

# Contrasting interannual and multidecadal NAO variability

Article

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## <sup>1</sup> Contrasting interannual and multidecadal NAO

## <sup>2</sup> variability

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Abstract Decadal and longer timescale variability in the winter North At-7 lantic Oscillation (NAO) has considerable impact on regional climate, yet it 8 remains unclear what fraction of this variability is potentially predictable. This 9 study takes a new approach to this question by demonstrating clear physical 10 differences between NAO variability on interannual-decadal (<30 year) and 11 multidecadal (>30 year) timescales. It is shown that on the shorter timescale 12 the NAO is dominated by variations in the latitude of the North Atlantic 13 jet and storm track, whereas on the longer timescale it represents changes in 14 their strength instead. NAO variability on the two timescales is associated 15 with different dynamical behaviour in terms of eddy-mean flow interaction, 16 Rossby wave breaking and blocking. The two timescales also exhibit different 17 regional impacts on temperature and precipitation and different relationships 18 to sea surface temperatures. These results are derived from linear regression 19 analysis of the Twentieth Century and NCEP-NCAR reanalyses and of a high-20 resolution HiGEM General Circulation Model control simulation, with addi-21 tional analysis of a long sea level pressure reconstruction. Evidence is presented 22 for an influence of the ocean circulation on the longer timescale variability of 23 the NAO, which is particularly clear in the model data. As well as provid-24 ing new evidence of potential predictability, these findings are shown to have 25 implications for the reconstruction and interpretation of long climate records. 26

- 27 Keywords North Atlantic Oscillation · Jet variability · Atmosphere-ocean
- 28 interaction · Climate reconstructions

#### 29 1 Introduction

As the leading pattern of atmospheric circulation variability over the North 30 Atlantic, the North Atlantic Oscillation<sup>1</sup> (NAO) has a strong influence on sur-31 face climate across the Atlantic basin and beyond (Thompson and Wallace 32 2001). Interest in the NAO has been partly motivated by the prominence of 33 its decadal-scale variability in winter (Stephenson et al. 2000). The increase 34 of the winter NAO index from the 1960s to the 1990s gained particular atten-35 tion (Hurrell 1995) but decadal variability is also evident in longer records of 36 the NAO (Pinto and Raible 2012). In the likely absence of atmospheric mem-37 ory from one winter season to the next, influences from other components of 38 the climate system may have played a role. Evidence has been provided of 30 possible influences such as the extratropical (Rodwell et al. 1999; Czaja and 40 Frankignoul 1999; Mosedale et al. 2006; Gastineau and Frankignoul 2012) or 41 tropical oceans (Hoerling et al. 2001; Selten et al. 2004; Greatbatch et al. 42 2012), the sea-ice (Deser et al. 2004; Bader et al. 2011) and also forcings act-43 ing via the stratosphere, such as changes in stratospheric water vapour (Joshi 44 et al. 2006) or solar variability (Ineson et al. 2011). The associated potential 45 for predictability of NAO variability continues to drive research in this area 46 (Folland et al. 2012). 47

<sup>48</sup> Much recent work has focused on shorter, intraseasonal timescales in at-<sup>49</sup> tempts to understand the atmospheric dynamics underlying NAO variability,

<sup>&</sup>lt;sup>1</sup> Or equivalently, the Arctic Oscillation or Northern Annular Mode (Feldstein and Franzke 2006).

following Feldstein (2003) and Benedict et al. (2004). It is clear that some of 50 the NAO variability on decadal timescales could arise from so-called climate 51 noise, in which seasonal sampling of the strong intraseasonal variability can 52 lead to apparent power on interannual and longer timescales (Wunsch 1999; 53 Feldstein 2000b; Schneider et al. 2003; Raible et al. 2005). Various statistical 54 methods have been applied to estimate the fraction of variance on interannual 55 and longer timescales which could be explained simply as climate noise. How-56 ever, these methods differ widely in their findings (Feldstein 2000a; Keeley 57 et al. 2009; Franzke and Woollings 2011), so that the statistical significance of 58 low-frequency NAO variability, and hence the potential for seasonal-decadal 59 predictability is still unclear. In this paper we take a different and complemen-60 tary approach, by searching for physical differences between NAO variability 61 on short and long timescales. 62

The NAO is essentially a description of the preferred structure of variabil-63 ity in the North Atlantic eddy-driven jet stream (Thompson et al. 2002). This 64 deep tropospheric jet represents the net effect of westerly wind forcing by the 65 transient atmospheric eddies (Li and Wettstein 2012), and variations in its 66 strength and position affect regional temperatures and precipitation via vari-67 ations in the prevailing westerly winds and associated storm tracks. The NAO 68 is usually defined via patterns in surface pressure or geopotential height, using 69 methods such as principal component analysis. Physical quantities such as the 70 latitude and speed of the jet are generally not separable by these methods 71 (Monahan et al. 2009), and the NAO reflects variations in both of these quan-72

tities (Woollings et al. 2010). When height fields are linearly regressed onto 73 time series of the jet latitude and speed, both of the resulting spatial pat-74 terns project onto the NAO pattern (Woollings and Blackburn 2012). Despite 75 this, the jet latitude and speed are clearly distinct, having different annual cy-76 cles, power spectra and interannual variability (Woollings and Blackburn 2012; 77 Woollings et al. 2014). This suggests that variations in jet latitude and speed 78 have different physical mechanisms and drivers, and yet they are combined in 79 standard NAO analyses. 80

Here we highlight the contrasting nature of the jet variability associated 81 with the NAO on two different timescales, namely multidecadal and interannual-82 decadal, with periods greater or less than 30 years respectively. This study 83 is related to other approaches which focus on the non-stationarity of the 84 NAO pattern over time (Jung et al. 2003; Lu and Greatbatch 2002; Raible 85 et al. 2006; Wang et al. 2012; Moore et al. 2013), or multi-decadal changes 86 in regime activity (Casty et al. 2005; Franzke et al. 2011). Other studies have 87 highlighted non-stationary relationships between the NAO and regional im-88 pacts on temperature (Pozo-Vázquez et al. 2001; Haylock et al. 2007; Comas-89 Bru and McDermott 2013), precipitation (Vicente-Serrano and López-Moreno 90 2008; Raible et al. 2014) and storm activity (Luo et al. 2011; Lee et al. 91 2012), and the timescale dependence shown here may help in interpreting 92 this non-stationarity. Finally, there is evidence for distinct patterns of ocean-93 atmosphere variability on decadal/multi-decadal timescales in observations 94 and models (Deser and Blackmon 1993; Delworth and Mann 2000; Sutton and 95

<sup>96</sup> Hodson 2003; Shaffrey and Sutton 2006), with non-stationarity or frequency<sup>97</sup> dependence in the relationship between the NAO and sea surface temperatures
<sup>98</sup> (Raible *et al.* 2001; Walter and Graf 2002; Raible *et al.* 2005; Alvarez-Garcia
<sup>99</sup> *et al.* 2008). Hence we also examine the NAO-SST relationship on the two
<sup>100</sup> timescales as a preliminary study of the associated ocean-atmosphere interac<sup>101</sup> tion.

#### $_{102}$ 2 Methods

In this study we focus on the variability in wintertime (DJF) mean data from 103 atmospheric reanalyses. We use both the NCEP-NCAR reanalysis (Kalnay 104 and coauthors 1996) and the Twentieth Century reanalysis (20CR) (Compo 105 et al. 2011). The latter uses mean sea level pressure (MSLP) observations 106 only and takes an ensemble approach to quantify uncertainty, providing 56 107 estimates of atmospheric flow from 1871 to 2012. Unless otherwise stated, the 108 analyses presented here were performed individually for each of these ensemble 109 members and only averaged over the ensemble at the end of the analysis. 110

The NCEP-NCAR reanalysis is used over the period 1950-2012. Since this only provides 62 winters of data this record is short to examine the multidecadal behaviour, but as will be shown results are qualitatively similar to those in 20CR. This reanalysis is used particularly to investigate transient features such as the storm track and the transient eddy fluxes. These quantities have been examined in 20CR but found to give unphysical results in the low-frequency regressions, in particular at high latitudes. This is likely due to issues in the data-sparse period before 1920 (see e.g. Krueger *et al.* (2013)
for a general discussion). The storm track is characterised using data filtered
with a 2-6 day Chebyshev recursive filter to select only the synoptic timescales
(Cappellini 1978).

We also make use of a 100 year present-day control simulation of the high-122 resolution coupled General Circulation Model HiGEM (Shaffrey et al. 2009). 123 This has an atmospheric resolution of  $0.833^{\circ} \times 1.25^{\circ}$  in longitude-latitude 124 with 38 levels, and an ocean resolution of  $\frac{1}{3}^{\circ}$  with 40 levels. This model shows 125 improved simulation of the climatology and variability of North Atlantic cli-126 mate compared to the standard resolution HadGEM1.2 (Shaffrey et al. 2009; 127 Keeley et al. 2012; Hodson and Sutton 2012). Some limited transient diagnos-128 tics have been derived from the HiGEM simulation and these agree well with 129 the results from the NCEP-NCAR reanalysis. 130

The NAO was defined in all datasets as the leading Empirical Orthogonal 131 Function (EOF) and associated principal component time series of monthly 132 mean wintertime (DJF) mean sea level pressure over the Atlantic sector (90°W-133 30°E, 30-90°N). In the 20CR data the NAO was calculated separately in each 134 ensemble member, and the resulting average spatial pattern is shown in Fig-135 ure 5a of Woollings et al. (2014, W14 hereafter). The monthly NAO index was 136 averaged up to seasonal mean values for analysis. As described in W14, indices 137 of jet latitude and speed were derived using the zonal wind at 850 hPa. The 138 method essentially averages the daily zonal wind over 0-60 °W and smoothes 139 it using a 10-day low pass filter before locating the maximum wind speed 140

(Woollings *et al.* 2010). The resulting daily values of jet latitude and speed
were averaged over each winter season to derive seasonal mean values.

To separate the different timescales we apply Empirical Mode Decomposi-143 tion (EMD), as in Franzke and Woollings (2011), to the seasonal mean time 144 series of the NAO and jet indices. This approach empirically decomposes a 145 time series into Intrinsic Mode Functions (IMFs) of different average periods. 146 See Franzke and Woollings (2011) for more description and an example of the 147 method. Here we focus on two timescales: the interannual-decadal, formed by 148 isolating the IMFs with average periods less than 30 years, and the multi-149 decadal, with IMF periods greater than 30 years. The sum of the two filtered 150 time series is exactly equal to the full unfiltered series. These two timescales 151 were chosen after experimentation to best represent the contrasting NAO be-152 haviour (for example the 10-30 year band of timescales behave similarly to the 153 1-10 year band). Note that the general results presented here are reproduced 154 using other filtering methods such as running means, but the EMD results are 155 presented due to their smoothness and objectivity. 156

The general approach taken here is to linearly regress various fields onto the NAO time series at the two different timescales. After averaging the monthly data up to seasonal means, the NAO series is re-normalised so that the series of winter mean values has a mean of zero and a standard deviation of one. Maps therefore show the anomalies associated with one standard deviation of the full unfiltered winter NAO. As described below, this makes the magnitudes of the patterns on the two timescales comparable. However, it is important to note that the long timescale anomaly patterns then have larger amplitude
than is experienced in practise.

#### <sup>166</sup> **3 Jet Characteristics**

We begin by comparing the NAO and jet indices from 20CR in Figure 1. The 167 decadal variability of the NAO is clear, with high NAO values dominating 168 in the early and late twentieth century, and low NAO values dominating in 169 the middle of the century. In contrast, the jet latitude shows mostly interan-170 nual variability, and as shown by W14 it is the jet speed which exhibits the 171 strongest decadal variability. W14 used a Monte Carlo statistical test to assess 172 the probability that the observed variability in the decadal means of these jet 173 series could arise from a white noise process. The results showed that this was 174 quite plausible for the jet latitude (p=0.19) but very unlikely for the jet speed 175 (p=0.01).176

W14 also found that the jet latitude and speed series are uncorrelated (r=-177 0.07), yet both are related to the NAO. This is shown in Figure 2 which cor-178 relates these series with the NAO series on the different timescales obtained 179 using the EMD filtering. Both 20CR and NCEP-NCAR results are plotted, 180 with errorbars reflecting the uncertainty across the ensemble in 20CR. On 181 timescales shorter than 30 years the NAO is dominated by variations in jet 182 latitude. However, on the multidecadal timescale the reverse is true for 20CR 183 at least; the jet speed is more highly correlated with the NAO. NCEP-NCAR 184 shows high correlations for both jet speed and jet latitude on this timescale. If 185

the 20CR analysis is restricted to the time period of the NCEP-NCAR reanalysis, this gives very similar correlations to the NCEP-NCAR data (asterisks in Figure 2), suggesting that this difference is largely due to the short time period.

These correlations suggest a change in the nature of NAO variability on 190 long timescales, with variations in jet speed becoming more important. This 191 impression is confirmed in Figure 3 which shows the 850 hPa zonal wind 192 anomalies associated with NAO variability on the two timescales. On the 193 shorter timescale the wind anomalies generally straddle the mean jet, indi-194 cating a meridional shift, although the anomalies exhibit weaker meridional 195 tilt and are focused downstream of the mean wind maximum. On the mul-196 tidecadal timescale, however, the anomalies overlie the mean jet, indicating 197 a clear increase in jet speed during positive NAO variations. The increase in 198 speed is also shifted towards the eastern end of the jet, highlighting an ex-199 tension of the jet towards central Europe. Similar patterns are seen in the 200 NCEP-NCAR data, despite the difference in correlation on the long timescale 201 in Figure 2. Although the multidecadal anomalies are weaker in NCEP-NCAR, 202 the zonal wind is strengthened along the jet core as seen in 20CR. The same 203 behaviour is also seen in the HiGEM model control simulation, which suggests 204 the result is not a coincidence of the recent observed period. The similarity 205 of the patterns across the three datasets adds considerable confidence to the 206 result. These other datasets will be used in particular to analyse the storm 207 activity and ocean-atmosphere interaction on the two timescales, since these 208

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are two aspects which have proved problematic in the analysis of the 20CR
data.

Following the jet analyses of W14, we performed a simple statistical test 211 of the decadal NAO variability against a white noise hypothesis. For each of 212 the 56 ensemble members, 1000 surrogate white noise NAO time series were 213 generated with the same standard deviation, and then decadal means were 214 calculated from these. The surrogate series were then used to determine the 215 likelihood of the observed level of variability in decadal means occurring from 216 the noise. This analysis showed that the decadal NAO variability in 20CR 217 is very unlikely to occur in a white noise model (p=0.01). We then applied 218 multiple linear regression to express the NAO as a linear combination of the jet 219 indices (which explained 71% of the NAO variance). This enabled us to remove 220 the influence of jet latitude and speed in turn and recalculate the likelihood 221 of the resulting decadal NAO variability. Removing the contribution of jet 222 latitude variations resulted in an NAO series which was still very unlikely in 223 the noise model (p=0.01), but removing the contribution of jet speed gave a 224 value of p=0.19, so that the resulting decadal NAO variability was no longer 225 significantly different from that expected from white noise. This exercise shows 226 that it is the variations in jet speed which are responsible for the elevated power 227 of the NAO on decadal timescales. 228

It would clearly be beneficial to verify the contrasting NAO behaviour on the two timescales in a longer observational dataset. Several attempts have been made to reconstruct atmospheric flow fields beyond the last century using

instrumental records, with considerable success. We have analysed the Küttel 232 et al. (2010) reconstruction of MSLP over Europe and the eastern North At-233 lantic back to 1750, which uses both terrestrial pressure and marine wind data. 234 The result is that this dataset does not exhibit the distinct nature of multi-235 decadal variability shown in the other datasets. Our analysis (described in 236 section 8) suggests that this may be at least partly an artefact of an assump-237 tion of stationarity in the method used to derive the reconstruction. Given 238 the strong agreement between the other datasets, we conclude that the recon-239 struction likely underestimates the timescale dependence of NAO variability. 240

#### 241 4 Regional Impacts

NAO variations are of particular interest because of their strong influence on 242 regional surface climate. These connections have obvious societal impact and 243 are also often used to reconstruct indices of the NAO back in time. Figure 4 244 shows the patterns of near surface air temperature and precipitation associated 245 with the NAO. These impacts are notably different on the two timescales, espe-246 cially over Europe. On interannual timescales this analysis gives the canonical 247 patterns of a quadrupole in temperature anomalies and a north-south dipole 248 in precipitation. On the decadal timescale, however, these patterns are shifted 249 south, so that both temperature and precipitation anomalies are focused on 250 western-central Europe. This southward shift is consistent with the role of the 251 East Atlantic pattern (the second EOF) which also describes changes in jet 252 speed and can be interpreted as acting to shift the NAO circulation pattern 253

north and south (Woollings *et al.* 2010). Figure 4 also shows strong differences
in Arctic climate on the two timescales, though this should be treated with
caution. If the analysis is restricted to the first 100 years of the period (18711970) then the Arctic signal is greatly reduced. This suggests that this signal
arises from a correlation with the recent Arctic warming trend, which may be
unrelated to the NAO.

Some consideration of the variance of the NAO on the different timescales 260 should be taken in interpreting these impacts. This applies to all of the re-261 gression maps shown in this paper. As described above, the NAO time series 262 was normalised and then split into the two components. The variance of the 263 full time series is 1.0, but the two components have very different variances: 264 0.90 for the interannual-decadal series and only 0.08 for the multidecadal se-265 ries. This means that anomalies of the magnitude of those in the multidecadal 266 regressions are never realised in practise; the units of these regressions are 267 changes per standard deviation of the full NAO series, while the multidecadal 268 changes are much smaller than this. Despite this, the anomalies are of consid-269 erable importance when compared to the level of variability on this timescale. 270 Over western Europe the multidecadal NAO regression accounts for up to 271 50% of the variance in decadal mean zonal wind, and similarly 30% of the 272 temperature and 60% of the precipitation variance. 273

#### 274 5 Hemispheric connections

In this section we investigate whether the NAO on the two timescales has different links to remote regions, in particular the Pacific. This is motivated by discussion over the hemispheric or regional nature of the NAO (Wallace 2000) and also by evidence that interaction between the Atlantic and Pacific sectors might be non-stationary or timescale dependent (Raible *et al.* 2001; Castanheira and Graf 2003; Pinto *et al.* 2011; Lee *et al.* 2012).

Figure 5 shows the MSLP associated with NAO variations on the two 281 timescales. Despite the clear differences in jet behaviour, the MSLP patterns 282 in the Atlantic are only subtly different, indicating that MSLP anomalies are 283 hard to interpret in terms of jet characteristics. The most robust difference 284 is an eastward shift of the equatorward Atlantic centre of action on the long 285 timescale. The two reanalyses and the model are in good agreement over the 286 structure of variability on the shorter timescale, including a weak centre of 287 action in the eastern North Pacific. Climate models have historically overes-288 timated the NAO teleconnection to the North Pacific (McHugh and Rogers 289 2005), but HiGEM appears to perform well in this regard (at least on the short 290 timescale). On the longer timescale there is little agreement between the three 291 datasets in the Pacific sector (though again the two reanalyses are similar if 292 only the NCEP-NCAR period is used; not shown). This lack of agreement be-293 tween datasets limits the confidence we can have in hemispheric connections 294 on the long timescale. 295

To further investigate the Atlantic-Pacific links in the observations, we 296 analyse the storm track variability in the NCEP-NCAR reanalysis. The upper 297 level storm track is summarised by the mean of  $v^{\prime 2}$  at 200 hPa, which is re-298 gressed onto the NAO in Figure 6. In the Atlantic sector the contrast in storm 299 track behaviour is very clear. The positive phase of the NAO is associated with 300 a northward shift and extension of the storm track on interannual-decadal 301 timescales, whereas it is associated with a strengthening of the storm track on 302 multidecadal timescales. These features do extend upstream into the Pacific 303 on both timescales. This is particularly clear on the multidecadal timescale, 304 where a strong increase in storm activity is seen over the eastern North Pacific. 305 This is consistent with the results of Lee et al. (2012) who found similar long 306 term changes in the Atlantic and Pacific storm tracks over recent decades. 307 Although weaker, the Pacific storm track signal on the shorter timescale is 308 again consistent with the Atlantic flow, since it indicates a weakening on the 309 southern side of the storm track. It is also consistent with studies which have 310 noted a latitudinal shift of the Pacific storm track accompanying an Atlantic 311 shift (Franzke et al. 2004; Strong and Magnusdottir 2008). Figure 6 also shows 312 corresponding results from HiGEM. As in the MSLP analysis, there seems to 313 be a Pacific-Atlantic storm track link on the shorter timescale which agrees 314 well with that in the reanalysis. On the long timescale there is good agreement 315 between the model and the reanalysis over the Atlantic sector, but not over 316 the Pacific. 317

To summarise, there is evidence that NAO variability on both timescales has links to the Pacific sector. This is particularly clear on the shorter timescale, where both MSLP and storm tracks show good agreement between the different datasets. In contrast, confidence in Atlantic-Pacific links on the longer timescale is limited by the large differences between the three datasets.

#### 323 6 Eddy-mean flow interaction

In this section we present further dynamical diagnostics of the atmospheric cir-324 culation differences on the two timescales. We use the NCEP-NCAR reanalysis 325 for this analysis because of higher confidence in its transient fields. Figure 7 326 shows the vertical structure of the zonal wind anomalies along a section at 327 30 °W. On both timescales the wind anomalies are equivalent barotropic with 328 maxima in the upper troposphere. The differences between the two timescales 329 seen at 850 hPa are clearly evident in the eddy-driven jet through the depth 330 of the troposphere, and are not just surface features. This suggests that the 331 wind anomalies are accompanied by changes in transient eddy driving of the 332 zonal flow, as expected from the storm track changes shown in Figure 6. Fig-333 ure 7 also shows an interesting contrast in subtropical jet variability. This is 334 opposite to the eddy-driven jet, in that the subtropical jet strengthens and 335 weakens on the short timescale but shifts meridionally on the long timescale. 336 Transient baroclinic eddies influence the large-scale flow via both heat 337 and momentum fluxes. The top panels of Figure 8 show the lower tropo-338 spheric transient eddy heat fluxes (v'T'). As in the other fields there is a clear 339

change between the timescales from a largely shifting pattern of variability to a strengthening one. The transient eddy heat fluxes drive a residual overturning circulation with Coriolis torque acting to accelerate the westerly flow at the latitude of the maximum in v'T'. Figure 8 then shows that the changes in transient eddy heat fluxes act to support the zonal flow variations in each case, helping to shift the surface westerlies in the interannual-decadal variability and strengthen them in the multidecadal variability.

To summarise the effects of the transient eddy momentum fluxes, we follow 347 Raible et al. (2010) in calculating  $\mathbf{E} \cdot \mathbf{D}$  where  $\mathbf{E} = ((v'^2 - u'^2)/2, -u'v')$  is 348 similar to the E-vector of Hoskins *et al.* (1983) and  $\mathbf{D} = (U_x - V_y, V_x + U_y)$ 349 is the deformation vector of the time mean flow. Here u' and v' are the 2-6 350 day band-pass filtered wind components and U and V are the wind compo-351 nents averaged over the relevant winter season. This diagnostic describes the 352 exchange of kinetic energy between the eddies and the background flow (Mak 353 and Cai 1989). Regressions of  $\mathbf{E} \cdot \mathbf{D}$  on the NAO are shown in the lower panels 354 of Figure 8. The climatology of  $\mathbf{E} \cdot \mathbf{D}$  features positive values over North Amer-355 ica, implying that eddies grow there at the expense of the background state. 356 Over the Atlantic Ocean the climatology is negative, showing that the eddies 357 lose kinetic energy to the background state there. The regression patterns are 358 again very different on the two timescales. The multidecadal regression shows 359 a strengthening of the conversion from eddy to background state kinetic en-360 ergy, consistent with increased eddy driving of the stronger jet stream. On 361 the interannual-decadal timescale the pattern is more complicated. While the 362

region of maximum eddy forcing is shifted northward by the anomalies, the pattern also shows a meridional tightening of the eddy forcing over the ocean and a general strengthening downstream. The upstream part of this pattern may be related to the strengthening of the subtropical jet seen in Figure 7.

The effect of eddy forcing on the mean flow of the NAO is increasingly de-367 scribed with regard to the breaking of transient Rossby waves (Benedict et al. 368 2004; Franzke et al. 2004; Rivière and Orlanski 2007; Martius et al. 2007; Kunz 369 et al. 2009; Archambault et al. 2010). Here we use the index of Barnes and 370 Hartmann (2012) which identifies wave breaking via the latitudinal overturn-371 ing of vorticity contours. The index outputs the centroid of the wave breaking 372 event, counting each event once only (with a median lifetime of events of two 373 days) and discriminates between cyclonic and anticyclonic breaking based on 374 the morphology of the overturning region. Regressions of the occurrence of 375 wave breaking on the NAO are shown in Figure 9. On the interannual-decadal 376 timescale, the positive NAO is associated with a reduction in cyclonic wave 377 breaking on the poleward flank of the jet (to the south of Greenland) and an 378 increase in wave breaking of both types on the equatorward flank of the jet. 379 There is also a decrease in anticyclonic breaking in the subtropics, suggest-380 ing that this region of wave breaking shifts north along with the jet. These 381 patterns are consistent with the picture that Rossby wave breaking acts to de-382 celerate the westerly winds locally, so that breaking on the equatorward side 383 pushes the jet polewards and vice versa (Gabriel and Peters 2008). 384

On the multidecadal timescale the positive NAO is instead associated with 385 increased wave breaking on both sides of the jet, which is consistent with 386 the strengthening and extension of the jet. Such large-scale conditions are 387 known to foster the occurrence of extreme windstorms over Western Europe 388 (Hanley and Caballero 2012; Gómara et al. 2014). The strongest signals are 389 increased cyclonic breaking to the north and increased anticyclonic breaking 390 to the south, though the two types of breaking also show weaker increases 391 on the opposite side of the jet. On both timescales the behaviour is therefore 392 consistent with Strong and Magnusdottir (2008), in that the latitude of the 393 breaking seems more important than its direction (e.g. the breaking on the 394 equatorward side of the jet may be cyclonic as well as anticyclonic). 395

Finally in this section, we analyse the relationship between the NAO and 396 blocking on both timescales. Blocking is a synoptic situation in which the west-397 erly winds and storm tracks are blocked by a persistent, usually anticyclonic, 398 flow anomaly. Blocking is itself related to wave-breaking (Pelly and Hoskins 399 2003; Altenhoff et al. 2008), though the requirements of spatial scale and per-400 sistence separate it from more transient wave breaking (Masato et al. 2009). 401 Here we define blocking using the index of Scherrer  $et \ al.$  (2006) which is a 402 two-dimensional extension of the classical Tibaldi and Molteni (1990) index. 403 A blocking pattern is identified at a point if 1) the meridional 500 hPa geopo-404 tential height gradient is reversed and 2) the flow is westerly to the north of 405 the point, with a height gradient stronger than 10 m per degree of latitude. A 406 5-day persistence criterion is then applied at each gridpoint. 407

Figure 10 shows the regressions of blocking activity on the NAO. On the 408 interannual-decadal timescale, a positive NAO is associated with strongly re-409 duced blocking over Greenland and the northern North Atlantic, as in Shabbar 410 et al. (2001); Croci-Maspoli et al. (2007) and Woollings et al. (2008). This is 411 also consistent with the reduction in cyclonic wave-breaking seen in Figure 9. 412 The increase in blocking to the south of the jet and over western Europe is 413 also consistent with previous studies (Davini et al. 2013). Essentially the jet 414 shifts southward due to blocking on its northern flank and northward due to 415 blocking on its southern flank. 416

On the multidecadal timescale, blocking anomalies are less strongly re-417 lated to the NAO, with only very weak anomalies at high and low latitudes. 418 The implication is that the effect of blocking is largely to shift the jet stream 419 whereas transient wave breaking can act both to shift or strengthen the jet 420 depending on its position. Interestingly, there is an increase in blocking at the 421 jet exit over the British Isles, despite the strengthening of the westerly winds 422 there under the positive NAO. This may be a consequence of the storm track 423 changes, since a strong storm track upstream is favourable for block main-424 tenance (Shutts 1983; Nakamura and Wallace 1993). Häkkinen et al. (2011) 425 demonstrated multidecadal variability of Atlantic-European blocking associ-426 ated with Atlantic Ocean variability. Such basin-wide variations in blocking 427 do not appear in the analysis presented here, suggesting that the NAO is not 428 a good description of that variability. 429

#### 430 7 Ocean-atmosphere interaction

The distinct physical characteristics of decadal NAO variability suggest an 431 influence external to the atmosphere. The slowly varying ocean circulation is 432 one potential forcing and a natural candidate is the Atlantic variability de-433 scribed by the Atlantic Multidecadal Oscillation (Knight et al. 2005; Sutton 434 and Dong 2012). Correlations of the low-frequency NAO and jet indices with 435 a smoothed AMO index are given in Figure 1. The AMO index was obtained 436 from the NOAA ESRL website and was derived as in Enfield et al. (2001), in-437 cluding detrending and smoothing with a 121 month smoother. Annual means 438 are plotted in Figure 1. The correlations show that the NAO is weakly anti-439 correlated with the AMO and that most of this correlation likely comes from 440 the decadal variability in jet speed, which gives a slightly higher correlation of 441 -0.48. Another potential candidate for ocean forcing of decadal NAO variabil-442 ity is the slow evolution of ocean temperatures in the tropical western Pacific 443 (Kucharski et al. 2006; Manganello 2008). 444

These potential links are investigated in Figure 11 by correlating winter 445 mean sea surface temperatures (SSTs) with the NAO at the two timescales. 446 The SST data comes from the HadISST dataset (Rayner et al. 2003) and the 447 correlation uses the ensemble mean NAO from the complete period of the 448 20CR data. Only gridpoints where the correlation is significant at the 95%449 level are shown. On the short timescale the SSTs show the familiar tripole 450 pattern of anomalies which is largely a response to NAO variability. On the 451 longer timescale the SSTs show a more global pattern, with significant values 452

outside of the Atlantic basin. The North Atlantic is generally cool, as expected
from the negative correlation with the AMO, but the pattern is noisy and the
large values elsewhere are hard to interpret and may not be physically related.
The tropical western Pacific does, however, show a perturbed meridional SST
gradient, as found by Kucharski *et al.* (2006).

Taking a similar approach in correlating the SSTs with the time series of 458 jet speed from 20CR gives a clearer pattern on the multidecadal timescale, 459 comprising a cold subpolar gyre in the North Atlantic and warm anomalies 460 elsewhere, largely confined to the southern hemisphere (lower panels of Fig-461 ure 11). Both of these features are reminiscent of AMO behaviour, and lends 462 support to the potential role of Atlantic Ocean circulation in influencing NAO 463 changes on the multidecadal timescale (e.g. Omrani et al. 2014; Peings and 464 Magnusdottir 2014). There is also indication of a potential influence from the 465 tropical Indian Ocean as suggested by Bader and Latif (2003). 466

In order to provide further evidence of an ocean influence, we examine the 467 ocean-atmosphere coupling in more detail in the HiGEM simulation, where 468 data availability and quality is not a limiting factor. Hodson and Sutton (2012) 469 previously investigated the North Atlantic ocean-atmosphere coupling in this 470 model with a focus on the shorter timescale. Figure 12 shows the correlation 471 of winter (DJF) mean SST with the long timescale NAO for the model. This 472 shows a distinct cold North Atlantic subpolar gyre; a pattern which is at least 473 qualitatively similar to that related to jet speed in the observations (Figure 11). 474 Figure 13a shows time series of the subpolar gyre temperature and the long 475

timescale NAO variability, which show reasonable covariability on the longtimescale.

In order to determine an influence of the ocean on the atmosphere, Fig-478 ure 13b presents the heat content budget for the subpolar gyre region. The 479 sum of the heat contributions due to individual fluxes is shown by the dotted 480 black line, and this agrees well with the total heat content in the thick black 481 line, showing that all terms have been accounted for. The budget shows that 482 it is the ocean heat flux convergence (in red) which is driving the heat content 483 changes of the subpolar gyre, with the atmospheric fluxes (latent, sensible 484 heating) acting to damp the changes in heat content. This agrees with the 485 observational study of Gulev et al. (2013) which shows that ocean-atmosphere 486 surface heat fluxes are driven by the ocean on multidecadal timescales. 487

The variations in ocean heat flux convergence into the subpolar gyre may 488 be driven by a number of factors. Figure 13c demonstrates that these variations 489 are closely related to variations in the Meridional Overturning Circulation at 490 45 °N, suggesting that variations in meridional ocean transport are responsi-491 ble. These variations in turn arise in response to west to east ocean pressure 492 gradient across the Atlantic basin, which is dominated by ocean density vari-493 ations on the deep western Atlantic boundary (Figure 13c: green line). The 494 resulting picture is that the decadal variations in the SST of the subpolar gyre 495 (see Figure 13a) are driven by changes in the MOC in the model, which are 496 in turn driven by variations in the density within the deep western boundary 497 current. Such density variations are ultimately generated at the ocean surface 498

<sup>499</sup> in small regions of intense ocean cooling, such as the Labrador Sea (Marshall
<sup>500</sup> and Schott 1999). NAO variability is thought to be a significant factor in driv<sup>501</sup> ing variations in ocean cooling in these regions (Eden and Willebrand 2001).
<sup>502</sup> Density anomalies generated by this process then slowly propagate southwards
<sup>503</sup> along the western Atlantic boundary.

In summary, subpolar SST anomalies in HiGEM arise due to changes in 504 ocean heat convergence. These subpolar SST anomalies then in turn influence 505 the atmosphere, likely by changing the meridional temperature gradient and 506 hence the baroclinicity across the storm track. With a cold subpolar gyre the 507 meridional gradient is strengthened which is expected to lead to stronger storm 508 activity as seen in Figures 6 and 8. This in turn leads to increased acceleration 509 of the westerly flow (Figure 8) and a stronger jet. Evidence of this mechanism 510 of ocean influence on the atmosphere has been found in the natural variability 511 of other models (Gastineau and Frankignoul 2012), in the context of model 512 biases (Keeley et al. 2012) and in the response of models to climate change 513 (Woollings et al. 2012). 514

#### 515 8 Implications for long climate reconstructions

In this section we analyse the Küttel *et al.* (2010) reconstruction of Atlantic/European MSLP back to 1750 for evidence of contrasting NAO behaviour on the short and long timescales. The NAO in the reconstructed data is defined as the first EOF of winter mean MSLP over the region (0-40 °W, 20-70 °N), which roughly comprises the North Atlantic portion of the data domain. The surface

geostrophic zonal wind  $u_g$  is then derived from the MSLP and this is regressed 521 onto the NAO at the two different timescales. The results are shown in the 522 top panels of Figure 14. In contrast to the other datasets, the differences be-523 tween the two regression patterns are small, with the NAO largely describing 524 a jet shift on both timescales. To test whether this is due to the use of  $u_q$ 525 rather than  $u_{850}$ , we apply the same procedure to derive  $u_q$  from 20CR and 526 the results are shown in the middle panels of Figure 14. The results for 20CR 527 resemble the difference in  $u_{850}$  found between a jet shift on short timescales 528 and an increase in speed on longer timescales (Figure 3), suggesting that  $u_q$ 529 from the reconstruction should be capable of capturing this behaviour. 530

The reconstruction method used in Küttel et al. (2010) is the multivariate 531 principal component regression technique which relies on the assumption of 532 stationarity. EOFs of both observed MSLP fields and pressure-sensitive proxy 533 data (e.g., early measurements and documentary data such as ship log books) 534 are combined with a multiple linear regression technique for the observational 535 period to project local proxy information onto regional patterns. The linear 536 relation is then assumed to be stationary over time and used to reconstruct 537 MSLP fields further back in time. It is possible that the similarity of the 538 NAO regressions on the two timescales is a consequence of this assumption of 539 stationarity. To investigate this possibility we have performed a simple test on 540 the 20CR data, treating it in an analogous way to the reconstruction method. 541

Firstly an EOF analysis of the MSLP is performed on the last 30 years of the 20CR data. Only the first four EOFs are retained, which explain 91% of

the variance of this sample. These EOFs are then used as a basis to truncate 544 the full 136 year dataset: a multiple linear regression technique considering the 545 anomaly maps for each year, the four leading EOFs and the associated principal component time series are used to derive MSLP fields by only including 547 the projection on these four EOFs. The resulting pseudo-reconstructed MSLP 548 fields are then regressed on the NAO at the two timescales and finally  $u_q$  is 549 calculated from these. The results, shown in the lower panels of Figure 14, in-550 dicate some differences between the two timescales but these are substantially 551 weaker than those in the original 20CR analysis (middle panels of Figure 14). 552 This is particularly clear over Europe where the latitude of the wind anomalies 553 is quite different in the middle panels but not in the lower panels. Retaining 554 more than four EOFs (e.g. 10 EOFs which comprises 99% of the variance) does 555 not significantly alter these findings (not shown). While this procedure is anal-556 ogous but not identical to the technique of Küttel et al. (2010), it does suggest 557 that the lack of a distinct multidecadal NAO signature in the reconstruction 558 could be at least partly due to the assumption of stationarity in the method. 559 This effect could be compounded by the relatively low density of proxy data in 560 the jet stream region over the ocean considered in the reconstruction of Küttel 561 et al. (2010), and non-climatic noise intrinsic to proxy data. 562

#### 563 9 Conclusions

This study shows that the multidecadal variability of the NAO represents very different variations in atmospheric circulation from the interannual-decadal variability. The faster variability is dominated by meridional shifts of the jet
stream and associated storm track, while the slower variability is dominated
by changes in the speed of the jet and the strength of the storm track.

Variations on both timescales are supported by forcing from the transient 569 eddies, but the nature of this forcing is different. The interannual-decadal 570 variations are associated with shifts of the transient eddy forcing and with the 571 occurrence of blocking weather patterns. Other work suggests this variability 572 represents variations in the occurrence of different synoptic circulation regimes 573 such as preferred jet positions (W14 and references therein). In contrast, the 574 multidecadal variability is associated with changes in strength of the eddy 575 forcing and with in-phase changes in the occurrence of transient Rossby wave 576 breaking on both sides of the jet. 577

The patterns of influence of the NAO on regional temperatures, wind 578 speeds and precipitation are different on the two timescales, and this has clear 579 implications for the interpretation of proxy or reconstructed records of past at-580 mospheric variability in this region. The variations on multidecadal timescales 581 may not be well represented by the canonical NAO pattern, especially since 582 the shorter timescale variability dominates the variance of the NAO index. A 583 potential example of this has been given, by analysing a long MSLP recon-584 struction. In contrast to the other datasets, this does not exhibit a difference 585 in NAO character on short and long timescales, and it is suggested that the 586 stationarity assumption commonly used in reconstruction methods is at least 587

partly responsible for this. These findings have implications for the interpre tation of climate reconstructions and long climate records.

These results also provide strong evidence for the presence of some forc-590 ing on the decadal NAO from more slowing varying components of the climate 591 system than the atmosphere. Some evidence of links to Atlantic Ocean variabil-592 ity were revealed, although other factors may also contribute. This evidence 593 is particularly clear in the HiGEM GCM, where variations in the Atlantic 594 Meridional Overturning Circulation lead to significant SST anomalies in the 595 subpolar gyre region which are then damped by the heat fluxes to the atmo-596 sphere. 597

For the emerging discipline of decadal prediction these results are an en-598 couraging sign of potential predictability of the winter NAO on multidecadal 599 timescales. Furthermore, the multidecadal component of NAO variability has 600 a clear and distinct influence on surface temperatures and precipitation, es-601 pecially in Europe, so that decadal forecasts of this variability could be of 602 practical use. However, the contrasting behaviour on interannual and decadal 603 timescales suggests that the potential sources of skill may be different for 604 decadal forecasts than for seasonal forecasts. 605

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Fig. 1 Ensemble mean indices of the winter mean NAO, jet latitude and jet speed from 20CR, with the multidecadal (>30 year) component also shown. The shading indicates the  $\pm 2$  standard deviation range across the ensemble. The AMO is shown in the bottom panel, and in each other panel the correlation of the respective low-frequency timescale with this is given.



Fig. 2 Correlations of jet indices with the NAO on short and long timescales, showing NCEP-NCAR results as crosses and 20CR as circles, with errorbars giving the  $\pm 2$  standard deviation range across the ensemble. Asterisks give 20CR results for the NCEP-NCAR period.



U850 regressed on NAO

Fig. 3 Regression patterns of anomalies in 850 hPa zonal wind on the NAO at the two timescales, using 20CR, NCEP-NCAR and HiGEM. The wind climatology is shown in black contours at 7.5 and  $10.5 \,\mathrm{ms}^{-1}$ .



## Surface air temperature regressed on NAO



## Precipitation regressed on NAO



Fig. 4 Regression patterns of near surface air temperature (on the  $\sigma$  =0.995 level) and precipitation on the NAO at the two timescales. This analysis was performed on the ensemble mean fields from 20CR.



Fig. 5 MSLP from both reanalyses and from the HiGEM model regressed onto the NAO on both timescales. Contours are drawn every 1 hPa with negative contours in blue. All ensemble members are used for 20CR and the results averaged.



#### Storm track regressed on NAO (NCEP)

Fig. 6 Regressions of the storm track activity on the NAO in the NCEP-NCAR reanalysis and HiGEM, using the square of the meridional wind anomalies after applying a 2-6 day filter. The climatology is contoured at 60 and 100  $m^2s^{-2}$ .



U(30W) regressed on NAO (NCEP)

Fig. 7 Regression patterns of anomalies in zonal wind at 30  $^{\circ}\mathrm{W}$  on the NAO at the interannual and decadal timescales, using the NCEP reanalysis. The wind climatology is contoured every  $5 \,\mathrm{ms}^{-1}$ .



High-pass vT850 regressed on NAO



Fig. 8 Top: Regressions on the NAO of the 2-6 day v'T' at 850 hPa. The climatology is contoured at 4, 7 and 10 Km s<sup>-1</sup>. Bottom: Regressions on the NAO of the eddy forcing diagnostic  $\mathbf{E} \cdot \mathbf{D}$  at 250 hPa. The climatology is contoured every  $5 \,\mathrm{m^2 s^{-3}}$ , with negative contours dashed and the zero contour omitted. In all cases the data is from the NCEP-NCAR reanalysis.



Fig. 9 Regressions on the NAO of the transient Rossby wave breaking occurrence, split into cyclonic (CWB) and anticyclonic (AWB). The NCEP-NCAR reanalysis is used.

### Blocking regressed on NAO



 ${\bf Fig. \ 10} \ {\rm Regressions \ of \ blocking \ occurrence \ on \ the \ NAO, \ using \ the \ NCEP-NCAR \ reanalysis.}$ 

The climatology is shown in black contours every 2%.



Fig. 11 Correlation value r of winter mean SSTs on the NAO (top) and jet speed (bottom) at the two timescales. Only values which are significant at the 95% level have been plotted.



Fig. 12 Correlation value r of Sea Surface Temperature (SST) correlated with multidecadal NAO index in HiGEM. Shaded areas are significant at the 95% level (p < 0.05). Both SST and NAO were detrended before correlation.



Fig. 13 Decadal NAO variability and Sub Polar Gyre heating in HiGEM Control simulation. A) Black: Mean Atlantic Sub Polar Gyre (SPG) Sea Surface Temperature (SST) (75:0°W, 45:60°N- box in Figure 12). Red: Detrended decadal component of the NAO in HiGEM multiplied by -1, extracted using EMD as for observations. Both indices have been standardized to have unit variance. B) Heat budget for the SPG region. Black solid: upper ocean heat content within the SPG region (0:500 m depth). Other lines - Heat content in the SPG due to: Ocean Heat convergence (Red), Surface Latent (Purple) and Sensible (Green) Heat fluxes and Longwave (Light Blue) and Shortwave (Dark Blue) surface radiation fluxes. All surface fluxes are defined positive into the ocean. Black dotted line: the sum of all contributions to the heat content. All indices have been detrended. Units are  $10^7$  PJ. Black (solid and dotted) lines have been multipled by 2 to aid comparison with SPG SST in panel A. C) Black: Ocean Heat Convergence Flux into the SPG region (45:60 °N). Red: Atlantic Meridional Overturing Circulation (AMOC) at 45 °N (AMOC is the integral of southward meridional ocean velocity between 1000:7000 m across the Atlantic Basin). Green: Mean Ocean Density on the Deep western Atlantic Boundary (1500:3000 m 59:58 °W 44:45 °N). All indices have been detrended and standardized to have unit variance.



Fig. 14 As Figure 3 but showing the surface geostrophic zonal wind using the Küttel reconstruction and the 20CR data. The wind climatology is contoured in black at  $\pm 5$ ,  $7 \text{ ms}^{-1}$ . The reconstruction only covers the region shown. See text for further details.