b^{p}$ for $b_{0}^{p}$ and $L^{p}$ for $L_{0}^{p}$.

The Toeplitz operator $T_{\psi}$ with symbol $\psi$ is defined by

$$
\begin{equation*}
T_{\psi} f=P_{\lambda}(\psi f), \tag{1.1}
\end{equation*}
$$

where $P_{\lambda}$ is the orthogonal projection of $L_{\lambda}^{2}$ onto $b_{\lambda}^{2}$ (see (3.3)), $\psi$ is a measurable function on $\mathbb{B}_{n}$, and $f$ is harmonic in $\mathbb{B}_{n}$. It is known that $P_{\lambda}$ can be extended to a bounded projection from $L_{\lambda}^{p}$ onto $b_{\lambda}^{p}$ for $1<p<\infty$ (see [3]). It follows that $T_{\psi}$ is well defined and bounded on $b_{\lambda}^{p}$ whenever $\psi$ is bounded.

[^0]For $f \in L_{\lambda}^{1}, P_{\lambda} f$ is still well defined as a function (see (3.3) below), but we may have $P_{\lambda} f \notin L_{\lambda}^{1}$ in general.

The study of $T_{\psi}$ on $b^{p}$ goes back to [8], where it was proven for nonnegative symbols $\psi \in L^{1}$ that $T_{\psi}$ is bounded (compact) on $b^{p}$ if and only if the averaging function

$$
\begin{equation*}
x \mapsto \frac{1}{\left|E_{r}(x)\right|} \int_{E_{r}(x)} \psi \mathrm{d} V \tag{1.2}
\end{equation*}
$$

is bounded on $\mathbb{B}_{n}$ (vanishes as $|x| \rightarrow 1$ ), where $E_{r}(x)=\left\{y \in \mathbb{B}_{n}:|y-x|<\right.$ $r(1-|x|)\}$ with $r \in(0,1)$ and $\left|E_{r}(x)\right|$ denotes its volume. Further, it was shown that for symbols $\psi$ continuous on $\overline{\mathbb{B}}_{n}, T_{\psi}$ is compact on $b^{p}$ if and only if $\psi=0$ on $\partial \mathbb{B}_{n}$, which was generalized to $b_{\lambda}^{2}$ in [8]. It is easy to see that these conditions for boundedness (compactness) can also be formulated in terms of boundedness (vanishing) of the Berezin transform of $\psi$ or variants of the averaging function (1.2). Moreover, applying (1.2) to the positive and negative parts of $\operatorname{Re} \psi$ and $\operatorname{Im} \psi$ shows that if

$$
\begin{equation*}
\sup _{x \in \mathbb{B}_{n}} \frac{1}{\left|E_{r}(x)\right|} \int_{E_{r}(x)}|\psi| \mathrm{d} V<\infty \tag{1.3}
\end{equation*}
$$

then $T_{\psi}$ is bounded. However, the modulus in the integrand makes (1.3) far from being necessary if the symbol is not positive. In this paper we propose averages over different sets, certain spherical boxes, but also with the modulus outside of the integral; see (2.3). In particular, we give a new, weaker sufficient condition for the boundedness and compactness of Toeplitz operators on the weighted harmonic Bergman spaces $b_{\lambda}^{p}$ in Theorems 2.3 and 2.5. In Corollary 2.6 we present the corresponding results in the case of weighted Bergman spaces $A_{\lambda}^{p}$ of analytic functions on the unit ball of $\mathbb{C}^{n}$.

Most results about Toeplitz operators on $b_{\lambda}^{p}$ are generalizations from the setting of analytic Bergman spaces $A_{\lambda}^{p}$ using similar but also additional ideas related to non-analyticity. As in $b_{\lambda}^{p}$, characterizing bounded Toeplitz operators is an open problem in $A^{p}(\mathbb{D})$ (the unweighted Bergman space of the unit disk $\mathbb{D}$ of the complex plane $\mathbb{C}$ ) even in the case $p=2$. Compared with (1.3), much more general and weaker sufficient conditions were found for Toeplitz operators on $A^{p}(\mathbb{D})$ in [9] with further improvements in [10]. Roughly speaking, the conditions resemble (1.3), but the modulus appears outside the integral. This means that the modulus of a (wildly oscillating) symbol may be very large, but it may still induce a bounded operator due to cancellation in the integral; for an illustrative example, see [10]. Recently, the same idea was used to introduce weak BMO and VMO type conditions in [11]. Here, we prove similar results for the space $b_{\lambda}^{p}$ of the unit ball with standard weights for the sake of maximal generality, which carries a number of technical challenges. We note that recently in [12] it was shown that the mentioned sufficient conditions in $A^{p}$ are not necessary for the boundedness. We also refer to [12] for a concise account of the study of boundedness of Toeplitz operators on Bergman spaces.

As for the contents of this paper, the main results are formulated in Sect. 2 and their proofs are prepared in Sects. 3 and 4, which contain some
more necessary notation, definitions and preliminary lemmas. The proofs of the main results are completed in Sect. 5, and in Sect. 6 we construct examples of radial, oscillating symbols, the modulus of which may grow arbitrarily fast, but which still induce bounded Toeplitz operators.

## 2. The Main Results

In what follows we generalize the results of $[9,10]$ from the case of analytic functions in $\mathbb{D}$ to the case of harmonic functions in $\mathbb{B}_{n}$ and also consider weighted norms. Our analogous sufficient condition is a rather weak requirement of the boundedness of certain averages of $\psi$ over spherical boxes (see Theorem 2.3). We present the main results, Theorems 2.3-2.5 and Corollary 2.6 , with a minimal amount of notation.

It seems likely that analogous results hold for little Hankel operators on $b_{\lambda}^{p}$, too, but we do not consider this question here; cf. [10].

Let $\mathbb{Q}_{n}:=[0,1) \times[0, \pi]^{n-2} \times[0,2 \pi]$ and $\sigma: \mathbb{Q}_{n} \rightarrow \mathbb{B}_{n}$ be defined by

$$
\begin{aligned}
& \sigma\left(r, \theta_{2}, \ldots, \theta_{n-1}, \theta_{n}\right) \\
& =\left(r \cos \theta_{2}, r \sin \theta_{2} \cos \theta_{3}, r \sin \theta_{2} \sin \theta_{3} \cos \theta_{4}, \ldots\right.
\end{aligned} \quad \begin{aligned}
& \left.r \sin \theta_{2} \cdots \sin \theta_{n-1} \cos \theta_{n}, r \sin \theta_{2} \cdots \sin \theta_{n-1} \sin \theta_{n}\right) .
\end{aligned}
$$

We note that $\sigma$ is surjective, and injective almost everywhere (i.e. for almost every $x \in \mathbb{Q}_{n}, \sigma(x)=\sigma(y)$ implies $\left.x=y\right)$. For $x, y \in \mathbb{Q}_{n}$ we define the boxes

$$
Q(x, y):=\left\{z \in \mathbb{Q}_{n}: z_{j} \in\left[\min \left\{x_{j}, y_{j}\right\}, \max \left\{x_{j}, y_{j}\right\}\right]\right\}
$$

and $B(x, y):=\sigma(Q(x, y))$. In order to distinguish between $B(x, y) \subset \mathbb{B}_{n}$ and $Q(x, y) \subset \mathbb{Q}_{n}$, we call $B(x, y)$ a spherical box. For the following, we need to single out certain spherical boxes. Let

$$
\mathcal{K}:=\left\{\left(m, k_{2}, k_{3}, \ldots, k_{n}\right) \in \mathbb{N}_{0}^{n}: 0 \leq k_{n} \leq k_{n-1} \leq \cdots \leq k_{2} \leq 2^{m}-1\right\}
$$

For $k \in \mathcal{K}$ we define

$$
\begin{aligned}
C_{k}:=[ & \left.-2^{-m}, 1-2^{-m-1}\right] \times\left[\frac{\pi}{2} k_{2} 2^{-m}, \frac{\pi}{2}\left(k_{2}+1\right) 2^{-m}\right] \times\left[\frac{\pi}{2} \frac{k_{3}}{k_{2}+1}, \frac{\pi}{2} \frac{k_{3}+1}{k_{2}+1}\right] \\
& \times \cdots \times\left[\frac{\pi}{2} \frac{k_{n-1}}{k_{n-2}+1}, \frac{\pi}{2} \frac{k_{n-1}+1}{k_{n-2}+1}\right] \times\left[2 \pi \frac{k_{n}}{k_{n-1}+1}, 2 \pi \frac{k_{n}+1}{k_{n-1}+1}\right] .
\end{aligned}
$$

Note that $\bigcup_{k \in \mathcal{K}} C_{k}=[0,1) \times\left[0, \frac{\pi}{2}\right]^{n-2} \times[0,2 \pi]$. If we add all the sets that can be obtained from the sets $C_{k}$ by repeatedly using the reflections $\theta_{2} \mapsto$ $\pi-\theta_{2}, \ldots, \theta_{n-1} \mapsto \pi-\theta_{n-1}$, we obtain a cover of $\mathbb{Q}_{n}$. Let $\mathcal{Q}$ denote the collection of these sets. Enumerate the elements of $\mathcal{Q}$ by $Q_{1}, Q_{2}, \ldots$ and define $B_{j}:=\sigma\left(Q_{j}\right)$ for $j \in \mathbb{N}$. Then $\bigcup_{j \in \mathbb{N}} B_{j}$ covers $\mathbb{B}_{n}$ and $B_{j} \cap B_{k}$ is a null set whenever $j \neq k$. If $|x| \in\left[1-2^{-m}, 1-2^{-m-1}\right]$ for all $x \in B_{j}$, then $B_{j}$ is called a dyadic box of generation $m$. For a dyadic box $B_{j}$ of generation $m$, we denote $B_{j}^{*}:=B_{j}+2^{-m-2} \mathbb{B}_{n}$. The collection of these spherical boxes $B_{j}$ has the property that they are all of size comparable to $2^{-m n}$ and there is a constant $N \in \mathbb{N}$ such that every $x \in \mathbb{B}_{n}$ is contained in at most $N$ of the sets $B_{j}^{*}$. We will prove these facts in Lemma 4.3 below.

We remark that it may be tempting and more natural to extend the decomposition of the unit disk in [9] to higher dimensions using boxes of the form

$$
C_{k}=\left[1-2^{-m}, 1-2^{-m-1}\right] \times \prod_{j=2}^{n}\left[2^{-m} k_{j} \pi, 2^{-m}\left(k_{j}+1\right) \pi\right] \subset \mathbb{Q}_{n}
$$

where $k_{j}=1, \ldots, 2^{m}-1$ for $j=2, \ldots, n-1$ and $k_{n}=1, \ldots, 2^{m+1}-1$. After mapping them to $\mathbb{B}_{n}$ and some reordering, these boxes also provide a countable, pairwise essentially disjoint cover $\left\{B_{j}: j \in \mathbb{N}\right\}$ of $\mathbb{B}_{n}$. However, it turns out that the size of these boxes is not always comparable to $2^{-m n}$. This is due to the fact that the Jacobian of $\sigma$ is singular if any of $\theta_{2}, \ldots, \theta_{n-1}$ is equal to 0 or $\pi$. It also means that too many of the sets $B_{j}^{*}$ overlap close to the singular points. As these two properties are crucial for our analysis, we had to be more careful in the definition of the sets $B_{j}$. Having said that, it is highly expected that many other decompositions also lead to the same result. Our choice is not special in that regard, but it is somewhat natural as it leads to a definition of generalized Toeplitz operators that is independent of the decomposition; cf. Theorem 2.4.

The set of locally integrable functions on $\mathbb{B}_{n}$ will be denoted by $L_{\text {loc }}^{1}$. Note that $L_{\text {loc }}^{1}$ does not depend on $\lambda$. The characteristic function of a measurable set $E \subset \mathbb{B}_{n}$ will be denoted by $\chi_{E}$.
Definition 2.1. Let $\psi \in L_{\mathrm{loc}}^{1}, 1<p<\infty$ and assume that the series

$$
\begin{equation*}
T_{\psi} f(x):=\sum_{j=1}^{\infty} T_{\psi}\left(\chi_{B_{j}} f\right)(x)=\sum_{j=1}^{\infty} P_{\lambda}\left(\psi \chi_{B_{j}} f\right)(x) \tag{2.1}
\end{equation*}
$$

converges for almost every $x \in \mathbb{B}_{n}$ and all $f \in b_{\lambda}^{p}$. Then $T_{\psi}$ is called a generalized Toeplitz operator.

As $P_{\lambda}$ is bounded on $L_{\lambda}^{p}$, it is clear that $T_{\psi} f=P_{\lambda}(\psi f)$ whenever $\psi f \in$ $L_{\lambda}^{p}$. In particular, if $\psi$ is bounded, then $T_{\psi}$ is just the usual Toeplitz operator as defined in the introduction. However, without any further assumptions it is a priori not even clear if $T_{\psi} f \in b_{\lambda}^{p}$. Our first main result is that $T_{\psi}$ is a welldefined bounded linear operator under the "weak" Carleson-type condition (2.4).

For any measurable set $B \subset \mathbb{B}_{n}$ we denote the weighted volume by $|B|_{\lambda}:=\int_{B} \mathrm{~d} V_{\lambda}$. Moreover, $|B|:=|B|_{0}$.

Let $x^{(j)}$ and $y^{(j)}$ denote the smallest, respectively largest, element of $Q_{j}$ with respect to the partial order

$$
\begin{align*}
x \leq y \Longleftrightarrow & x_{1} \leq y_{1},\left|\frac{\pi}{2}-x_{2}\right| \geq\left|\frac{\pi}{2}-y_{2}\right|, \ldots,\left|\frac{\pi}{2}-x_{n-1}\right| \geq\left|\frac{\pi}{2}-y_{n-1}\right|, \\
& x_{n} \leq y_{n} \tag{2.2}
\end{align*}
$$

imposed on $\mathbb{Q}_{n}$.
Remark 2.2. Note that for $x, y \in[0,1) \times\left[0, \frac{\pi}{2}\right]^{n-2} \times[0,2 \pi]$ this is just the usual partial order of points in $\mathbb{R}^{n}$, which is then mirrored to all of $\mathbb{Q}_{n}$ to account for the construction of the sets $Q_{j}$ and $B_{j}$. In particular, the $x^{(j)}$ and $y^{(j)}$ are two opposite corners of $Q_{j}$ and we have $B_{j}=B\left(x^{(j)}, y^{(j)}\right)$. In the
following, we will only work in the subset $[0,1) \times\left[0, \frac{\pi}{2}\right]^{n-2} \times[0,2 \pi]$ and then use that everything is symmetric around $\theta_{2} \mapsto \pi-\theta_{2}, \ldots, \theta_{n-1} \mapsto \pi-\theta_{n-1}$. This simplifies the notation significantly and also means that we do not have to worry about the slightly unusual order (2.2).

Given a function $\psi \in L_{\text {loc }}^{1}$, we define

$$
\begin{equation*}
\widehat{\psi}_{j}:=\sup _{y \in B_{j}}\left|\int_{B\left(x^{(j)}, y\right)} \psi \mathrm{d} V_{\lambda}\right| . \tag{2.3}
\end{equation*}
$$

Theorem 2.3. Let $\psi \in L_{\mathrm{loc}}^{1}, 1<p<\infty$ and the family $\left(B_{j}\right)_{j \in \mathbb{N}}$ be as above. If there exists a constant $C_{\psi} \geq 0$ such that

$$
\begin{equation*}
\widehat{\psi}_{j} \leq C_{\psi}\left|B_{j}\right|_{\lambda} \tag{2.4}
\end{equation*}
$$

for all $j \in \mathbb{N}$, then the series (2.1) converges almost everywhere and in $L_{\lambda}^{p}$. Moreover, $T_{\psi}$ defines a bounded linear operator on $b_{\lambda}^{p}$ and there is a constant $C \geq 0$ independent of $\psi$ such that $\left\|T_{\psi}\right\| \leq C C_{\psi}$.

For $\psi \in L_{\text {loc }}^{1}$ and $0<\rho<1$, we define $\psi_{\rho}(z)=\psi(z)$ for $|z| \leq \rho$ and $\psi_{\rho}(z)=0$ for $\rho<|z|<1$. Note that $T_{\psi_{\rho}}$ is bounded on $b_{\lambda}^{p}$ by the previous theorem (see also Lemma 3.2 below). Our next theorem gives an alternative definition of Toeplitz operators with $L_{\text {loc }}^{1}$-symbol that does not depend on the decomposition $\left(B_{j}\right)_{j \in \mathbb{N}}$.
Theorem 2.4. Let $1<p<\infty$ and $1 / p+1 / q=1$, and suppose that $\psi \in L_{\text {loc }}^{1}$ satisfies (2.4). Then

$$
T_{\psi} f=\lim _{\rho \rightarrow 1} T_{\psi_{\rho}} f
$$

for all $f \in b_{\lambda}^{p}$ and $T_{\psi}^{*}: b_{\lambda}^{q} \rightarrow b_{\lambda}^{q}$ can be expressed as

$$
T_{\psi}^{*} f=\lim _{\rho \rightarrow 1} T_{\bar{\psi}_{\rho}} f
$$

for $f \in b_{\lambda}^{q}$.
Here, $T_{\psi}^{*}$ denotes the adjoint of the operator $T_{\psi}$ with respect to the standard duality of $b_{\lambda}^{p}$-spaces.

Concerning the compactness of the Toeplitz operator, we formulate the result in terms of the corresponding vanishing weak Carleson condition.

Theorem 2.5. Let the symbol $\psi$ satisfy the assumptions of Theorem 2.3 and in addition the condition

$$
\begin{equation*}
\lim _{j \rightarrow \infty} \frac{\widehat{\psi}_{j}}{\left|B_{j}\right|_{\lambda}}=0 \tag{2.5}
\end{equation*}
$$

Then $T_{\psi}: b_{\lambda}^{p} \rightarrow b_{\lambda}^{p}$ is compact for all $1<p<\infty$.
We emphasize that these sufficient conditions do not concern the modulus of the symbol $\psi$, only that of the integral.

Finally, we consider Toeplitz operators, denoted by $T_{\psi}^{\text {an }}$, in the case of the weighted Bergman spaces $A_{\lambda}^{p}=A_{\lambda}^{p}\left(\mathbf{B}_{n}\right)$ of the open unit ball $\mathbf{B}_{n}$ of $\mathbb{C}^{n}$. Here, $\mathbb{C}^{n}$ is identified in the canonical way with $\mathbb{R}^{2 n}$, hence $\mathbf{B}_{n}=\mathbb{B}_{2 n}$, and the

Bergman space $A_{\lambda}^{p}$ is the closed subspace of $L_{\lambda}^{p}\left(\mathbb{B}_{2 n}\right)$ consisting of analytic functions of $n$ complex variables in the domain $\mathbf{B}_{n}$. The Toeplitz operator $T_{\psi}^{\text {an }}$ with symbol $\psi \in L_{\text {loc }}^{1}\left(\mathbb{B}_{2 n}\right)$ is defined as in (1.1) by replacing the projection $P_{\lambda}$ by the orthogonal projection $P_{\lambda}^{\text {an }}$ from $L_{\lambda}^{p}\left(\mathbb{B}_{2 n}\right)$ onto $A_{\lambda}^{p}\left(\mathbf{B}_{n}\right)$, and the generalized Toeplitz operator corresponding to (2.1) is defined analogously.

Corollary 2.6. Let the assumptions of Theorem 2.3 (respectively, Theorem 2.5) be valid for the symbol $\psi$ in the domain $\mathbb{B}_{2 n}, n \in \mathbb{N}$. Then, the Toeplitz operator $T_{\psi}^{\mathrm{an}}: A_{\lambda}^{p}\left(\mathbf{B}_{n}\right) \rightarrow A_{\lambda}^{p}\left(\mathbf{B}_{n}\right)$ is bounded (resp. compact). The statement corresponding to Theorem 2.4 is also valid in the space $A_{\lambda}^{p}\left(\mathbf{B}_{n}\right)$.

The proofs of the main results will be given in Sect. 5 .

## 3. Further Notation and Definitions

By $C, c, C^{\prime}$ etc. we mean positive constants which may vary from place to place. If the constant depends on some parameter, say $n$, this is shown as $C(n)$. For positive valued expressions $f$ and $g$ depending on some variables or parameters, the notation $f \approx g$ (respectively, $f \lesssim g$ ) means the existence of constants $c, C$ such that $c f \leq g \leq C f$ (resp. $f \leq C g$ ) for all values of the variables or parameters. By $E(a, r) \subset \mathbb{R}^{n}$ we denote the open Euclidean ball with center $a \in \mathbb{R}^{n}$ and radius $r>0$.

On $\mathbb{B}_{n}$ we define the standard weight function

$$
\begin{equation*}
w(x)=1-|x|^{2} . \tag{3.1}
\end{equation*}
$$

Given $\lambda>-1$ and $1 \leq p<\infty$, the norm of the weighted space $L_{\lambda}^{p}$ and harmonic Bergman space $b_{\lambda}^{p} \subset L_{\lambda}^{p}$ is defined by

$$
\|f\|_{p, \lambda}:=\left(\int_{\mathbb{B}_{n}}|f|^{p} \mathrm{~d} V_{\lambda}\right)^{1 / p}
$$

where

$$
\begin{equation*}
\mathrm{d} V_{\lambda}:=c(n, \lambda) w^{\lambda} \mathrm{d} V:=\frac{2}{n} \cdot \frac{\Gamma(n / 2+\lambda+1)}{\Gamma(n / 2) \Gamma(\lambda+1)} w^{\lambda} \mathrm{d} V \tag{3.2}
\end{equation*}
$$

The orthogonal projection $P_{\lambda}$ from $L_{\lambda}^{2}$ onto $b_{\lambda}^{2}$ (the harmonic Bergman projection) can be expressed as an integral operator

$$
\begin{equation*}
P_{\lambda} f(x)=\int_{\mathbb{B}_{n}} f(y) R_{\lambda}(x, y) \mathrm{d} V_{\lambda}(y), \tag{3.3}
\end{equation*}
$$

where, for $x, y \in \mathbb{B}_{n}$,

$$
R_{\lambda}(x, y)=\frac{\Gamma(n / 2)}{\Gamma(n / 2+\lambda+1)} \sum_{k=0}^{\infty} \frac{\Gamma(k+n / 2+\lambda+1)}{\Gamma(k+n / 2)} Z_{k}(x, y)
$$

and $Z_{k}(x, y)$ denote the extended zonal harmonics of order $k$. We refer to [1, Chapter 5] for the definition of these functions and related facts. This series converges absolutely and uniformly on $K \times \overline{\mathbb{B}}_{n}$ for every compact set $K \subset \mathbb{B}_{n}$; see [7, Proposition 2.6]. In particular, $R_{\lambda}$ is a smooth bounded function on
$K \times \mathbb{B}_{n}$ and also on $\mathbb{B}_{n} \times K$ by symmetry. For $1<p<\infty$ the boundedness of $P_{\lambda}: L_{\lambda}^{p} \rightarrow b_{\lambda}^{p}$, defined as in (3.3), is proven in [3, Theorem 3.1].

The following pointwise estimate follows directly from the mean value property of harmonic functions.

Lemma 3.1. Let $\lambda>-1,1 \leq p<\infty$, and $f \in b_{\lambda}^{p}$. Then

$$
|f(x)| \lesssim \frac{\|f\|_{p, \lambda}}{w(x)^{(n+\lambda) / p}}
$$

for all $x \in \mathbb{B}_{n}$.
For compactly supported symbols we have the following result, which is well known to experts; however, we include its proof for completeness because we do not know a reference.

Lemma 3.2. Let $\psi \in L_{\lambda}^{1}$ have compact support, that is, $\operatorname{supp} \psi \subseteq r \overline{\mathbb{B}_{n}}$ for some $r<1$. Then, $P_{\lambda} \psi$ is harmonic and there is a constant $C=C(r)$ such that

$$
\left\|P_{\lambda} \psi\right\|_{p, \lambda} \leq C\|\psi\|_{1, \lambda}
$$

for all $1<p<\infty$. In particular, $T_{\psi}$ is well-defined and bounded on $b_{\lambda}^{p}$.
Proof. We have

$$
\left|P_{\lambda} \psi(x)\right| \leq \int_{r \overline{\mathbb{B}}_{n}}\left|\psi(y)\left\|R_{\lambda}(x, y)\left|\mathrm{d} V_{\lambda}(y) \leq C \int_{r \mathbb{\mathbb { B }}_{n}}\right| \psi(y) \mid \mathrm{d} V_{\lambda}(y)=C\right\| \psi \|_{1, \lambda}\right.
$$

for all $x \in \mathbb{B}_{n}$ because $R_{\lambda}(x, y)$ is bounded on $\mathbb{B}_{n} \times r \overline{\mathbb{B}}_{n}$ as noted above. This implies that $P_{\lambda} \psi$ is harmonic and $\left\|P_{\lambda} \psi\right\|_{p, \lambda} \leq C\|\psi\|_{1, \lambda}$ for all $p$.

By Lemma 3.1, $|f(y)| \leq \tilde{C}(r)\|f\|_{p, \lambda}$ for $|y| \leq r$. The same estimate as above thus yields the boundedness of $T_{\psi}$.

The maximal harmonic Bergman projection $P_{\lambda}^{M}$ is the nonlinear operator

$$
P_{\lambda}^{M} f(x)=\int_{\mathbb{B}_{n}}|f(y)|\left|R_{\lambda}(x, y)\right| \mathrm{d} V_{\lambda}(y)
$$

For $1<p<\infty$ this is a well-defined mapping $L_{\lambda}^{p} \rightarrow L_{\lambda}^{p}$, and it is also bounded in the sense that, for some constant

$$
\begin{equation*}
\left\|P_{\lambda}^{M} f\right\|_{p, \lambda} \leq C\|f\|_{p, \lambda} \tag{3.4}
\end{equation*}
$$

for $f \in L_{\lambda}^{p}$; see again [3].
Recall that for each $a \in \mathbb{B}_{n}$, the Möbius transformation $\varphi_{a}: \mathbb{B}_{n} \rightarrow \mathbb{B}_{n}$ is defined by the formula

$$
\begin{equation*}
\varphi_{a}(x)=\frac{|x-a|^{2} a-\left(1-|a|^{2}\right)(x-a)}{[x, a]^{2}} \tag{3.5}
\end{equation*}
$$

where

$$
\begin{equation*}
[x, a]:=\left(1-2 x \cdot a+|x|^{2}|a|^{2}\right)^{1 / 2} \tag{3.6}
\end{equation*}
$$

As is well known, $\varphi_{a}$ is an automorphism (analytic bijection) of $\mathbb{B}_{n}$ onto itself, which maps the point $a$ to the origin, and it is also an involution, i.e., $\varphi_{a} \circ \varphi_{a}(z)=z$ for all $z \in \mathbb{B}_{n}$.

The technical challenge of our paper arises from the decomposition of the unit ball into spherical boxes and a tricky integration by parts argument. Accordingly, it is important to introduce suitable combinatorial notation. First, we will use standard multi-index notation so that for a multi-index $\alpha=$ $\left(\alpha_{1}, \ldots, \alpha_{n}\right) \in \mathbb{N}_{0}^{n}$, we denote $|\alpha|:=\alpha_{1}+\cdots+\alpha_{n}$, and if $\beta=\left(\beta_{1}, \ldots, \beta_{n}\right) \in$ $\mathbb{N}_{0}^{n}$, then $\beta \leq \alpha$ means that $\beta_{k} \leq \alpha_{k}$ for all $k$. Given a multi-index $\alpha$, the corresponding partial differential operator acting on functions with $n$ real variables is defined by

$$
D^{\alpha}=\partial_{1}^{\alpha_{1}} \ldots \partial_{n}^{\alpha_{n}}
$$

where $\partial_{k}=\partial / \partial x_{k}$ for all $k$. Differentiation will be performed both in Cartesian and spherical coordinates in the following.

Moreover, in order to perform the integration by parts, we need some notation for certain parts of the boundary. For $x, y \in \mathbb{Q}_{n}$ and multi-indices $\alpha \in\{0,1\}^{n}$, let

$$
\begin{equation*}
Q_{\alpha}(x, y)=\prod_{\substack{k=1, \ldots, n \\ \alpha_{k}=1}}\left[x_{k}, y_{k}\right] \times \prod_{\substack{k=1, \ldots,, n \\ \alpha_{k}=0}}\left\{y_{k}\right\} . \tag{3.7}
\end{equation*}
$$

For $z \in Q(x, y)$ we denote by $z_{\alpha}$ the element in $Q(x, y)$ with coordinates

$$
\left(z_{\alpha}\right)_{k}= \begin{cases}z_{k} & \text { if } \alpha_{k}=1,  \tag{3.8}\\ y_{k} & \text { if } \alpha_{k}=0\end{cases}
$$

## 4. Preparatory Lemmas

### 4.1. Remarks Concerning the Hyperbolic Metric

The following facts about Möbius transforms $\varphi_{a}$ are well known and easy to check from the definitions, cf. (3.1), (3.5), (3.6).
Lemma 4.1. Let $a \in \mathbb{B}_{n}$. The identities

$$
\begin{equation*}
1-\left|\varphi_{a}(x)\right|^{2}=\frac{w(a) w(x)}{[x, a]^{2}} \quad \text { and } \quad\left|\varphi_{a}^{\prime}(x)\right|=\frac{w(a)}{[x, a]^{2}} \tag{4.1}
\end{equation*}
$$

hold for every $x \in \mathbb{B}_{n}$. Moreover, for all $x, y \in \mathbb{B}_{n}$,

$$
\begin{equation*}
\left[\varphi_{a}(x), \varphi_{a}(y)\right]^{2}=\left|\varphi_{a}^{\prime}(x)\right|\left|\varphi_{a}^{\prime}(y)\right|[x, y]^{2} \tag{4.2}
\end{equation*}
$$

The distance of two points $a, b \in \mathbb{B}_{n}$ in the hyperbolic (or Poincaré) metric is given by

$$
d(a, b)=\log \left(\frac{1+\left|\varphi_{a}(b)\right|}{1-\left|\varphi_{a}(b)\right|}\right)
$$

and, in particular, $d(0, a)=\log ((1+|a|) /(1-|a|))$.
Lemma 4.2. Let $a, b \in \mathbb{B}_{n}$. Then

$$
e^{-d(a, b)} \leq \frac{[x, a]}{[x, b]} \leq e^{d(a, b)}
$$

for all $x \in \mathbb{B}_{n}$.

Proof. We write $x=\varphi_{a}(y)$ and $b=\varphi_{a}(c)$. By (4.2), we have

$$
\frac{[x, a]^{2}}{[x, b]^{2}}=\frac{\left[\varphi_{a}(y), \varphi_{a}(0)\right]^{2}}{\left[\varphi_{a}(y), \varphi_{a}(c)\right]^{2}}=\frac{\left|\varphi_{a}^{\prime}(y) \| \varphi_{a}^{\prime}(0)\right|[y, 0]^{2}}{\left|\varphi_{a}^{\prime}(y)\right|\left|\varphi_{a}^{\prime}(c)\right|[y, c]^{2}}=\frac{[a, c]^{2}}{[y, c]^{2}}
$$

using (4.1) for $\left|\varphi_{a}^{\prime}(0)\right|$ and $\left|\varphi_{a}^{\prime}(c)\right|$. Since $c=\varphi_{a}(b)$, we have

$$
\frac{[x, a]}{[x, b]}=\frac{[a, c]}{[y, c]} \leq \frac{1+|c|}{1-|c|}=\frac{1+\left|\varphi_{a}(b)\right|}{1-\left|\varphi_{a}(b)\right|}=e^{d(a, b)} .
$$

The other inequality follows from interchanging the roles of $a$ and $b$.
Next we prove some crucial properties of the dyadic boxes $B_{j}$ and their enlargements $B_{j}^{*}$. Recall that the $B_{j}$ are essentially disjoint and $B_{j}^{*}=B_{j}+$ $2^{-m-2} \mathbb{B}_{n}$.

Lemma 4.3. Let $B_{j}$ be a dyadic box of generation $m$. Then
(i) $\left|B_{j}\right| \approx 2^{-m n}$,
(ii) $\operatorname{diam}\left(B_{j}\right) \approx 2^{-m}$,
(iii) $\left|B_{j}^{*}\right| \approx 2^{-m n}$,
(iv) there is a constant $N \in \mathbb{N}$ such that every $x \in \mathbb{B}_{n}$ is contained in at most $N$ of the sets $B_{j}^{*}$,
(v) $1-|x| \approx 1-|x|^{2} \approx 2^{-m}$ whenever $x \in B_{j}^{*}$,
(vi) $[x, a] \approx[x, b]$ for all $x \in \mathbb{B}_{n}$ and $a, b \in B_{j}^{*}$.

Proof. Without loss of generality, we may always assume that $B_{j}=\sigma\left(C_{k}\right)$ for some $k \in \mathcal{K}$.
(i) A direct computation gives

$$
\begin{aligned}
& r^{n-1} \sin ^{n-2} \theta_{2} \ldots \sin \theta_{n-1} \mathrm{~d} \theta_{n} \ldots \mathrm{~d} \theta_{2} \mathrm{~d} r \\
& \approx \int_{1-2^{-m}}^{1-2^{-m-1}} \int_{\frac{\pi}{2} k_{2} 2^{-m}}^{\frac{\pi}{2}\left(k_{2}+1\right) 2^{-m}} \cdots \int_{\frac{\pi}{2} \frac{k_{n-1}}{k_{n-2}+1}}^{\frac{\pi}{2} \frac{k_{n-1}+1}{k_{n-2}+1}} \int_{2 \pi \frac{k_{n}}{k_{n-1}+1}}^{2 \pi \frac{k_{n}+1}{k_{n-1}+1}} \\
& r^{n-1} \theta_{2}^{n-2} \ldots \theta_{n-1} \mathrm{~d} \theta_{n} \ldots \mathrm{~d} \theta_{2} \mathrm{~d} r \\
& \approx\left(\left(1-2^{-m-1}\right)^{n}-\left(1-2^{-m}\right)^{n}\right)\left(\left(\left(k_{2}+1\right) 2^{-m}\right)^{n-1}-\left(k_{2} 2^{-m}\right)^{n-1}\right) \\
& \times \prod_{l=3}^{n}\left(\left(\frac{k_{l}+1}{k_{l-1}+1}\right)^{n-l+1}-\left(\frac{k_{l}}{k_{l-1}+1}\right)^{n-l+1}\right) \\
& \approx 2^{-m}\left(\left(k_{2}+1\right)^{n-1}-k_{2}^{n-1}\right) 2^{-m(n-1)} \\
& \times \prod_{l=3}^{n} \frac{1}{\left(k_{l-1}+1\right)^{n-l+1}}\left(\left(k_{l}+1\right)^{n-l+1}-k_{l}^{n-l+1}\right) \\
& \approx 2^{-m n} \prod_{l=2}^{n-1} \frac{\left(k_{l}+1\right)^{n-l+1}-k_{l}^{n-l+1}}{\left(k_{l}+1\right)^{n-l}} \\
& \approx 2^{-m n} \text {. }
\end{aligned}
$$

(ii) The edges of the dyadic boxes are coordinate lines $\gamma_{l}$ of the coordinate transform $\sigma$. Those are of the form

$$
\begin{aligned}
\gamma_{l}(t)=\sigma\left(r, \theta_{2}, \ldots,\right. & \left., \ldots, \theta_{n}\right) \\
& \stackrel{\uparrow}{\theta_{l}}-\text { coordinate }
\end{aligned}
$$

The length $\ell$ of the edges can therefore be computed directly:

$$
\begin{aligned}
\ell\left(\gamma_{1}\right) & =\int_{1-2^{-m}}^{1-2^{-m-1}}\left\|\gamma_{1}^{\prime}(t)\right\| \mathrm{d} t=\int_{1-2^{-m}}^{1-2^{-m-1}} 1 \mathrm{~d} t=2^{-m-1} \\
\ell\left(\gamma_{2}\right) & =\int_{\frac{\pi}{2} k_{2} 2^{-m}}^{\frac{\pi}{2}\left(k_{2}+1\right) 2^{-m}}\left\|\gamma_{2}^{\prime}(t)\right\| \mathrm{d} t=\int_{\frac{\pi}{2} k_{2} 2^{-m}}^{\frac{\pi}{2}\left(k_{2}+1\right) 2^{-m}} r \mathrm{~d} t \lesssim 2^{-m}, \\
& \vdots \\
\ell\left(\gamma_{l}\right) & =\int_{\frac{\pi}{2} \frac{k_{l}}{\frac{\pi}{2} \frac{k_{l}+1}{k_{l-1}+1}}\left\|\gamma_{l}^{\prime}(t)\right\| \mathrm{d} t=\int_{\frac{\pi}{2}}^{\frac{\pi}{2} \frac{k_{l}+1}{k_{l-1}+1}} \frac{k_{l}}{k_{l-1}+1}}^{k_{l} \sin \theta_{2} \ldots \sin \theta_{l-1} \mathrm{~d} t} \\
& \lesssim\left(k_{2}+1\right) 2^{-m} \cdot \frac{k_{3}+1}{k_{2}+1} \cdot \ldots \cdot \frac{k_{l-1}+1}{k_{l-2}+1} \cdot \frac{1}{k_{l-1}+1}=2^{-m}, \\
& \vdots \\
\ell\left(\gamma_{n}\right) & =\int_{2 \pi \frac{k_{n}}{\frac{\pi}{2} \frac{k_{n}+1}{k_{n-1}+1}}\left\|\gamma_{n}^{\prime}(t)\right\| \mathrm{d} t=\int_{\frac{\pi}{2}}^{\frac{\pi}{2} \frac{k_{n}+1}{k_{n-1}+1}} \frac{k_{n}}{k_{n-1}+1}}^{r \sin \theta_{2} \ldots \sin \theta_{n-1} \mathrm{~d} t} \\
& \lesssim\left(k_{2}+1\right) 2^{-m} \cdot \frac{k_{3}+1}{k_{2}+1} \cdot \ldots \cdot \frac{k_{n-1}+1}{k_{n-2}+1} \cdot \frac{1}{k_{n-1}+1}=2^{-m} .
\end{aligned}
$$

It follows $\operatorname{diam}\left(B_{j}\right) \approx 2^{-m}$.
(iii) We only need to show $\left|B_{j}^{*}\right| \lesssim 2^{-m n}$. Clearly, $B_{j}^{*}$ is contained in a ball of diameter at most $\operatorname{diam}\left(B_{j}\right)+2^{-m-1} \approx 2^{-m}$. Hence $\left|B_{j}^{*}\right| \lesssim 2^{-m n}$.
(iv) Let $x \in \mathbb{B}_{n}$ with $|x| \in\left[1-2^{-m}, 1-2^{-m-1}\right]$ and denote the set of indices $j$ such that $B_{j} \cap E\left(x, 2^{-m-2}\right) \neq \emptyset$ by $\mathcal{J}_{x}$. It suffices to show that the cardinality of $\mathcal{J}_{x}$ is bounded independently of $x$. Clearly, $B_{j}$ is of generation $m-1, m$ or $m+1$ for every $j \in \mathcal{J}_{x}$. Therefore, by (i) and (ii), $\left|B_{j}\right| \approx 2^{-n m}$ and $\operatorname{diam}\left(B_{j}\right) \approx 2^{-m}$ for every $j \in \mathcal{J}_{x}$. It follows that there is a constant $C$ (only depending on $n$ ) such that $B_{j} \subset E\left(x, C 2^{-m}\right)$ for every $j \in \mathcal{J}_{x}$. Since $\left|E\left(x, C 2^{-m}\right)\right| \approx 2^{-m n}, E\left(x, C 2^{-m}\right)$ can only contain finitely many essentially disjoint sets of volume $\approx 2^{-m n}$. This finite number is the constant $N$ we are looking for; it only depends on $n$. This finishes the proof.
(v) As $1-2^{-m} \leq|x| \leq 1-2^{-m-1}$ for $x$ in a dyadic box of generation $m$, this is clear.
(vi) In view of (ii), there is a constant $C$ such that $|a-b| \leq C 2^{-m}$ for all $a, b \in B_{j}^{*}$. We also have $1-|a|^{2} \geq 2^{-m-2}$ and $1-|b|^{2} \geq 2^{-m-2}$. Hence,
by (3.6) and (4.1),

$$
\begin{aligned}
\frac{1}{1-\left|\varphi_{a}(b)\right|^{2}} & =\frac{[b, a]^{2}}{w(a) w(b)}=\frac{w(a) w(b)+|a|^{2}+|b|^{2}-2 a \cdot b}{w(a) w(b)} \\
& =1+\frac{|a-b|^{2}}{w(a) w(b)} \leq 1+16 C^{2}
\end{aligned}
$$

and therefore,

$$
d(a, b)=\log \left(\frac{\left(1+\left|\varphi_{a}(b)\right|\right)^{2}}{1-\left|\varphi_{a}(b)\right|^{2}}\right) \leq \log \left(\frac{4}{1-\left|\varphi_{a}(b)\right|^{2}}\right) \leq \log \left(4+64 C^{2}\right)
$$

The assertion (vi) now follows from Lemma 4.2.

### 4.2. Basic Facts About the Space $b_{\lambda}^{p}$

We will need a number of results on the spaces $b_{\lambda}^{p}$. The following was proven in [3, Lemma 2.8].

Lemma 4.4. For all multi-indices $\alpha \in \mathbb{N}_{0}^{n}$ and all $\lambda>-1$,

$$
\left|D_{x}^{\alpha} R_{\lambda}(x, y)\right| \lesssim[x, y]^{-n-\lambda-|\alpha|}, x, y \in \mathbb{B}_{n} .
$$

To prove the next lemma, we use the Forelli-Rudin type estimate

$$
\begin{equation*}
\int_{\mathbb{B}_{n}} \frac{\left(1-|y|^{2}\right)^{t}}{[x, y]^{n+s+t}} \mathrm{~d} V(y) \approx\left(1-|x|^{2}\right)^{-s} \tag{4.3}
\end{equation*}
$$

where $s>0$ and $t>-1$ (see [4, Proposition 2.2]).
Lemma 4.5. Let $\lambda>-1$ and $1<p<\infty$. Then the integral operator defined by

$$
\Lambda_{\lambda} f(x):=\int_{\mathbb{B}_{n}} \frac{f(y)}{[x, y]^{n+\lambda}} \mathrm{d} V_{\lambda}(y), \quad x \in \mathbb{B}_{n},
$$

is bounded on $L_{\lambda}^{p}$.
Proof. Define $h(x)=\left(1-|x|^{2}\right)^{\alpha}$ and choose $\frac{-1-\lambda}{\max \{p, q\}}<\alpha<0$. Then, by (4.3), when $1 / p+1 / q=1$, we get

$$
\int_{\mathbb{B}_{n}} \frac{h(y)^{q}}{[x, y]^{n+\lambda}} \mathrm{d} V_{\lambda}(y)=\int_{\mathbb{B}_{n}} \frac{\left(1-|y|^{2}\right)^{\alpha q+\lambda}}{[x, y]^{n+\lambda}} \mathrm{d} V(y) \lesssim h(x)^{q}
$$

and similarly $\int_{\mathbb{B}_{n}} \frac{h(x)^{p}}{[x, y]^{n+\lambda}} \mathrm{d} V_{\lambda}(x) \lesssim h(y)^{p}$. Now Schur's test completes the proof.

Lemma 4.6. Let $\lambda>-1,1<p<\infty$ and $k \in \mathbb{N}$. If $\alpha$ is a multi-index with $|\alpha|=k$ and $f \in b_{\lambda}^{p}$, then $w^{k} D^{\alpha} f \in L_{\lambda}^{p}$, and

$$
\left\|w^{k} D^{\alpha} f\right\|_{p, \lambda} \lesssim\|f\|_{p, \lambda}
$$

Proof. In view of Lemma 3.1, $b_{\lambda}^{p} \subset b_{\gamma}^{2}$ for sufficiently large $\gamma$, and hence

$$
f(x)=\int_{\mathbb{B}_{n}} R_{\gamma}(x, y) f(y) \mathrm{d} V_{\gamma}(y)
$$

Differentiating under the integral sign, we obtain

$$
\begin{equation*}
w(x)^{|\alpha|} D^{\alpha} f(x)=\int_{\mathbb{B}_{n}} K(x, y) f(y) \mathrm{d} V_{\lambda}(y) \tag{4.4}
\end{equation*}
$$

where

$$
K(x, y):=w(x)^{|\alpha|} w(y)^{\gamma-\lambda} D_{x}^{\alpha} R_{\gamma}(x, y)
$$

By Lemma 4.4,

$$
|K(x, y)| \lesssim \frac{w(x)^{|\alpha|} w(y)^{\gamma-\lambda}}{[x, y]^{n+\gamma+|\alpha|}} \lesssim \frac{1}{[x, y]^{n+\lambda}}
$$

where the last inequality follows from

$$
[x, y] \geq 1-|x||y| \geq \max \{1-|x|, 1-|y|\} \geq \frac{1}{2} \max \{w(x), w(y)\}
$$

Using (4.4) and Lemma 4.5, we get

$$
\left\|w^{|\alpha|} D^{\alpha} f\right\|_{p, \lambda} \lesssim\|f\|_{p, \lambda},
$$

which completes the proof.

### 4.3. Integration by Parts in Spherical Boxes

The results in [9] were based on a tricky integration by parts. We have to exploit this procedure in even higher generality, hence, it is useful to expose the corresponding general integration-by-parts-formula. The set of $n$-times continuously differentiable functions on $\mathbb{B}_{n}$ will be denoted by $C^{n}\left(\mathbb{B}_{n}\right)$. In the following, the Jacobian determinant of the coordinate transform $\sigma$ will be denoted by $J_{\sigma}$.

Lemma 4.7. Let $f \in L_{\text {loc }}^{1}, g \in C^{n}\left(\mathbb{B}_{n}\right)$ and $x, y \in \mathbb{Q}_{n}$. Then, with $F=f \circ \sigma$ and $G=g \circ \sigma$,

$$
\begin{align*}
\int_{B(x, y)} f g \mathrm{~d} V_{\lambda}= & \int_{Q(x, y)} F(\gamma) G(\gamma) J_{\sigma}(\gamma) w(\sigma(\gamma))^{\lambda} \mathrm{d} \gamma \\
= & \sum_{\alpha \in\{0,1\}^{n}} s_{\alpha} \int_{Q_{\alpha}(x, y)}\left(\int_{Q\left(x, \gamma_{\alpha}\right)} F(\tau) J_{\sigma}(\tau) w(\sigma(\tau))^{\lambda} \mathrm{d} \tau\right) \\
& \times D^{\alpha} G\left(\gamma_{\alpha}\right) \mathrm{d} \gamma_{\alpha}, \tag{4.5}
\end{align*}
$$

where $s_{\alpha} \in\{ \pm 1\}$ for all $\alpha \in\{0,1\}^{n}$.
Recall that $Q_{\alpha}(x, y)=\prod_{\substack{k=1, \ldots, n \\ \alpha_{k}=1}}\left[x_{k}, y_{k}\right] \times \prod_{\substack{k=1, \ldots, n \\ \alpha_{k}=0}}\left\{y_{k}\right\}$, which means that the integration $\int_{Q_{\alpha}(x, y)} \mathrm{d} \gamma_{\alpha}$ is performed only in those variables $\gamma_{k}$ where $\alpha_{k}=1$. This is in concordance with the notation (3.8) since $\left(\gamma_{\alpha}\right)_{k}=y_{k}$ if $\alpha_{k}=0$.

Proof. Let $\mathbb{1}:=(1, \ldots, 1)$. We will use the following well-known formula

$$
\begin{equation*}
\left(D^{\mathbb{1}} u\right) \cdot v=\sum_{\alpha \in\{0,1\}^{n}}(-1)^{|\alpha|} D^{\mathbb{1}-\alpha}\left(u \cdot D^{\alpha} v\right) \tag{4.6}
\end{equation*}
$$

which is easily proven by induction. Now choose

$$
u(\gamma):=\int_{Q(x, \gamma)} F(\tau) J_{\sigma}(\tau) w(\sigma(\tau))^{\lambda} \mathrm{d} \tau, \quad \text { and } \quad v(\gamma):=G(\gamma)
$$

and observe that $u(\gamma)=0$ if $x_{k}=\gamma_{k}$ for some $k$. Integrating the formula (4.6) over $Q(x, y)$ and using the fundamental theorem of calculus multiple times yields the result.

## 5. Proofs of the Main Results

Before giving the proofs of our main results, we still need to consider four lemmas. The first one is used to fix a small flaw in [9]: in the reference, the inequality (3.8) is not true as such, but the integration domain has to be replaced by a larger set. This, however, is not difficult, and we use here the enlarged dyadic boxes $B_{j}^{*}$ to this end.

Lemma 5.1. Let $f$ be a harmonic function on $\mathbb{B}_{n}$ and let $B_{j}$ be one of our dyadic boxes. Then for each $j$,

$$
\begin{equation*}
|f(x)| \lesssim \frac{1}{\left|B_{j}\right|_{\lambda}} \int_{B_{j}^{*}}|f| \mathrm{d} V_{\lambda} \tag{5.1}
\end{equation*}
$$

for every $x \in B_{j}$.
Proof. Suppose that $B_{j}$ is of generation $m$. By definition of the sets $B_{j}^{*}$, we have $E\left(x, \delta 2^{-m-2}\right) \subset B_{j}^{*}$ for every $x \in B_{j}$. The mean value property of harmonic functions yields

$$
f(x)=\frac{1}{\left|E\left(x, 2^{-m-2}\right)\right|} \int_{E\left(x, 2^{-m-2}\right)} f(y) \mathrm{d} V(y) .
$$

Thus,

$$
\begin{aligned}
|f(x)| & \leq \frac{1}{\left|E\left(x, 2^{-m-2}\right)\right|} \int_{E\left(x, 2^{-m-2}\right)}|f(y)| \mathrm{d} V(y) \\
& \lesssim 2^{n m} \int_{B_{j}^{*}}|f(y)| \mathrm{d} V(y) \\
& \lesssim 2^{(n+\lambda) m} \int_{B_{j}^{*}}|f(y)| w(y)^{\lambda} \mathrm{d} V(y)
\end{aligned}
$$

where the last inequality follows from the fact that $w(y) \approx 2^{-m}$ for $y \in B_{j}^{*}$; see Lemma 4.3. Since $\left|B_{j}\right|_{\lambda} \approx 2^{-(n+\lambda) m}$, this proves (5.1).

In order to apply Lemma 4.7, we need to transfer partial derivatives from spherical coordinates to Cartesian coordinates. For this we need the following lemma. We note that there is an explicit formula for iterated partial derivatives, the Faà di Bruno formula, but since we do not need it in full generality, we decided to give a proof in simple terms.

Lemma 5.2. Let $g \in C^{n}\left(\mathbb{B}_{n}\right)$ and $\alpha \in\{0,1\}^{n},|\alpha| \geq 1$. Then for every $\beta \in \mathbb{N}_{0}^{n}$ with $1 \leq|\beta| \leq|\alpha|$ there is a $d_{\beta} \in C^{\infty}\left(\mathbb{Q}_{n}\right)$ such that

$$
D_{\gamma}^{\alpha} g(\sigma(\gamma))=\sum_{1 \leq|\beta| \leq|\alpha|} c_{\beta}(\gamma) d_{\beta}(\gamma)\left(D_{x}^{\beta} g\right)(x)
$$

where

$$
c_{\beta}(\gamma)=\prod_{j=2}^{n-1}\left(\sin \theta_{j}\right)^{\max \left\{0,|\beta|-\sum_{i=1}^{j} \alpha_{i}\right\}},
$$

$\gamma=\left(r, \theta_{2}, \ldots, \theta_{n}\right) \in \mathbb{Q}_{n}$ and $x=\sigma(\gamma)$.
Proof. We only prove that one can factor out $\left(\sin \theta_{2}\right)^{|\beta|-\alpha_{1}-\alpha_{2}}$. As will be clear from the proof, factoring out the other terms in the product can then be done in the same way. We also note that the power of each $\sin \theta_{j}$ in the expansion of $D_{\gamma}^{\alpha}(g \circ \sigma)$ can never get negative. It can therefore be assumed that $|\beta|-\alpha_{1}-\alpha_{2} \geq 1$; otherwise there is nothing to prove.

We prove the claim by induction over $|\alpha|$. First assume that $|\alpha|=1$. Then

$$
D_{\gamma}^{\alpha}(g \circ \sigma)=\frac{\partial(g \circ \sigma)}{\partial \gamma_{k}}=\sum_{l=1}^{n} \frac{\partial x_{l}}{\partial \gamma_{k}} \frac{\partial g}{\partial x_{l}}
$$

for some $k \in\{1, \ldots, n\}$. Now observe that $\sin \theta_{2}$ is necessarily a factor of $\frac{\partial x_{l}}{\partial \gamma_{k}}$ unless $k \in\{1,2\}$. This completes the proof for $|\alpha|=1$.

To complete the induction process, we compute

$$
\frac{\partial}{\partial \gamma_{k}} \sum_{1 \leq|\beta| \leq|\alpha|}\left(\sin \theta_{2}\right)^{|\beta|-\alpha_{1}-\alpha_{2}} d_{\beta} D_{x}^{\beta} g
$$

for each $k \in\{1, \ldots, n\}$ with $\alpha_{k}=0$ and show that we can factor out sufficiently many $\sin \theta_{2}$. First assume $k \neq 2$. Then

$$
\begin{aligned}
& \frac{\partial}{\partial \gamma_{k}} \sum_{1 \leq|\beta| \leq|\alpha|}\left(\sin \theta_{2}\right)^{|\beta|-\alpha_{1}-\alpha_{2}} d_{\beta} D_{x}^{\beta} g \\
& =\sum_{1 \leq|\beta| \leq|\alpha|}\left(\sin \theta_{2}\right)^{|\beta|-\alpha_{1}-\alpha_{2}}\left(\frac{\partial d_{\beta}}{\partial \gamma_{k}} D_{x}^{\beta} g+d_{\beta} \sum_{l=1}^{n} \frac{\partial x_{l}}{\partial \gamma_{k}} \frac{\partial}{\partial x_{l}} D_{x}^{\beta} g\right) .
\end{aligned}
$$

Let $\alpha^{\prime}$ be equal to $\alpha$ except that the $k$-th entry is flipped from 0 to 1 . If $k \neq 1$, we can factor out another $\sin \theta_{2}$ from $\frac{\partial x_{l}}{\partial \gamma_{k}}$ as above and therefore the power of $\sin \theta_{2}$ again matches the order of the derivative in each term. If $k=1$, then $\frac{\partial x_{1}}{\partial \gamma_{1}}=\cos \theta_{2}$, so there is no additional $\sin \theta_{2}$ to factor out. However, since
$\alpha_{1}^{\prime}=\alpha_{1}+1$, we still have the correct power of $\sin \theta_{2}$ in the expansion. Now let $k=2$. In this case

$$
\begin{aligned}
& \frac{\partial}{\partial \gamma_{k}} \sum_{1 \leq|\beta| \leq|\alpha|}\left(\sin \theta_{2}\right)^{|\beta|-\alpha_{1}-\alpha_{2}} d_{\beta} D_{x}^{\beta} g \\
& =\sum_{1 \leq|\beta| \leq|\alpha|}\left(\sin \theta_{2}\right)^{|\beta|-\alpha_{1}-\alpha_{2}}\left(\frac{\partial d_{\beta}}{\partial \gamma_{k}} D_{x}^{\beta} g+d_{\beta} \sum_{l=1}^{n} \frac{\partial x_{l}}{\partial \gamma_{k}} \frac{\partial}{\partial x_{l}} D_{x}^{\beta} g\right) \\
& \quad+\left(\sin \theta_{2}\right)^{|\beta|-\alpha_{1}-\alpha_{2}-1} \cos \theta_{2} d_{\beta} D_{x}^{\beta} g .
\end{aligned}
$$

As $\alpha_{2}^{\prime}=\alpha_{2}+1$, we are in the same situation as for $k=1$. We again have the correct power of $\sin \theta_{2}$ in each term. This completes the proof.

The functions $c_{\beta}$ have the following important property.
Lemma 5.3. Let $\alpha \in\{0,1\}^{n},|\alpha| \geq 1$ and $c_{\beta}(\gamma)$ be as above. Then

$$
\int_{Q_{\alpha}\left(x^{(j)}, y^{(j)}\right)} c_{\beta}(\gamma) \mathrm{d} \gamma_{\alpha} \lesssim 2^{-m|\beta|}
$$

for all $j \in \mathbb{N}$.
Proof. As usual, we may assume that $Q\left(x^{(j)}, y^{(j)}\right)=C_{k}$ for some $k \in \mathcal{K}$. Choose $h \in \mathbb{N}$ such that $|\beta|=\sum_{i=1}^{h} \alpha_{i}$. If $h=1$, then $|\beta|=\alpha_{1}=1$ as well as $c_{\beta}(\gamma)=1$. Therefore,

$$
\int_{Q_{\alpha}\left(x^{(j)}, y^{(j)}\right)} c_{\beta}(\gamma) \mathrm{d} \gamma_{\alpha} \lesssim 2^{-m}=2^{-m|\beta|}
$$

If $h \geq 2$, another direct computation shows

$$
\begin{aligned}
\int_{Q_{\alpha}\left(x^{(j)}, y^{(j)}\right)} c_{\beta}(\gamma) \mathrm{d} \gamma_{\alpha} \lesssim & 2^{-m \alpha_{1}}\left(k_{2}+1\right)^{|\beta|-\alpha_{1}-\alpha_{2}} 2^{-m\left(|\beta|-\alpha_{1}\right)} \\
& \times \frac{\left(k_{3}+1\right)^{|\beta|-\alpha_{1}-\alpha_{2}-\alpha_{3}}}{\left(k_{2}+1\right)^{|\beta|-\alpha_{1}-\alpha_{2}}} \cdots \frac{\left(k_{n-1}+1\right)^{|\beta|-\sum_{l=1}^{h} \alpha_{l}}}{\left(k_{n-2}+1\right)^{|\beta|-\sum_{l=1}^{h-1} \alpha_{l}}} \\
= & 2^{-m|\beta|}
\end{aligned}
$$

The following lemma is the most important technical step in the proof of our theorems.

Lemma 5.4. Let $\psi \in L_{\mathrm{loc}}^{1}$. For $z \in \mathbb{B}_{n}$ and every dyadic box $B_{j}$, we have

$$
\left|T_{\psi}\left(\chi_{B_{j}} f\right)(z)\right| \lesssim \frac{\widehat{\psi}_{j}}{\left|B_{j}\right|_{\lambda}} \sum_{|\beta| \leq n_{B_{j}^{*}}} \int \frac{w(y)^{|\beta|}\left|D^{\beta} f(y)\right|}{[z, y]^{n+\lambda}} \mathrm{d} V_{\lambda}(y)
$$

Proof. Suppose that the generation of $B_{j}$ is $m$. We may further assume that $B_{j}=\sigma\left(C_{k}\right)$ for some $k \in \mathcal{K}$. Let $\gamma=\left(r, \theta_{2}, \ldots, \theta_{n}\right) \in \mathbb{Q}_{n}$ be such that $y=\sigma(\gamma) \in B_{j}$ and $\alpha \in\{0,1\}^{n},|\alpha| \geq 1$. By Lemma 4.4, we have

$$
\begin{equation*}
\left|D_{y}^{\alpha} R_{\lambda}(z, y)\right| \lesssim[z, y]^{-n-\lambda-|\alpha|} \tag{5.2}
\end{equation*}
$$

for every $z \in \mathbb{B}_{n}$. Using Lemma 5.2, we obtain

$$
\begin{aligned}
\left|D_{\gamma}^{\alpha}\left(R_{\lambda}(z, \sigma(\gamma)) f(\sigma(\gamma))\right)\right| & \leq \sum_{1 \leq|\beta| \leq|\alpha|} c_{\beta}(\gamma)\left|d_{\beta}(\gamma)\right|\left|D_{y}^{\beta}\left(R_{\lambda}(z, y) f(y)\right)\right| \\
& \lesssim \sum_{1 \leq|\beta| \leq|\alpha|} c_{\beta}(\gamma) \sum_{\tilde{\beta} \leq \beta}\left|D_{y}^{\beta-\tilde{\beta}} R_{\lambda}(z, y)\right|\left|D_{y}^{\tilde{\beta}} f(y)\right| \\
& \lesssim \sum_{1 \leq|\beta| \leq|\alpha|} c_{\beta}(\gamma) \sum_{\tilde{\beta} \leq \beta} 2^{m(|\beta|-|\tilde{\beta}|)}[z, y]^{-n-\lambda}\left|D_{y}^{\tilde{\beta}} f(y)\right|
\end{aligned}
$$

where the last inequality follows from (5.2) and $[z, \sigma(\gamma)] \geq 1-|\sigma(\gamma)| \approx 2^{-m}$ by Lemma 4.3. We now apply Lemma 5.1 to the functions $D^{\tilde{\beta}} f$ and use (v) and (vi) of Lemma 4.3. This yields

$$
\begin{align*}
& \left|D_{\gamma}^{\alpha}\left(R_{\lambda}(z, \sigma(\gamma)) f(\sigma(\gamma))\right)\right| \\
& \quad \lesssim \sum_{1 \leq|\beta| \leq|\alpha|} c_{\beta}(\gamma) \sum_{\tilde{\beta} \leq \beta} 2^{m(|\beta|-|\tilde{\beta}|)}[z, y]^{-n-\lambda} \frac{1}{\left|B_{j}\right|_{\lambda}} \int_{B_{j}^{*}}\left|D^{\tilde{\beta}} f\right| \mathrm{d} V_{\lambda} \\
& \quad \lesssim \sum_{1 \leq|\beta| \leq|\alpha|} c_{\beta}(\gamma) \sum_{\tilde{\beta} \leq \beta} 2^{m|\beta|} \frac{1}{\left|B_{j}\right|_{\lambda}} \int_{B_{j}^{*}} \frac{w(y)^{|\tilde{\beta}|}\left|D^{\tilde{\beta}} f(y)\right|}{[z, y]^{n+\lambda}} \mathrm{d} V_{\lambda}(y) \\
& \quad \lesssim \sum_{1 \leq|\beta| \leq|\alpha|} c_{\beta}(\gamma) 2^{m|\beta|} \frac{1}{\left|B_{j}\right|_{\lambda}} \int_{B_{j}^{*}} \frac{w(y)^{|\beta|}\left|D^{\beta} f(y)\right|}{[z, y]^{n+\lambda}} \mathrm{d} V_{\lambda}(y) \tag{5.3}
\end{align*}
$$

for all $z \in \mathbb{B}_{n}$. We are now going to apply the integration by parts lemma, Lemma 4.7. Recall that $B_{j}=B\left(x^{(j)}, y^{(j)}\right)=\sigma\left(Q\left(x^{(j)}, y^{(j)}\right)\right)$. Integrating (5.3) over $Q_{\alpha}\left(x^{(j)}, y^{(j)}\right)$ and using Lemma 5.3 yields

$$
\begin{align*}
& \int_{Q_{\alpha}\left(x^{(j)}, y^{(j)}\right)}\left|D_{\gamma}^{\alpha}\left(R_{\lambda}(z, \sigma(\gamma)) f(\sigma(\gamma))\right)\right| \mathrm{d} \gamma_{\alpha} \\
& \quad \lesssim \sum_{1 \leq|\beta| \leq|\alpha|} \frac{1}{\left|B_{j}\right|_{\lambda}} \int_{B_{j}^{*}} \frac{w(y)^{|\beta|}\left|D^{\beta} f(y)\right|}{[z, y]^{n+\lambda}} \mathrm{d} V_{\lambda}(y) \tag{5.4}
\end{align*}
$$

We now apply (4.5) to the integral

$$
\begin{equation*}
T_{\psi}\left(\chi_{B_{j}} f\right)(z)=\int_{B_{j}} \psi(y) f(y) R_{\lambda}(z, y) \mathrm{d} V_{\lambda}(y) \tag{5.5}
\end{equation*}
$$

For the factor $f$ in (4.5) we take the function $\psi$ and for $g$ the function $R_{\lambda}(z, \cdot) f$ with a fixed $z \in \mathbb{B}_{n}$. This yields (see the remark just after this proof)

$$
\begin{align*}
& \left|\int_{B_{j}} \psi(y) f(y) R_{\lambda}(z, y) \mathrm{d} V_{\lambda}(y)\right| \\
& \quad \leq \sum_{\alpha \in\{0,1\}^{n}} \int_{Q_{\alpha}\left(x^{(j)}, y^{(j)}\right)}\left|\int_{Q\left(x^{(j)}, \gamma_{\alpha}\right)} \psi(\sigma(\tau)) w(\sigma(\tau))^{\lambda} J_{\sigma}(\tau) \mathrm{d} \tau\right| \\
& \quad \times\left|D_{\gamma}^{\alpha}\left(R_{\lambda}\left(z, \sigma\left(\gamma_{\alpha}\right)\right) f\left(\sigma\left(\gamma_{\alpha}\right)\right)\right)\right| \mathrm{d} \gamma_{\alpha} . \tag{5.6}
\end{align*}
$$

The change of the variables formula turns the integral over the set $Q\left(x^{(j)}, \gamma_{\alpha}\right)$ into

$$
\begin{equation*}
\left|\int_{B\left(x^{(j)}, y\right)} \psi \mathrm{d} V_{\lambda}\right| \leq \widehat{\psi}_{j} \tag{5.7}
\end{equation*}
$$

for $y=\sigma\left(\gamma_{\alpha}\right) \in B_{j}$, see (2.3). Applying (5.4), we thus see that (5.6) is bounded by a constant times

$$
\begin{equation*}
\frac{\widehat{\psi}_{j}}{\left|B_{j}\right|_{\lambda}} \sum_{|\beta| \leq n_{B_{j}^{*}}} \int \frac{w(y)^{|\beta|}\left|D^{\beta} f(y)\right|}{[z, y]^{n+\lambda}} \mathrm{d} V_{\lambda}(y) \tag{5.8}
\end{equation*}
$$

since $|\alpha| \leq n$. So the lemma follows by combining (5.5)-(5.8).
We emphasize that the deduction (5.6) is in the core of our result: By using integration by parts, it is possible to make estimates only in terms of the modulus of the integral of $\psi$, and a direct bound involving the modulus of $\psi$ can be avoided.

We are now in the position to prove our main results:
Proof of Theorem 2.3. Recall that the Toeplitz operator $T_{\psi}$ is defined with the help of the series

$$
T_{\psi} f(x)=\sum_{j=1}^{\infty} T_{\psi}\left(\chi_{B_{j}} f\right)(x)
$$

cf. (2.1). We show that the series converges absolutely for almost every $x \in \mathbb{B}_{n}$ and that the resulting operator is bounded. Indeed, by the previous lemma, assumption (2.4) and Lemma 4.3, we have

$$
\begin{align*}
\sum_{j=1}^{\infty}\left|T_{\psi}\left(\chi_{B_{j}} f\right)(x)\right| & \lesssim C_{\psi} \sum_{j=1}^{\infty} \sum_{|\beta| \leq n_{B_{j}^{*}}} \int \frac{w(y)^{|\beta|}\left|D^{\beta} f(y)\right|}{[x, y]^{n+\lambda}} \mathrm{d} V_{\lambda}(y) \\
& \lesssim C_{\psi} \sum_{|\beta| \leq n_{\mathbb{B}_{n}}} \int \frac{w(y)^{|\beta|}\left|D^{\beta} f(y)\right|}{[x, y]^{n+\lambda}} \mathrm{d} V_{\lambda}(y) \tag{5.9}
\end{align*}
$$

for all $x \in \mathbb{B}_{n}$. By Lemmas 4.5 and 4.6 , we see that for all $|\beta| \leq n$,

$$
\left\|\int_{\mathbb{B}_{n}} \frac{w(y)^{|\beta|}\left|D^{\beta} f(y)\right|}{[x, y]^{n+\lambda}} \mathrm{d} V_{\lambda}(y)\right\|_{p, \lambda} \lesssim\left\|w^{|\beta|} D^{\beta} f\right\|_{p, \lambda} \lesssim\|f\|_{p, \lambda} .
$$

This implies that the series $\sum_{j=0}^{\infty}\left|T_{\psi}\left(\chi_{B_{j}} f\right)(x)\right|$ is pointwise bounded by an $L_{\lambda}^{p}$-function, and thus it converges for almost all $x \in \mathbb{B}_{n}$. Moreover, the above argument implies

$$
\begin{equation*}
\left\|\sum_{j=1}^{\infty} T_{\psi}\left(\chi_{B_{j}} f\right)\right\|_{p, \lambda} \lesssim C_{\psi}\|f\|_{p, \lambda} \tag{5.10}
\end{equation*}
$$

By dominated convergence, the series also converges in $L_{\lambda}^{p}$. In particular, $T_{\psi} f \in b_{\lambda}^{p}$ for all $f \in b_{\lambda}^{p}$ and $\left\|T_{\psi}\right\| \lesssim C_{\psi}$.

Proof of Theorem 2.4. We first show that $T_{\psi_{\rho}} f \rightarrow T_{\psi} f$ for every $f \in b_{\lambda}^{p}$. As $\left(\widehat{\psi}_{\rho}\right)_{j} \leq \widehat{\psi}_{j}$, the estimate (5.9) is uniform in $\rho$. That is, for almost every $x \in \mathbb{B}_{n}$, the series $\sum_{j \in \mathbb{N}} T_{\psi_{\rho}}\left(\chi_{B_{j}} f\right)(x)$ converges absolutely and uniformly in $\rho$. In particular,

$$
\lim _{\rho \rightarrow 1} \sum_{j=1}^{\infty} T_{\psi_{\rho}}\left(\chi_{B_{j}} f\right)(x)=\sum_{j=1}^{\infty} \lim _{\rho \rightarrow 1} T_{\psi_{\rho}}\left(\chi_{B_{j}} f\right)(x)=\sum_{j=1}^{\infty} T_{\psi}\left(\chi_{B_{j}} f\right)(x)
$$

Moreover, $\left\|T_{\psi_{\rho}} f\right\|_{p, \lambda} \lesssim C_{\psi}\|f\|$ by (5.10). By dominated convergence, we get $T_{\psi_{\rho}} f \rightarrow T_{\psi} f$ as $\rho \rightarrow 1$. For the adjoint observe

$$
\left\langle T_{\psi} f, g\right\rangle=\sum_{j=1}^{\infty}\left\langle\psi \chi_{B_{j}} f, g\right\rangle=\sum_{j=1}^{\infty}\left\langle f, \bar{\psi} \chi_{B_{j}} g\right\rangle=\left\langle f, T_{\bar{\psi}} g\right\rangle
$$

and thus $T_{\bar{\psi}_{\rho}} g \rightarrow T_{\psi}^{*} g$ follows analogously.
Proof of Theorem 2.5. A routine normal families argument shows that the unit ball of $b_{\lambda}^{p}$ is compact in the topology of uniform convergence on compact subsets of $\mathbb{B}_{n}$. It therefore suffices to show that $\left\|T_{\psi} f_{k}\right\|_{p, \lambda} \rightarrow 0$ for all norm bounded sequences $\left(f_{k}\right)_{k=1}^{\infty} \subset b_{\lambda}^{p}$ which converge to zero uniformly on compacta of $\mathbb{B}_{n}$.

So, we fix such a sequence $\left(f_{k}\right)_{k=1}^{\infty}$ with $\left\|f_{k}\right\|_{p, \lambda} \leq 1$ for all $k$. Let $\varepsilon>0$ be arbitary. For all $j, k \in \mathbb{N}$ Lemma 5.4 implies

$$
\left|T_{\psi}\left(\chi_{B_{j}} f_{k}\right)(x)\right| \lesssim \frac{\widehat{\psi}_{j}}{\left|B_{j}\right|_{\lambda}} \sum_{|\beta| \leq n_{B_{j}^{*}}} \int \frac{w(y)^{|\beta|}\left|D^{\beta} f_{k}(y)\right|}{[x, y]^{n+\lambda}} \mathrm{d} V_{\lambda}(y)
$$

and the weak vanishing Carleson condition (2.5) implies

$$
\lim _{j \rightarrow \infty} \frac{\widehat{\psi}_{j}}{\left|B_{j}\right|_{\lambda}}=0
$$

We choose $N \in \mathbb{N}$ such that $\frac{\hat{\psi}_{j}}{\left|B_{j}\right|_{\lambda}}<\varepsilon$ for $j>N$.
As the sets $B_{j}^{*}$ are bounded, there exists a constant $C>0$ such that

$$
\sum_{j=1}^{N} \sum_{|\beta| \leq n_{B_{j}^{*}}} \int_{[x, y]^{n+\lambda}} \frac{w(y)^{|\beta|}}{\left[V_{\lambda}(y) \leq C\right.}
$$

for all $x \in \mathbb{B}_{n}$. Since the sequence $\left(f_{k}\right)_{k=1}^{\infty}$, as well as the sequences of the derivatives of $f_{k}$, converge to 0 uniformly on compact subsets, we may choose $M \in \mathbb{N}$ such that

$$
\left|D^{\beta} f_{k}(y)\right| \leq \varepsilon
$$

for all $k \geq M,|\beta| \leq n, y \in B_{j}^{*}$ and $j \leq N$.
For $k \geq M$ we get

$$
\begin{aligned}
\sum_{j=1}^{\infty}\left|T_{\psi}\left(\chi_{B_{j}} f_{k}\right)(x)\right| \lesssim & C_{\psi} \sum_{j=1}^{N} \sum_{|\beta| \leq n} \int_{B_{j}^{*}} \frac{w(y)^{|\beta|}\left|D^{\beta} f_{k}(y)\right|}{[x, y]^{n+\lambda}} \mathrm{d} V_{\lambda}(y) \\
& +\varepsilon \sum_{j=N+1}^{\infty} \sum_{|\beta| \leq n} \int_{B_{j}^{*}} \frac{w(y)^{|\beta|}\left|D^{\beta} f_{k}(y)\right|}{[x, y]^{n+\lambda}} \mathrm{d} V_{\lambda}(y) \\
\lesssim & C_{\psi} \varepsilon+\varepsilon \sum_{|\beta| \leq n_{\mathbb{B}_{n}}} \int \frac{w(y)^{|\beta|}\left|D^{\beta} f_{k}(y)\right|}{[x, y]^{n+\lambda}} \mathrm{d} V_{\lambda}(y)
\end{aligned}
$$

Arguing as in the proof of Theorem 2.3 we find that $\left\|T_{\psi} f_{k}\right\|_{p, \lambda} \lesssim \varepsilon$ for sufficiently large $k$.

Proof of Corollary 2.6. As $P_{\lambda}^{\text {an }}$ and $P_{\lambda}$ are orthogonal projections and obviously $A_{\lambda}^{p}\left(\mathbf{B}_{n}\right) \subset b_{\lambda}^{p}\left(\mathbb{B}_{2 n}\right)$, we have $P_{\lambda}^{\text {an }}=P_{\lambda}^{\text {an }} P_{\lambda}$ on $L_{\lambda}^{2}=L_{\lambda}^{2}\left(\mathbb{B}_{2 n}\right)$. Let $f \in L_{\lambda}^{1}$ have compact support. Choose a sequence $\left(f_{n}\right)_{n \in \mathbb{N}}$ in $L_{\lambda}^{2}$ with the same support as $f$ such that $\left\|f-f_{n}\right\|_{1, \lambda} \rightarrow 0$ as $n \rightarrow \infty$. Then

$$
\begin{equation*}
\left\|P_{\lambda}\left(f-f_{n}\right)\right\|_{p, \lambda} \lesssim\left\|f-f_{n}\right\|_{1, \lambda} \text { and }\left\|P_{\lambda}^{\mathrm{an}}\left(f-f_{n}\right)\right\|_{p, \lambda} \lesssim\left\|f-f_{n}\right\|_{1, \lambda} \tag{5.11}
\end{equation*}
$$

by Lemma 3.2 and the respective result for the analytic case, which can be proven verbatim. In particular,

$$
P_{\lambda}^{\mathrm{an}} P_{\lambda} f=P_{\lambda}^{\mathrm{an}} P_{\lambda}\left(f-f_{n}\right)+P_{\lambda}^{\mathrm{an}} P_{\lambda} f_{n}=P_{\lambda}^{\mathrm{an}} P_{\lambda}\left(f-f_{n}\right)+P_{\lambda}^{\mathrm{an}} f_{n} .
$$

By (5.11), the former converges to 0 and the latter converges to $P_{\lambda}^{\text {an }} f$ as $n \rightarrow \infty$. Hence $P_{\lambda}^{\text {an }} f=P_{\lambda}^{\text {an }} P_{\lambda} f$ for $f \in L_{\lambda}^{1}$ with compact support. Now let $\psi \in L_{\mathrm{loc}}^{1}$ and $f \in A_{\lambda}^{p}$. Then

$$
\sum_{j=1}^{\infty} T_{\psi}^{\mathrm{an}}\left(\chi_{B_{j}} f\right)=\sum_{j=1}^{\infty} P_{\lambda}^{\mathrm{an}}\left(\psi \chi_{B_{j}} f\right)=\sum_{j=1}^{\infty} P_{\lambda}^{\mathrm{an}} P_{\lambda}\left(\psi \chi_{B_{j}} f\right)
$$

As the series $\sum_{j=1}^{\infty} P_{\lambda}\left(\psi \chi_{B_{j}} f\right)$ converges in $L_{\lambda}^{p}$ by Theorem 2.3, and $P_{\lambda}^{\text {an }}$ is continuous, we obtain

$$
\sum_{j=1}^{\infty} P_{\lambda}^{\mathrm{an}} P_{\lambda}\left(\psi \chi_{B_{j}} f\right)=P_{\lambda}^{\mathrm{an}}\left(\sum_{j=1}^{\infty} P_{\lambda}\left(\psi \chi_{B_{j}} f\right)\right)=P_{\lambda}^{\mathrm{an}} T_{\psi} f .
$$

In short, $T_{\psi}^{\mathrm{an}}=\left.P_{\lambda}^{\mathrm{an}} T_{\psi}\right|_{A_{\lambda}^{p}}$. Properties such as boundedness and compactness therefore transfer from the harmonic to the analytic setting.

## 6. An Example

Sufficient conditions for the boundedness of Toeplitz operators with radial symbols in Bergman spaces $A_{\lambda}^{p}\left(\mathbf{B}_{n}\right)$ have been given for example in [2], where $p=2$, or in [5], where $n=1$. All such results concern symbols $\psi$ belonging at least to $L^{1}\left(\mathbf{B}_{n}\right)$. Modifying an example first presented in [9] we show here that there are radial symbols with arbitrarily fast growing modulus, still
satisfying the sufficient condition (2.4) of Theorem 2.3 and thus inducing bounded Toeplitz operators.

If $\lambda>-1$ and $\psi=\psi(r)$ is a radial symbol on the ball $\mathbb{B}_{n}$, then our condition (2.4) for boundedness simplifies to

$$
\begin{equation*}
\sup _{m \in \mathbb{N}_{0}} 2^{m(1+\lambda)} \sup _{\rho \in\left[1-2^{-m}, 1-2^{-m-1}\right]}\left|\int_{1-2^{-m}}^{\rho} r^{n-1} \psi(r)\left(1-r^{2}\right)^{\lambda} \mathrm{d} r\right|<\infty \tag{6.1}
\end{equation*}
$$

Let $f:[1, \infty) \rightarrow \mathbb{R}$ be any continuous function with $\inf _{x \in[1, \infty)} f(x) x^{\lambda}>0$. Define

$$
g:[1, \infty) \rightarrow[0, \infty), \quad g(x)=\int_{1}^{x} \frac{f(y)}{y^{1-\lambda}} \mathrm{d} y
$$

and note that $g^{\prime}(x)=\frac{f(x)}{x^{1-\lambda}}>0$ and $g(x) \rightarrow \infty$ as $x \rightarrow \infty$. In particular, $g$ is invertible. Define

$$
\psi(r)=r^{-n+1}\left(1-r^{2}\right)^{-\lambda} f\left((1-r)^{-1}\right) \exp \left(i \pi g\left((1-r)^{-1}\right)\right) .
$$

It is plain that the modulus of $\psi$ can be made to grow arbitrarily fast as $r \rightarrow 1$, nevertheless, we claim that $\psi$ satisfies (6.1) and thus by Theorem 2.3, $T_{\psi}: b_{\lambda}^{p} \rightarrow b_{\lambda}^{p}$ is bounded for all $1<p<\infty$.

With the substitution $s:=g\left((1-r)^{-1}\right)$, i.e., $1-r=1 / g^{-1}(s)$, we get for the modulus of imaginary part (the real part is treated in the same way)

$$
\begin{aligned}
& \left|\operatorname{Im}\left(\int_{1-2^{-m}}^{\rho} r^{n-1} \psi(r)\left(1-r^{2}\right)^{\lambda} \mathrm{d} r\right)\right| \\
& \quad=\left|\int_{1-2^{-m}}^{\rho}(1-r)^{-1+\lambda} g^{\prime}\left((1-r)^{-1}\right) \sin \left(\pi g\left((1-r)^{-1}\right)\right) \mathrm{d} r\right| \\
& \quad=\left|\int_{g\left(2^{m}\right)}^{g\left((1-\rho)^{-1}\right)}\left(g^{-1}(s)\right)^{-1-\lambda} \sin (\pi s) \mathrm{d} s\right|
\end{aligned}
$$

Let $a$ be the smallest integer such that $a>g\left(2^{m}\right)$ and $b$ the largest integer such that $g\left((1-\rho)^{-1}\right)>b$. Then, as $\frac{1}{g^{-1}(s)}$ is decreasing, we get

$$
\begin{aligned}
& \left|\int_{g\left(2^{m}\right)}^{g\left((1-\rho)^{-1}\right)}\left(g^{-1}(s)\right)^{-1-\lambda} \sin (\pi s) \mathrm{d} s\right| \\
& \quad \leq\left|\int_{g\left(2^{m}\right)}^{a}\left(g^{-1}(s)\right)^{-1-\lambda} \sin (\pi s) \mathrm{d} s\right| \\
& \quad+\left|\sum_{k=a}^{b-1}(-1)^{k} \int_{k}^{k+1}\left(g^{-1}(s)\right)^{-1-\lambda}\right| \sin (\pi s)|\mathrm{d} s| \\
& \quad+\left|\int_{b}^{g\left((1-\rho)^{-1}\right)}\left(g^{-1}(s)\right)^{-1-\lambda} \sin (\pi s) \mathrm{d} s\right| \\
& \quad \leq 3 \cdot 2^{-m(1+\lambda)},
\end{aligned}
$$

since we have the lower estimate $g^{-1}(s) \geq 2^{m}$ on all integration intervals and the series in the middle term is alternating with decreasing absolute values
of the terms. This shows that $T_{\psi}$ is bounded. Similarly, replacing $g$ by

$$
g:[1, \infty) \rightarrow[0, \infty), \quad g(x)=\int_{1}^{x} f(y) y^{\lambda} \mathrm{d} y
$$

and choosing $\psi$ as above, yields a compact Toeplitz operator $T_{\psi}$. These examples show that symbols of arbitrary growth can induce bounded or even compact Toeplitz operators.

Open Access. This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons licence, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons licence, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons licence and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this licence, visit http:// creativecommons.org/licenses/by/4.0/.

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

## References

[1] Axler, S., Bourdon, P., Ramey, W.: Harmonic Function Theory. Graduate Texts in Mathematics, vol. 137, 2nd edn. Springer, New York (2001)
[2] Grudsky, S., Karapetyants, A., Vasilevski, N.: Toeplitz operators on the unit ball in $\mathbb{C}^{n}$ with radial symbols. J. Oper. Theory 49, 325-346 (2003)
[3] Jevtić, M., Pavlovi'c, M.: Harmonic Bergman functions on the unit ball in $\mathbb{R}^{n}$. Acta Math. Hung. 85, 81-96 (1999)
[4] Liu, C., Shi, J.: Invariant mean-value property and $\mathscr{M}$-harmonicity in the unit ball of $\mathbb{R}^{n}$. Acta Math. Sin. (Engl. Ser.) 19(1), 187-200 (2003)
[5] Lusky, W., Taskinen, J.: Toeplitz operators on Bergman spaces and Hardy multipliers. Studia Math. 204(2), 137-154 (2011)
[6] Miao, J.: Toeplitz operators on harmonic Bergman spaces. Integral Equ. Oper. Theory 27, 426-438 (1997)
[7] Nishio, M., Tanaka, K.: Harmonic Bergman kernels and Toeplitz operators on the ball with radial measures. Rev. Roumaine Math. Pures Appl. 62, 155-169 (2017)
[8] Stroethoff, K.: Compact Toeplitz operators on weighted harmonic Bergman spaces. J. Aust. Math. Soc. Ser. A 64(1), 136-148 (1998)
[9] Taskinen, J., Virtanen, J.A.: Toeplitz operators on Bergman spaces with locally integrable symbols. Rev. Mat. Iberoam. 26, 693-706 (2010)
[10] Taskinen, J., Virtanen, J.A.: On generalized Toeplitz and little Hankel operators on Bergman spaces. Arch. Math. (Basel) 110(2), 155-166 (2018)
[11] Taskinen, J., Virtanen, J.A.: Weak BMO and Toeplitz operators on Bergman spaces. New York J. Math. (in press) arXiv:2106.11734
[12] Yan, F., Zheng, D.: Bounded Toeplitz operators on Bergman space. Banach J. Math. Anal. 13(2), 386-406 (2019)
R. Hagger and J. A. Virtanen ( $\boxtimes$ )

Department of Mathematics and Statistics
University of Reading
Reading RG6 6AX
England, UK
e-mail: j.a.virtanen@reading.ac.uk
R. Hagger
e-mail: r.t.hagger@reading.ac.uk
C. Liu

School of Mathematical Sciences
University of Science and Technology of China
Hefei 230026 Anhui
People's Republic of China
e-mail: cwliu@ustc.edu.cn
J. Taskinen

Department of Mathematics and Statistics
University of Helsinki
P.O 68 (Pietari Kalmin katu 5)

00014 Helsinki
Finland
e-mail: jari.taskinen@helsinki.fi
Received: November 9, 2020.
Revised: February 8, 2021.
Accepted: February 9, 2021.


[^0]:    The first author was supported by the European Union's Horizon 2020 research and innovation programme under the Marie Sklodowska-Curie Grant Agreement No 844451. The fourth author was supported in part by Engineering and Physical Sciences Research Council (EPSRC) Grant EP/T008636/1.

